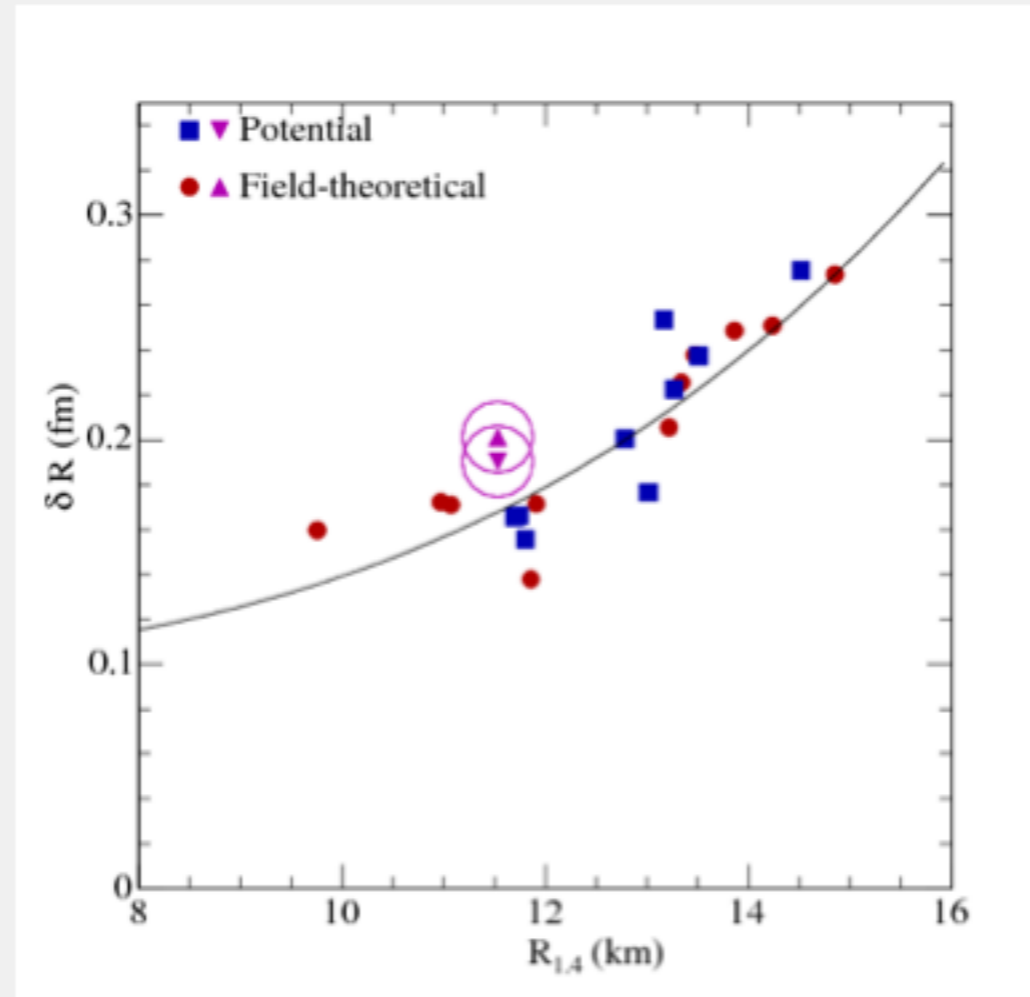


# Observations and the Nuclear Symmetry Energy

Andrew W. Steiner (INT/U. Washington)



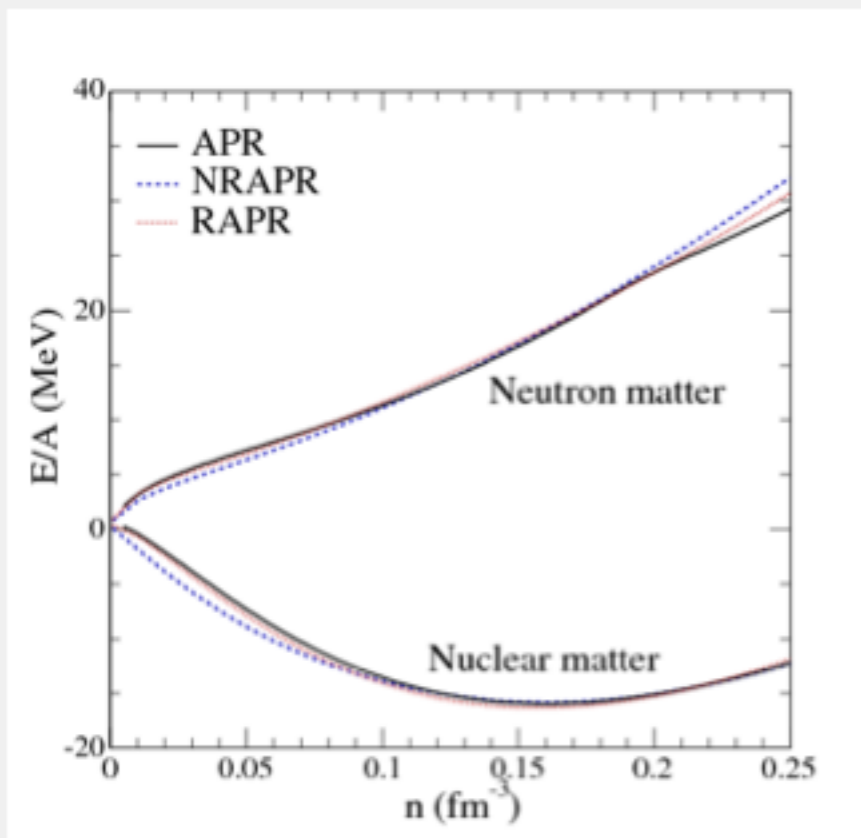
Steiner et al. (2005), adapted from Horowitz and Piekarewicz (2001)

July 18, 2013

With: Edward F. Brown (MSU), Alex T. Deibel (MSU), Tobias Fischer (Wroclaw)  
Matthias Hempel (Basel), Stefano Gandolfi (Los Alamos), and James M. Lattimer  
(Stony Brook)

- Definitions and Diversions
- Neutron Star Basics
  
- Cooling and Direct Urca
- Radii
- Magnetar Flares
- Supernova EOS
- Deep Crustal Heating
- Other Observables

# The Nuclear Symmetry Energy



Steiner et al. (2005)

- Define  $\alpha \equiv (n_n - n_p)/n$

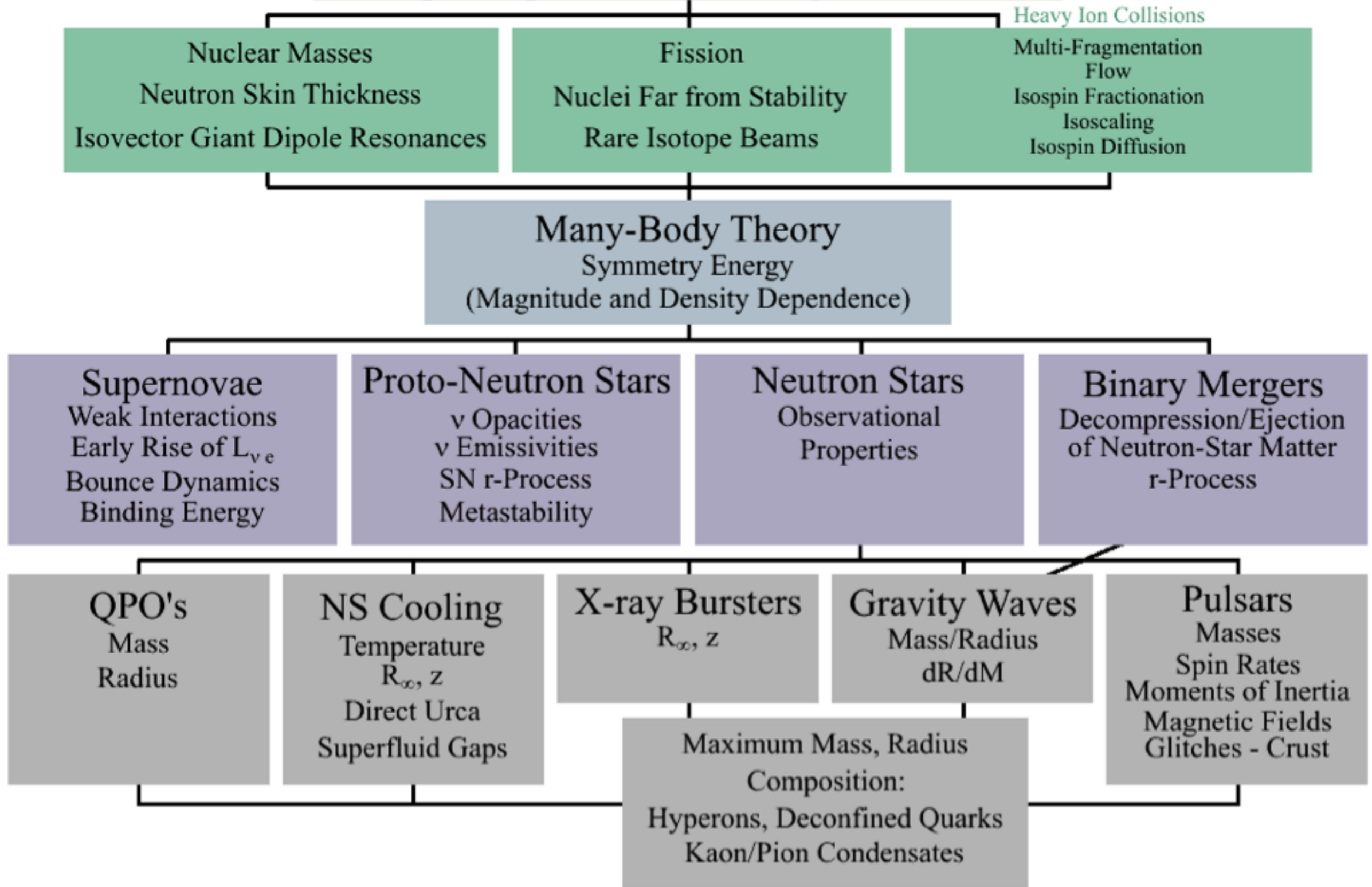
$$S(n) = \frac{1}{2n} \left( \frac{\partial(nE/A)}{\partial\alpha} \right)_{\alpha=0}$$

- $\tilde{S}(n) = (E/A)_{\text{neut}}(n) - (E/A)_{\text{nuc}}(n)$
- If  $E/A$  is quadratic in  $\alpha$ ,  $S(n) = \tilde{S}(n)$
- $S \equiv S(n_0)$
- $L$  is the derivative,  $L \equiv 3n_0 S'(n_0)$

- I define  $S$  and  $L$  entirely from homogeneous matter
- At low-densities one can include clusters
- In either case, I prefer to be consistent
- $4S(n) = \partial_\alpha \left[ \mu_n(n, \alpha) - \mu_p(n, \alpha) \right]$
- Useful definition because there are so many connections...

# There are many correlations...

## Isospin Dependence of Strong Interactions



# Gateway Quantities to the Symmetry Energy

Are  $S$  and  $L$  really the quantities of interest?

- Pressure of neutron matter near and above saturation
  - Easier to compute theoretically
  - Related to neutron stars
- Isovector dependence of the nucleon optical potential
  - Input for heavy-ion collisions
  - Relevant for transport in dense matter
- Isovector response of the ground state of a nucleus
  - Modification of the single particle energies
  - and the density distributions
- Isovector effective mass

Nevertheless, for now I stick with  $S$  and  $L$ .

# Neutron Star Composition

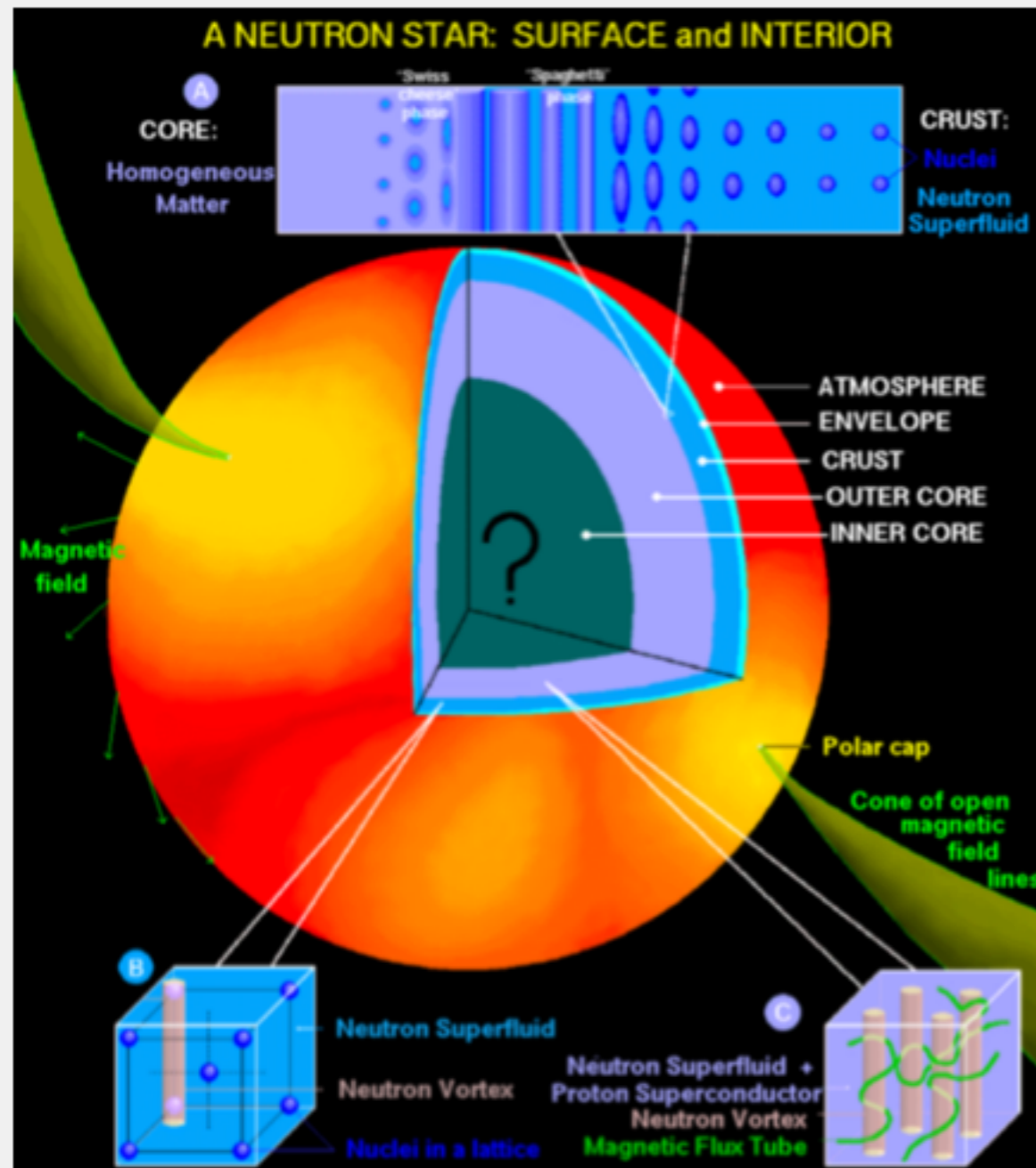


Figure by Dany Page

- In outer crust,  $\mu_e$  increases faster than  $\mu_{n,p}$ , higher densities more neutron-rich
- In inner crust,  $S$  determines EOS of neutron matter as well as properties of nuclei
- As one proceeds into the core  $\mu_{n,p}$  increase faster, tend to restore isospin symmetry
- High  $\mu_e$  can favor phase transitions, i.e.  $\mu_{\pi^-} = \mu_e$
- Relationship with hyperons more complicated
- When strange quarks appear, there is a hypercharge asymmetry energy

# Thermal Emission from Isolated Neutron Stars

$$C_V \frac{dT}{dt} = L_\nu + L_\gamma, \quad L_\gamma \sim T^{2+4\alpha}, \quad L_\nu \sim T^{6-8}, \quad C_V \sim CT$$

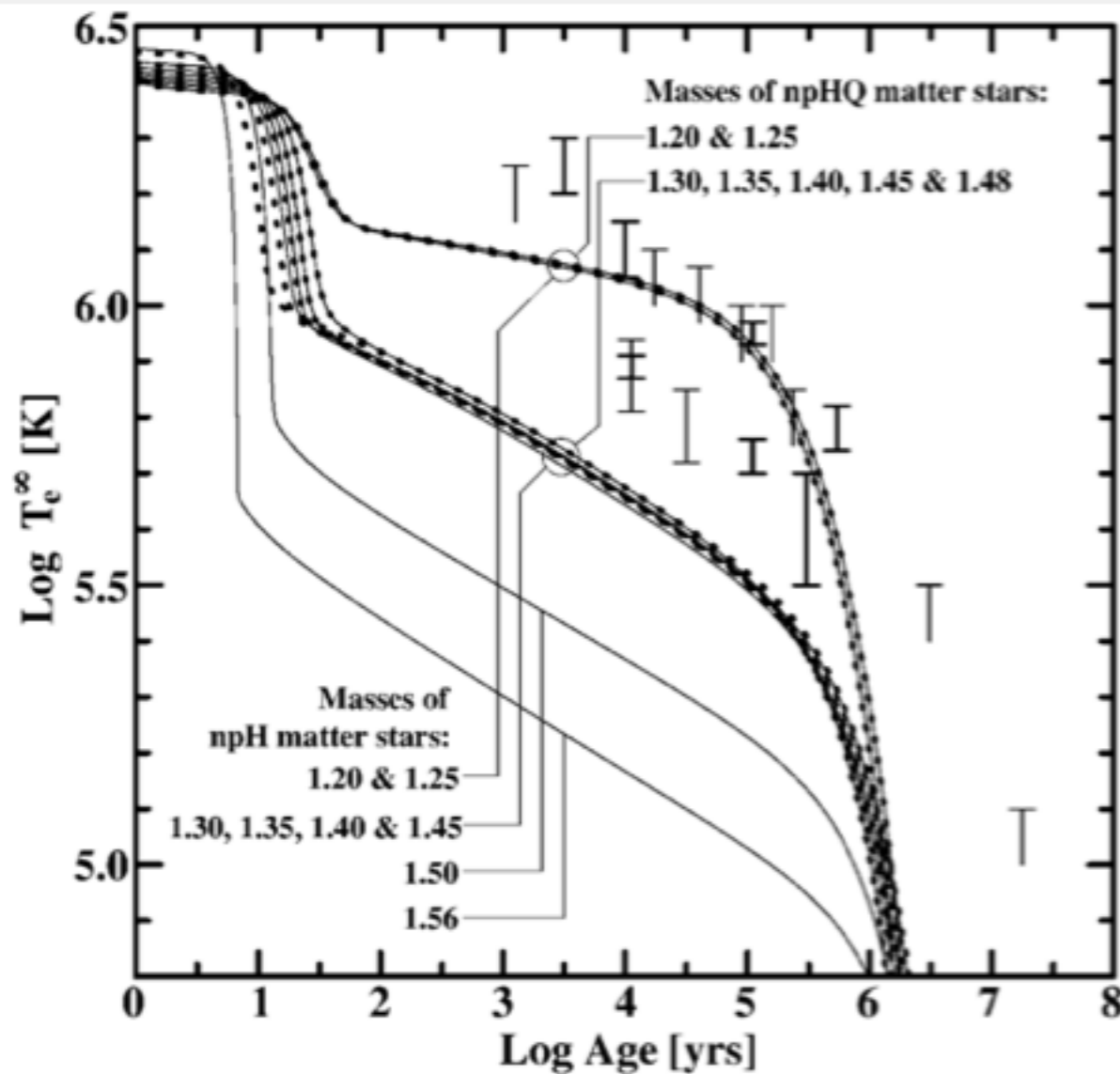


FIG. 3. Cooling of stars with  $npH$  (continuous lines) and  $npHQ$  matter (dotted lines) for various stellar masses (in  $M_\odot$ ).  $n^3P_2$  gaps are from case [c] while quark gaps, when present, are from model [C] of Fig. 1.

Page, et al. (2000)

- Fastest cooling from  $n \rightarrow p + e + \bar{\nu}_e$
- Only possible if enough protons are around

$$2s \equiv p_{Fn} + p_{Fp} + p_{Fe}$$

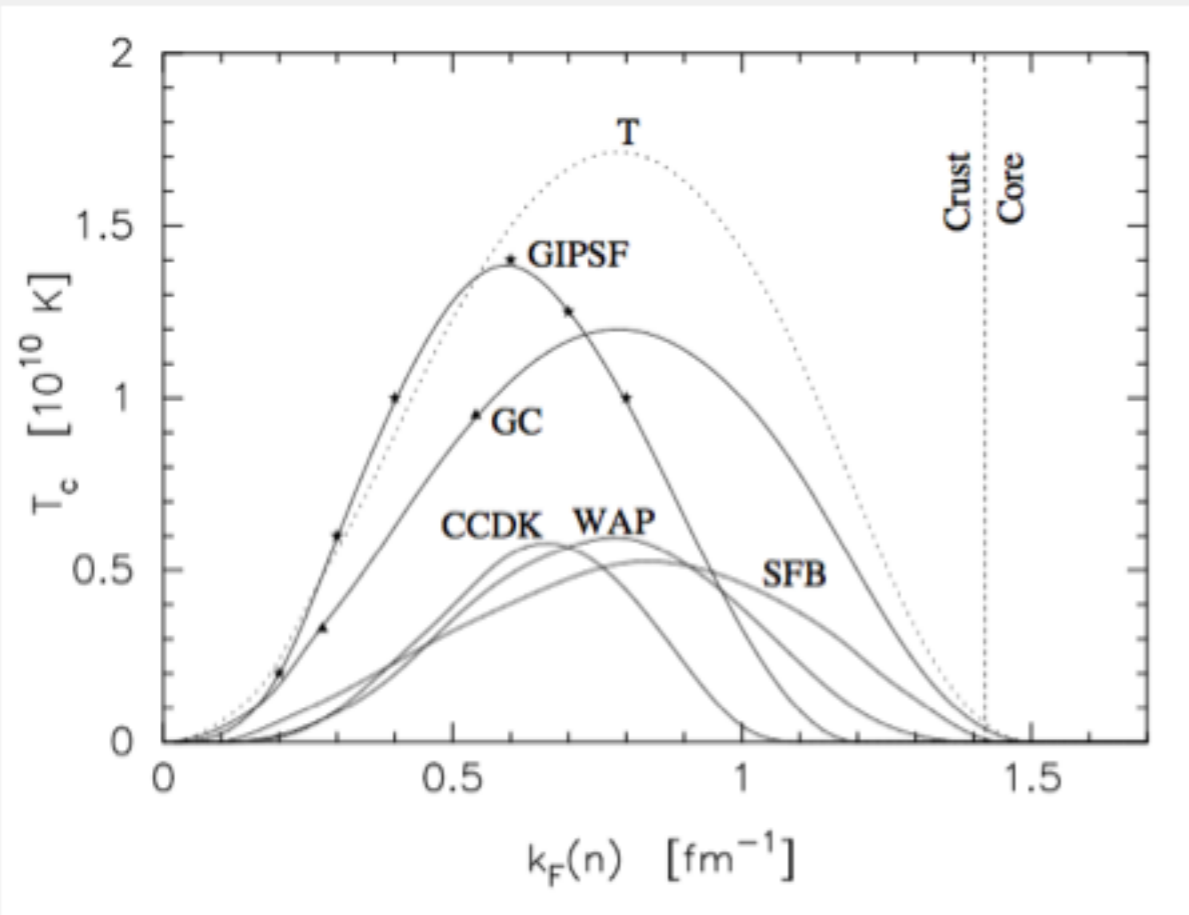
Condition for direct Urca is:

$$s(s - p_{Fn})(s - p_{Fp})(s - p_{Fe}) \geq 0$$

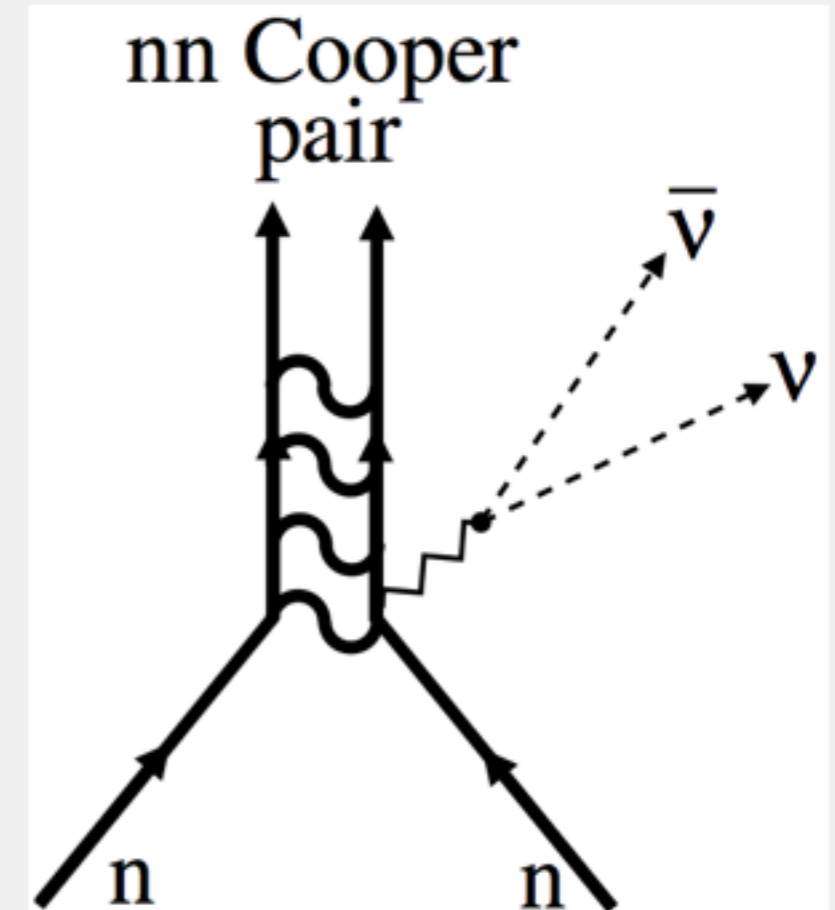
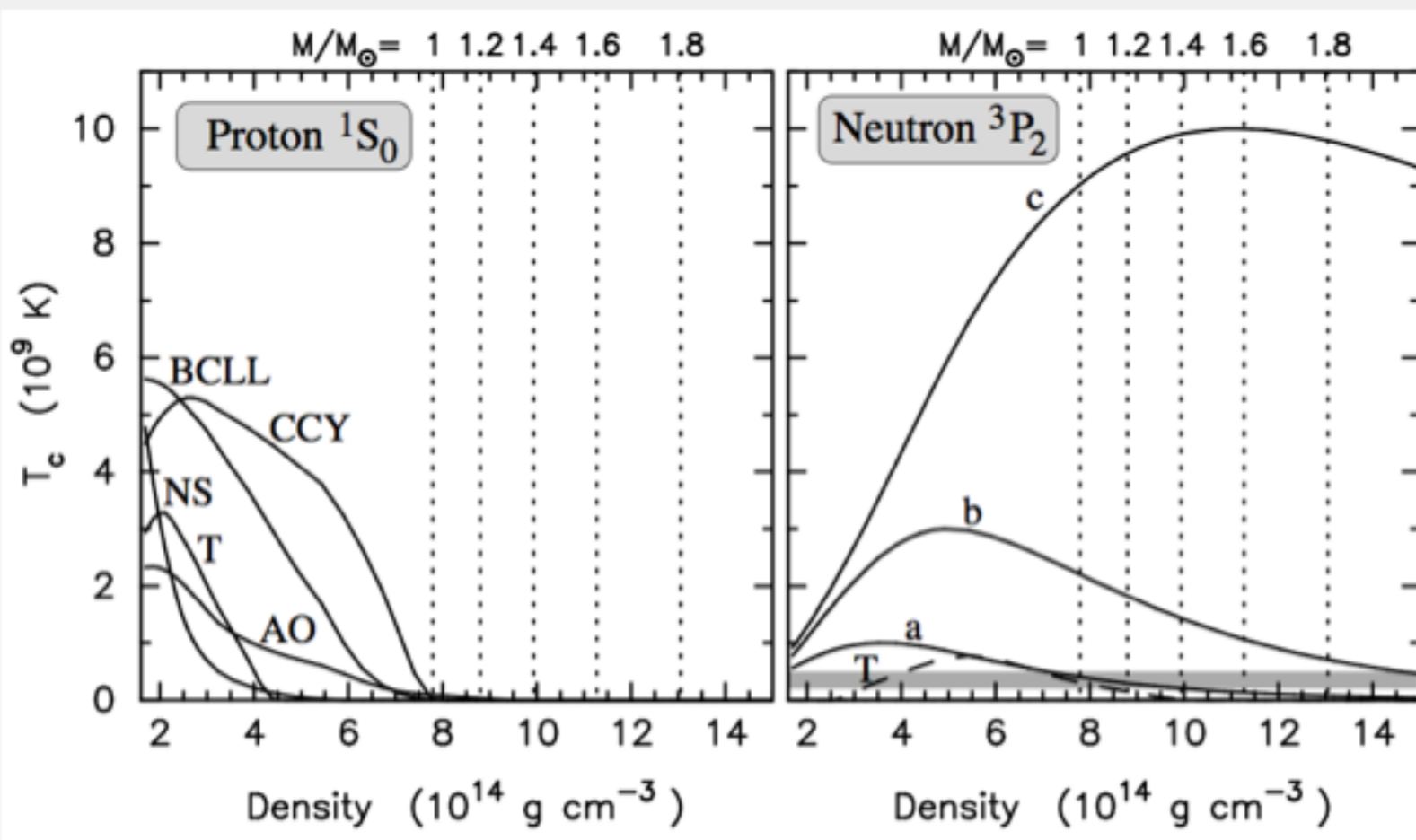
Steiner (2006)

- Alternative:  
 $n + n \rightarrow n + p + e + \bar{\nu}_e$

# Superfluidity and NS Cooling



- Superfluidity can block the direct Urca process
- ...but it opens up new cooling processes
- Superfluidity also implicitly dependent on  $S(n)$





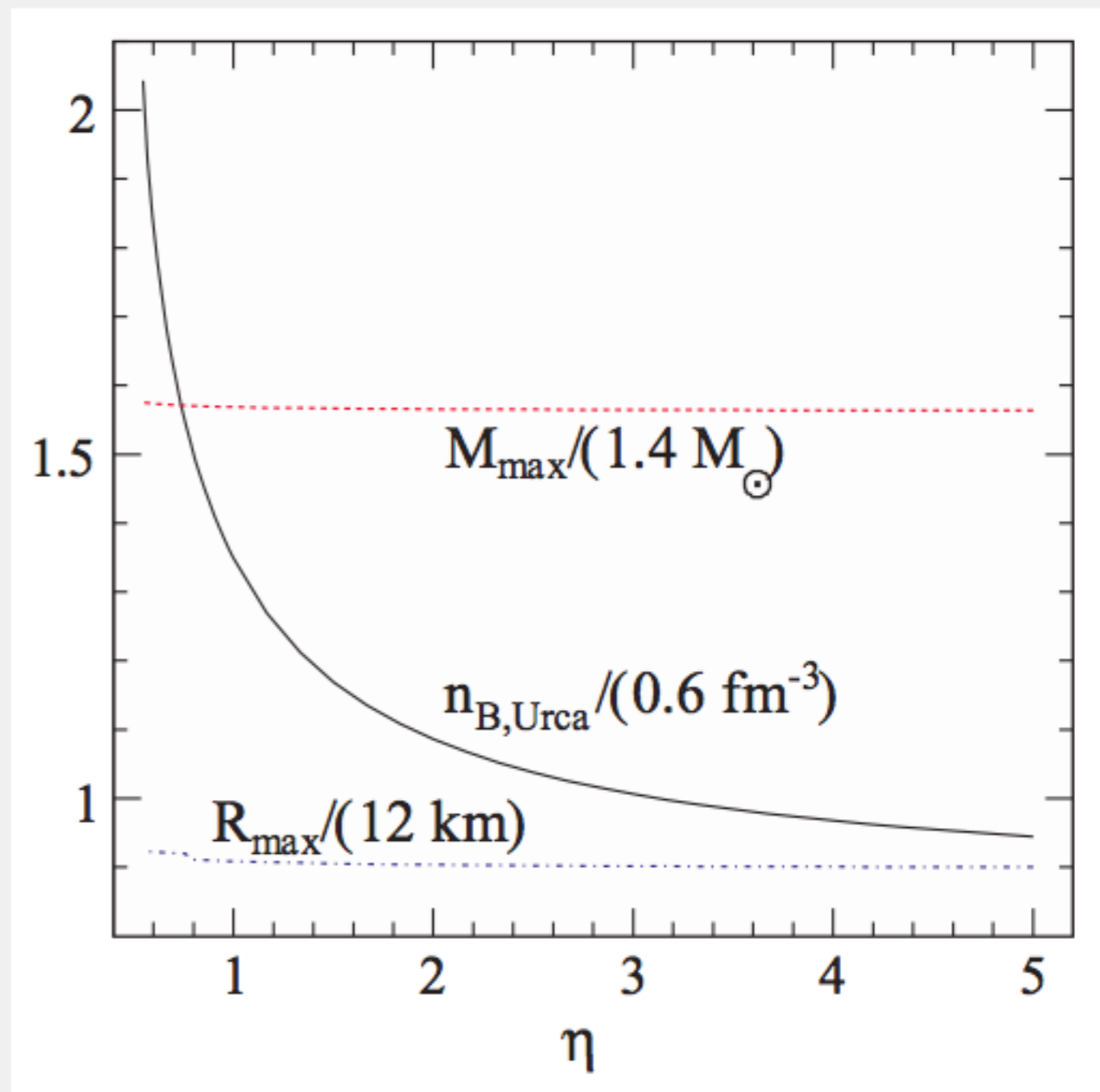
# Quartic Terms and the Direct Urca Process

$$(E/A)(n, \alpha) = (E/A)_{\text{nuc}}(n, \alpha) + \alpha^2 S(n) + \alpha^4 Q(n)$$

- Below saturation, quartic terms are likely "small", above saturation densities, they may be important

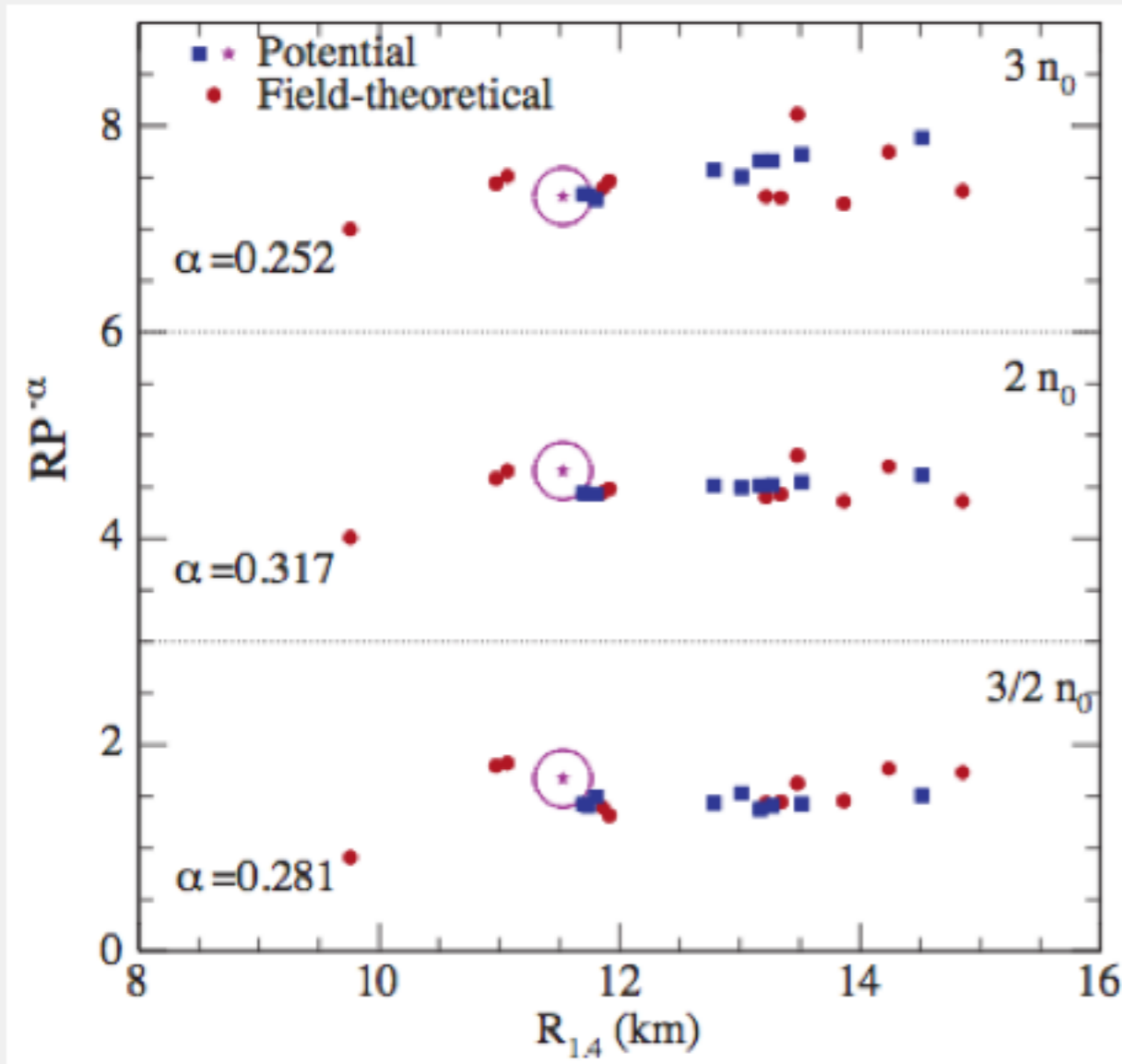
$$\eta(n) \equiv \frac{4S(n) + 5Q(n)}{4S(n) + Q(n)}$$

- $3/7 < \eta(n) < 5$
- Complicates connection between symmetry energy and direct Urca
- Superfluidity also very important, and depends on  $L$

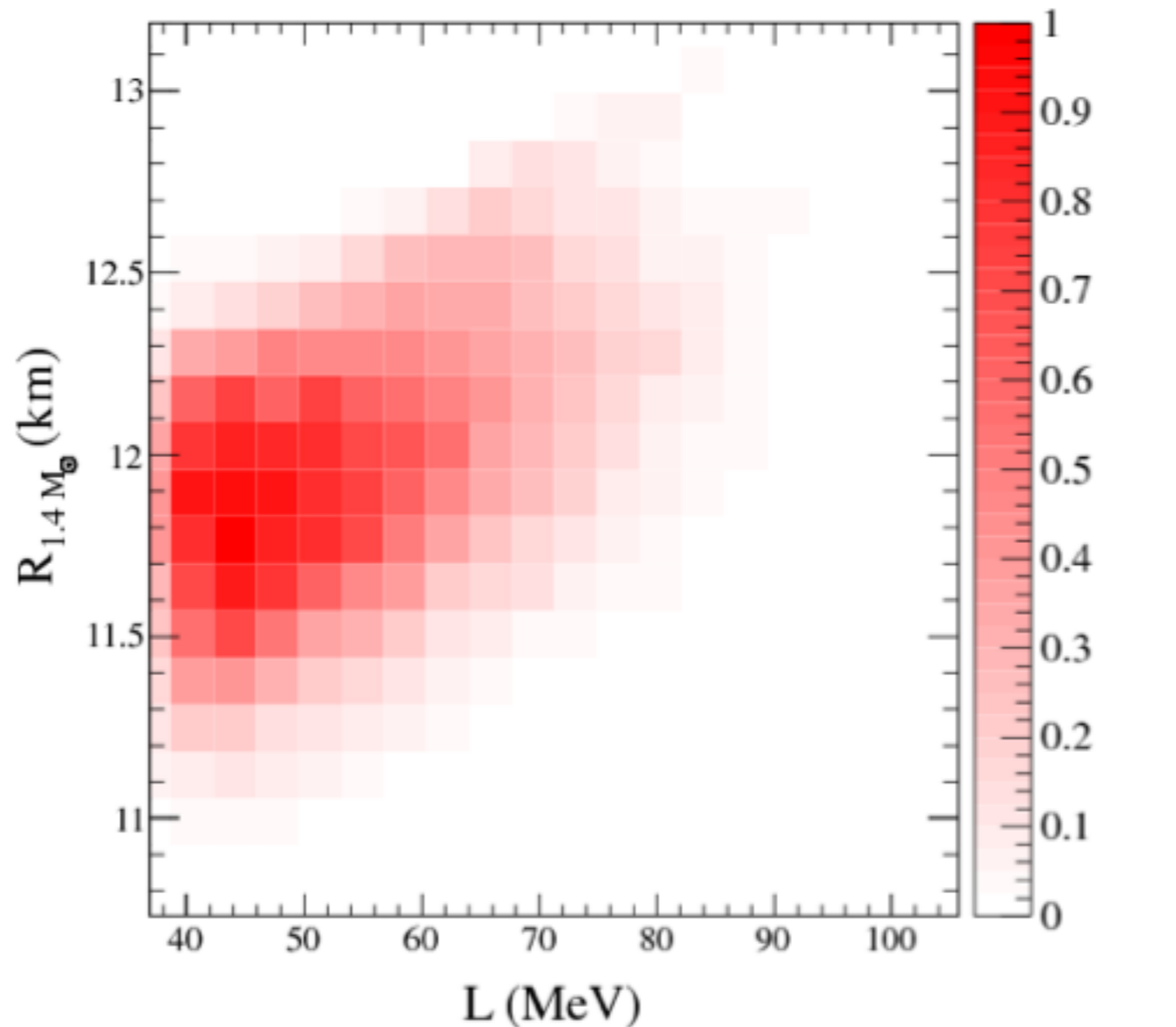


Steiner (2006)

# Connection to Neutron Star Radii



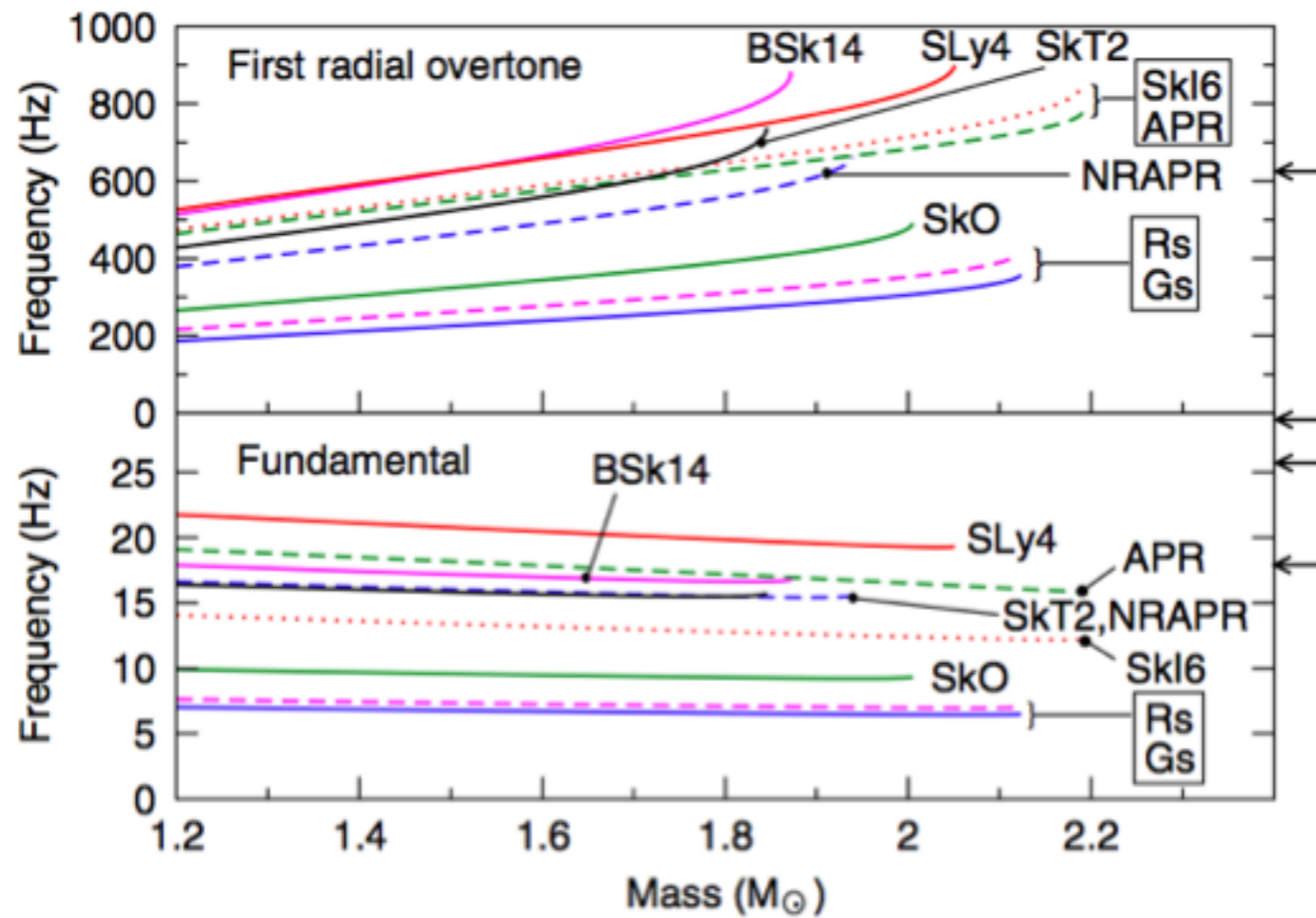
Steiner et al. (2005) based on Lattimer and Prakash (2001)



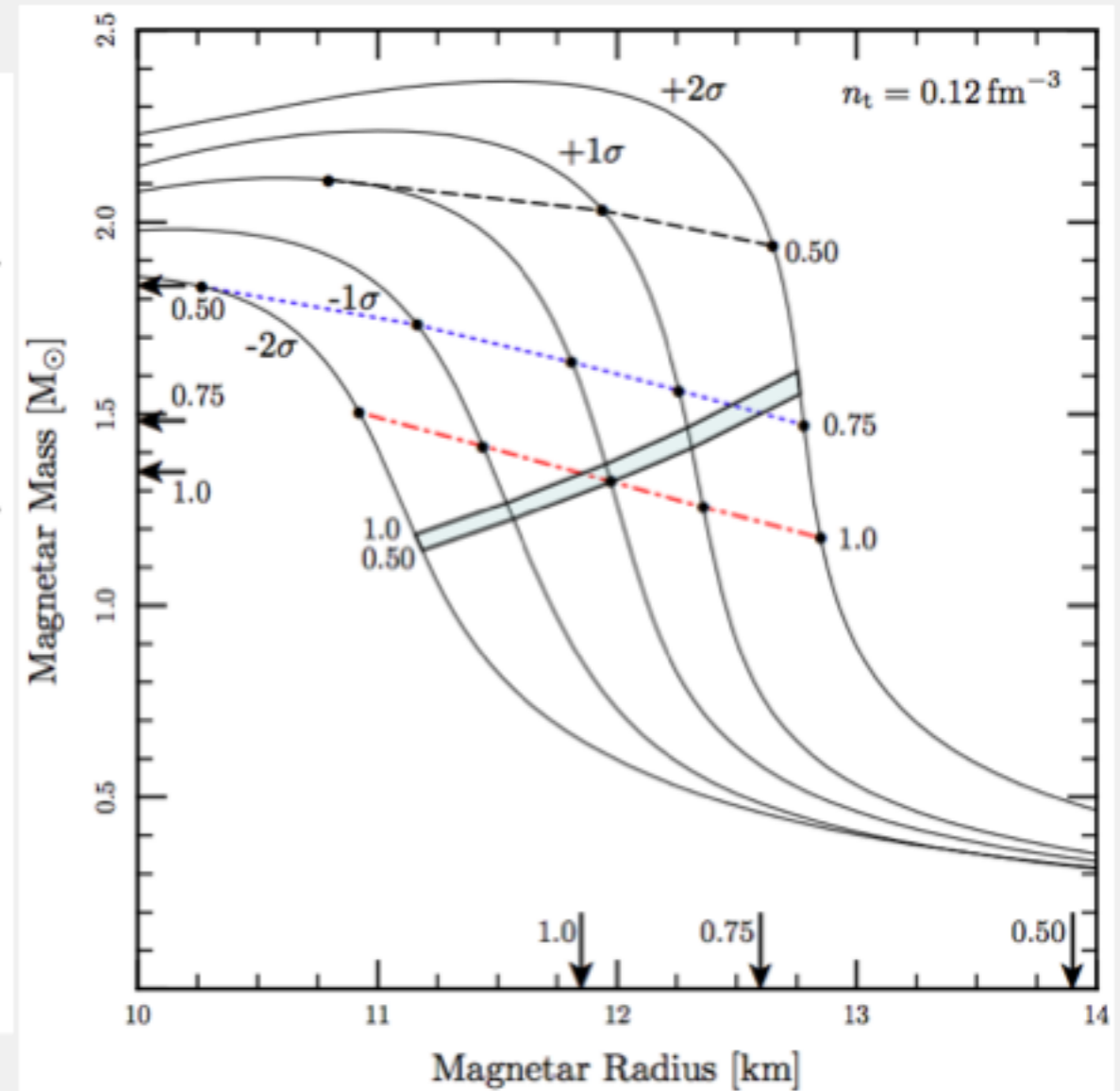
Based on Steiner et al. (2010)

- $L$  correlated with  $R$  for  $1.4 M_\odot$
- Unless there is a strong phase transition at low density
- $L$  not correlated with  $R$  for  $2.0 M_\odot$
- More about this next week...

# Magnetar Flares and the Symmetry Energy



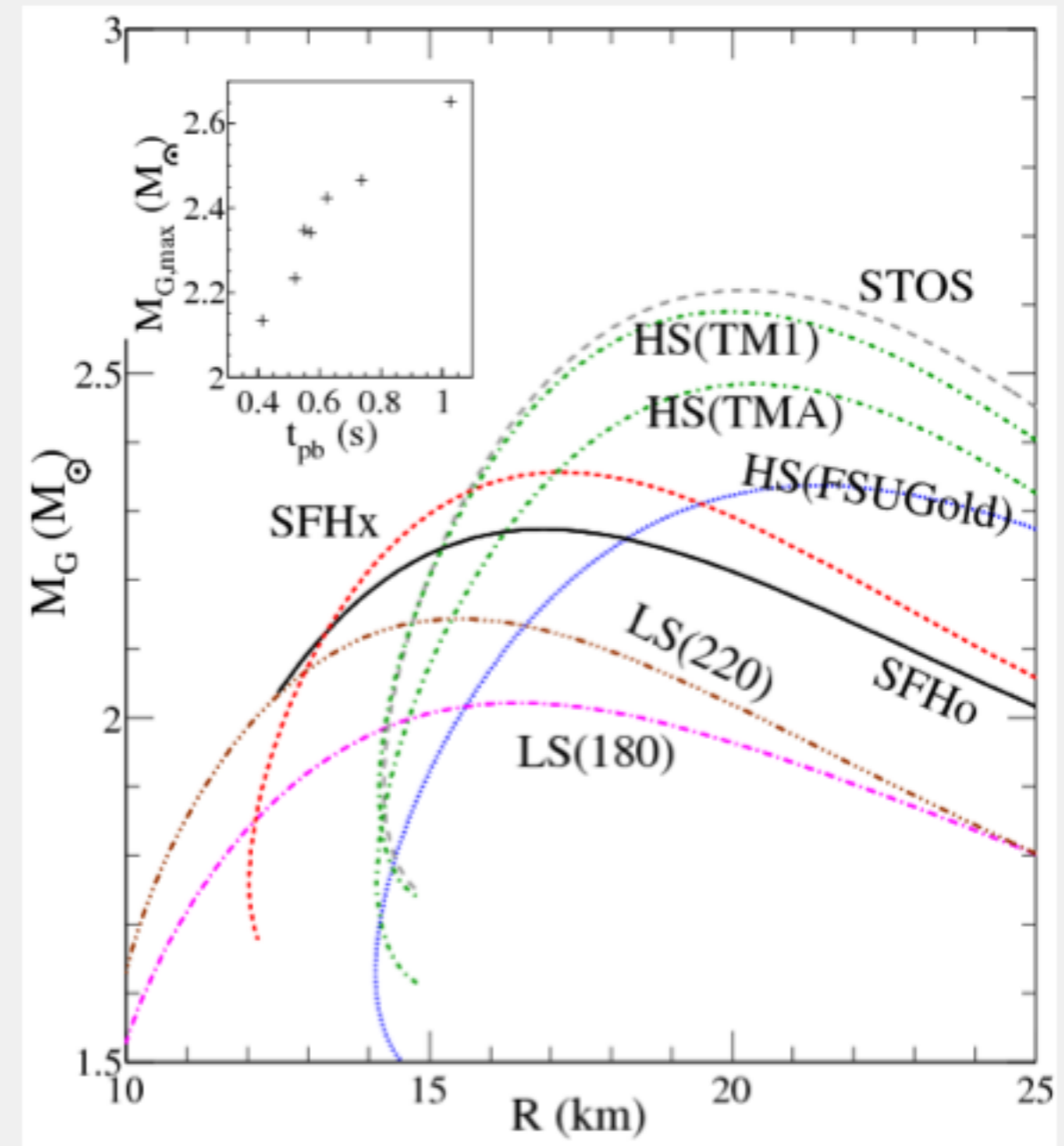
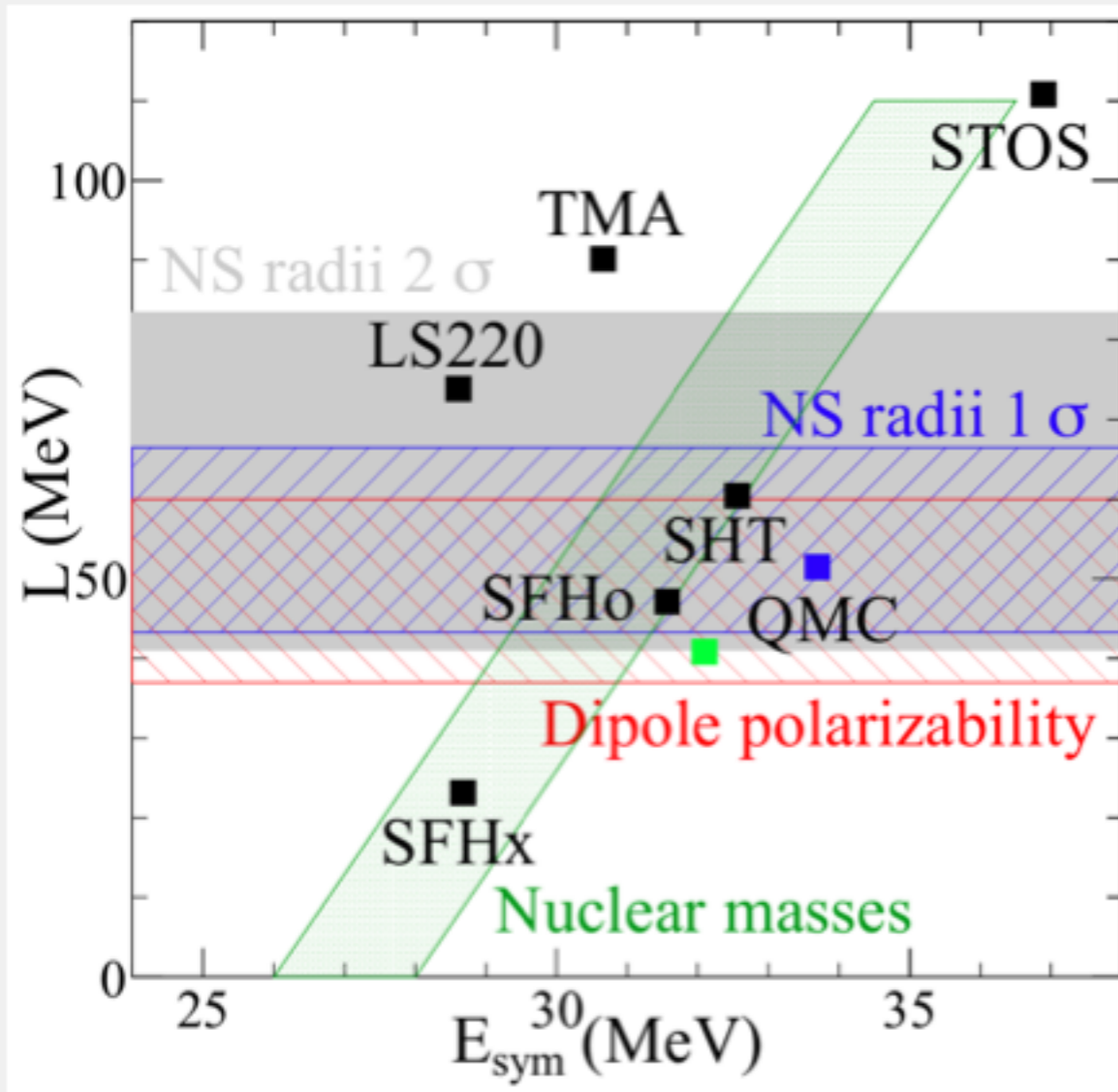
Steiner and Watts (2009), see talks by Iida and Sotani



Deibel et al. (2013)

- Seismic events excite torsional modes  $\leftarrow$  shear modulus  $\leftarrow$  composition  
composition  $\leftarrow S, L$
- Unknown mechanism for X-ray generation in magnetosphere
- May come from core modes (these may depend on R and thus L anyway)
- Entrainment is important - How does entrainment depend on L?

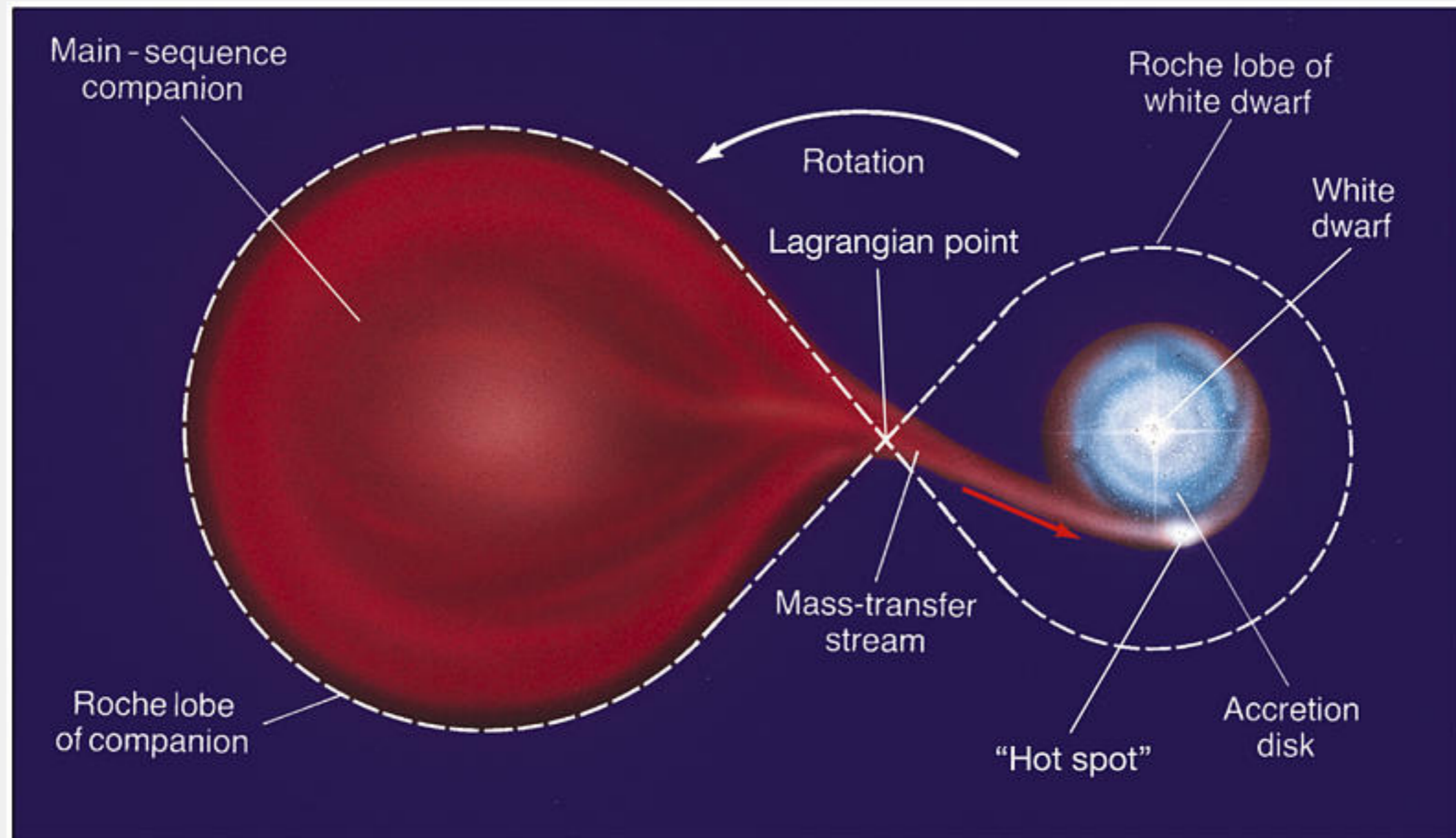
# Supernova EOS and the Symmetry Energy



Steiner, Hempel, and Fischer (2013)

- Limited number of supernova EOSs which satisfy  $S - L$  correlation
- Current EOS uncertainties too small to explain explosion
- Some correlations with the symmetry energy

# Accreting Neutron Stars: LMXBs

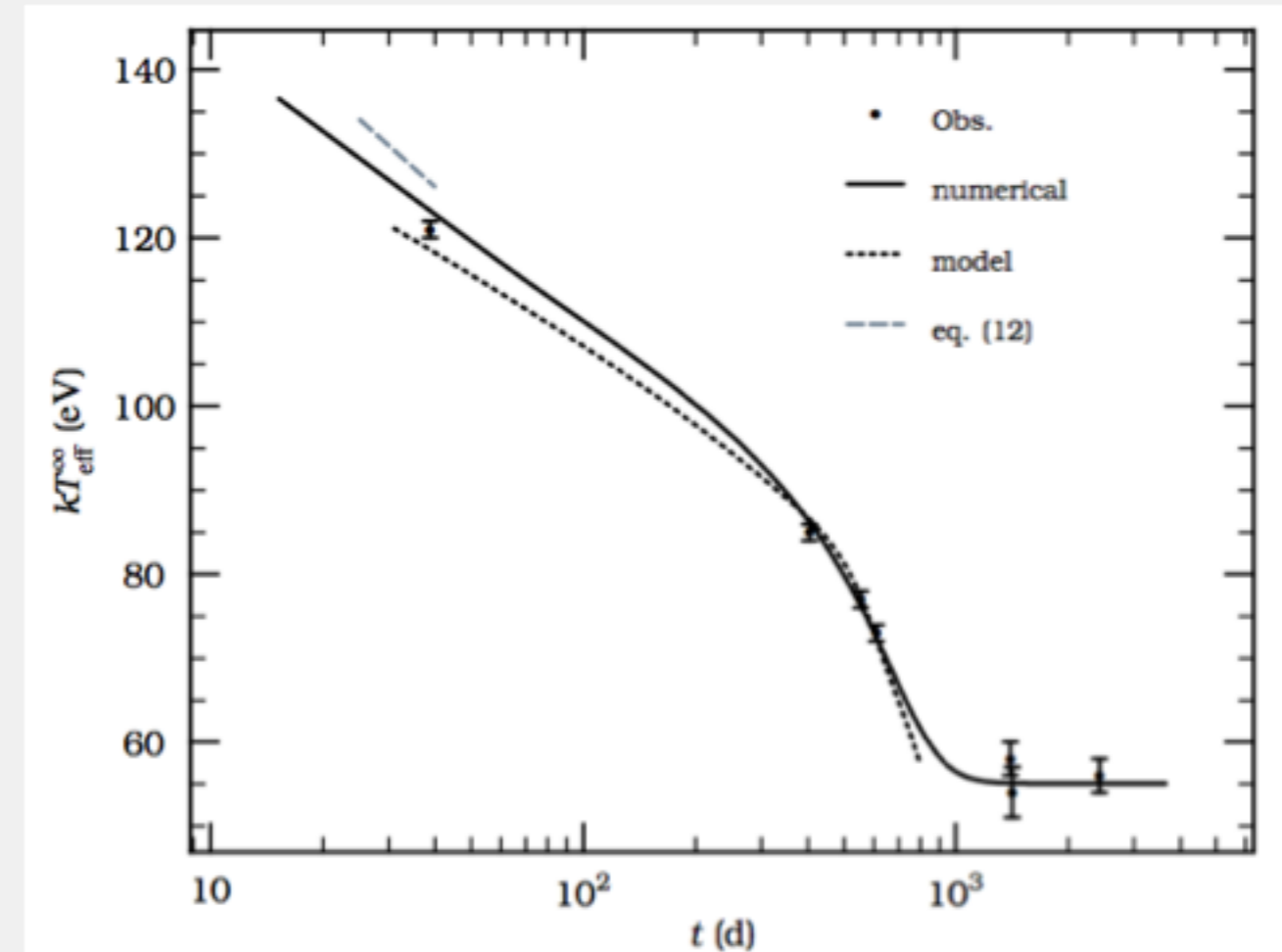


Copyright © 2005 Pearson Prentice Hall, Inc.

- Most stars have companions: neutron stars can have main-sequence companions
- Accretion heats the crust and is episodic
- At high enough density, H and He are unstable to thermonuclear explosions

# Accreted Neutron Star Crusts

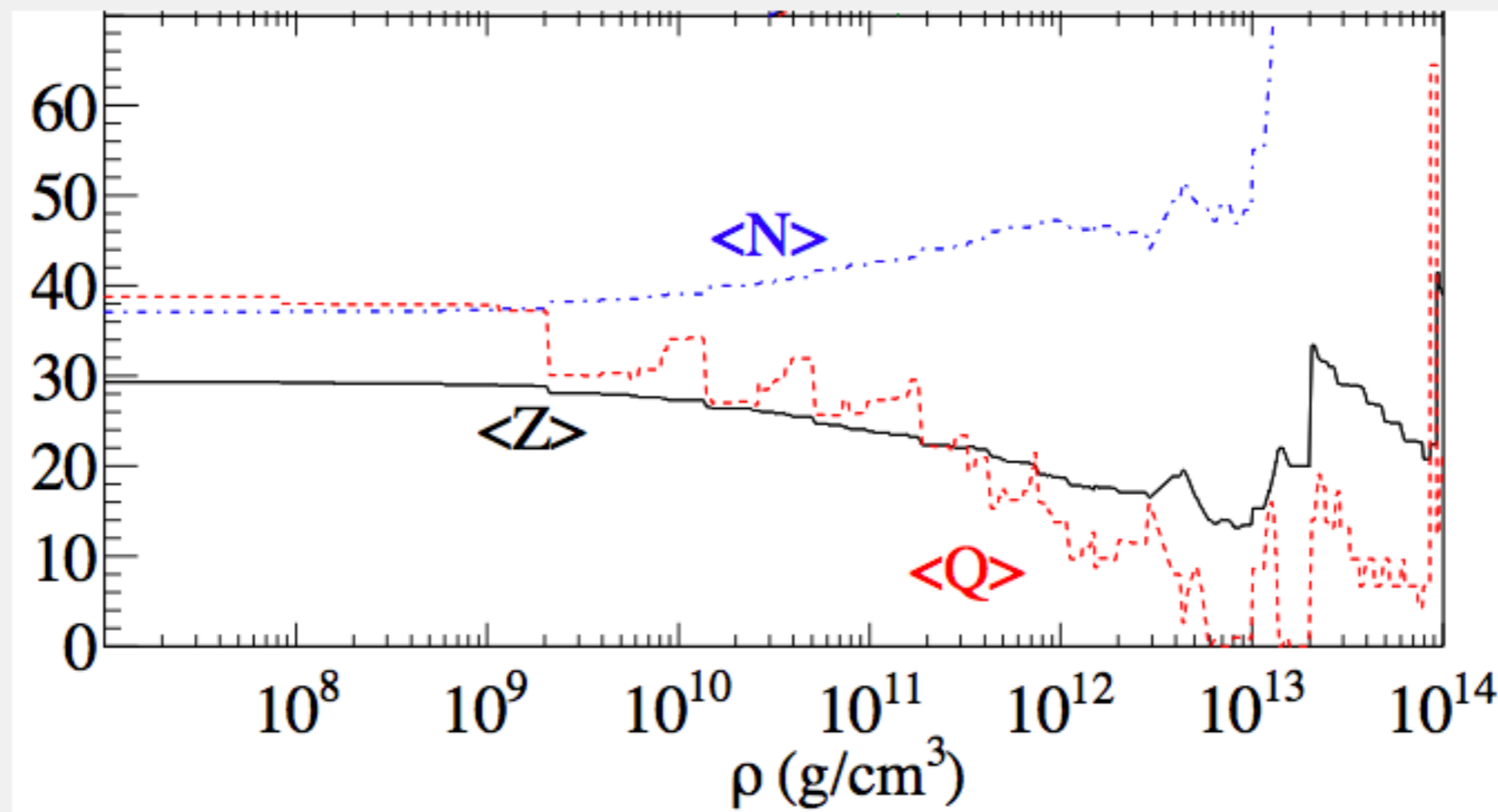
- In a cold neutron star, surface is usually taken to be  $\sim {}^{56}\text{Fe}$
- Nuclei become larger and more neutron rich with increasing density
- The surface of accreting neutron stars is H and He
- H and He is accreted and becomes unstable - X-ray burst
- X-ray burst ashes undergo electron captures, neutron emissions, and pycnonuclear fusions
- Deep crustal heating



Brown et al. 2009, Page et al. 2012

# A Multicomponent Model of the Deep Crust

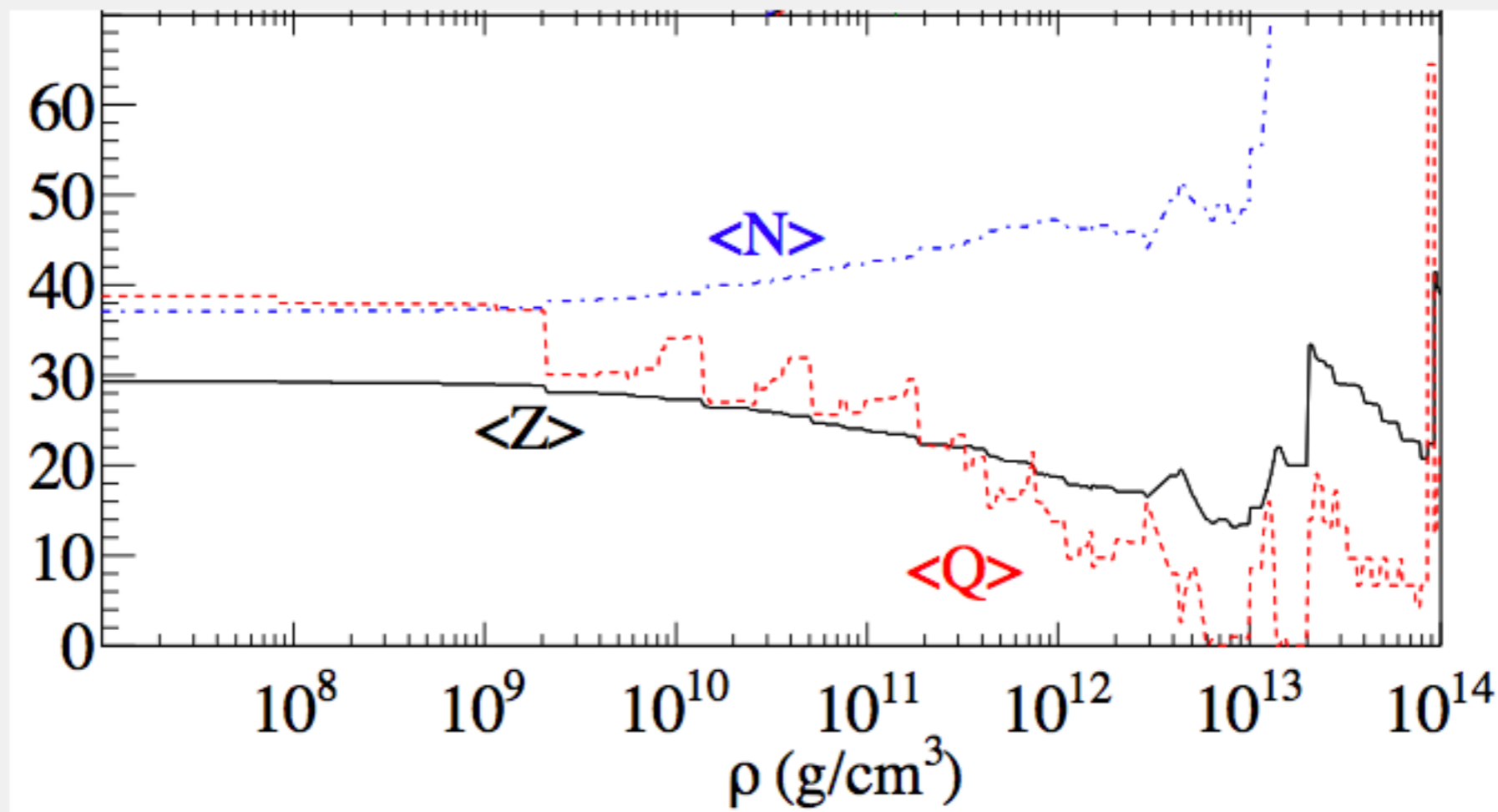
- Steiner (2012) is the first multi-component model of the accreted crust beyond neutron drip
- Employ a liquid-droplet model with an RMS deviation of 1.1 MeV
- Use quasi-statistical equilibrium instead of a full reaction network
- The multi-component model is important because it resolves reaction pathways that are impossible in single-component model



Steiner (2012)

# Symmetry Energy and Deep Crustal Heating

- Vary masses based on nucleon-nucleon interactions with different symmetry energies
- Skyrme models: SLy4 ( $L \sim 45$  MeV) and Gs ( $L \sim 90$  MeV)
- Begin with an initial composition of X-ray burst ashes
- Find that SLy4 gives 2.4 MeV per nucleon while Gs gives 4.8 MeV per nucleon



Steiner (2012)



# Other Connections to the Symmetry Energy

- Moment of inertia,  $I \sim MR^2$  (Talks by Newton and Fattoyev)
  - Potential for measurement in double pulsar J0737
  - Also  $L$  is correlated with  $I_{\text{crust}} / I$
  - Glitches: magnetic torques spin down lattice
  - Glitch happens when neutron superfluid catches up with crust
  - $I_{\text{sfluid}}$  must be large enough to explain glitches  
[Link, Epstein, and Lattimer \(1999\)](#)
  - Entrainment lowers  $I_{\text{sfluid}}$   
[Chamel](#)
- Tidal deformability,  $\lambda$ 
  - Measured by LIGO in NS-NS merger at  $\sim 400$  Hz
- $I/M^3$  correlated with  $\lambda/M^5$  ([Yagi and Yunes 2013](#))
- r-process nucleosynthesis ([Talks by Shen and Roberts](#))
  - Determines neutrino spectra in the wind
  - Determines proton fraction in ejected material in NS-NS merger
  - Drip-line may be important in fission cycling

# Summary

Symmetry energy is important for both nuclear physics and astrophysics

- There are gateway quantities that we may want to consider carefully
- Be careful with direct Urca,  $S(n)$ , and  $Q(n)$
- $L$  important for NS radii - Lots of fun next week on that
- Maybe important for flares, but this is a bit messy
- Never a bad time to use good microphysics in astrophysical simulations
- Deep crustal heating sensitive to  $L$