r-Process Powered Transients from Compact Object Mergers *r*-Process Powered Transients from Compact Object Mergers

> "I'm partial to the name 'blingnova' to describe this kind of event, since what we are seeing is basically an ostentatious glimmering of riches," Kasen said. - from The Washington Post

Outline

- Introduction
- Mass ejection during mergers
- Evolution of the ejecta
 - Nuclear Processing
 - Radioactive decay
- Optical or infrared transients
- Implications for r-process production
- Observational constraints

Multi-Messenger Events

- Gravitational Waves (LIGO, VIRGO, etc.)
- Neutrinos
- Gamma Rays (progenitors of short GRBs?)



- Optical?
- Chemical Evolution



Multi-Messenger Events

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Chemical Evolution

Lee & Ramirez-Ruiz (2007)



Multi-Messenger Events

- Gravitational Waves (LIGO, VIRGO, etc.)
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Lee & Ramirez-Ruiz (2007)





Chemical Evolution





• In-spiral driven by gravitational wave emission:

$$\tau_{mrg} = 8 \times 10^9 P_d^{8/3} M_{2.8}^{-5/3} (1+q)(1+1/q) \,\mathrm{yr}$$

• Tidal effects produce minimum stable angular momentum for polytropic binaries. Once more angular momentum is lost, binaries become dynamically unstable (Lai et al.'94).

Merger Dynamics

Rosswog & Ramirez-Ruiz (2002)

Merger Dynamics



Merger Mass Ejection

- Material is tidally ejected through the outer Lagrange points
- In GR, material is also ejected from the collision region
- Significant variation in the amount of unbound mass
- Significant variation between Newtonian and GR models



Ejected Mass Distribution



 $M_{ejected} \approx 0 - 0.1 M_{\odot}$

Homologous evolution sets in very quickly after coalescence

EoS Dependence of Mass Ejection



- Smaller radius -> larger velocity at collision -> increased mass ejection
 - Hotokezaka EoSs: APR4, ALF2, H4, and MS1 Bauswein EoSs: Finite temperature supernova EoSs

From One to Two Tails

	Type M		Primary Mass (M_{\odot})	Ejected Mass (M_{\odot})	Ejecta Velocity (c)			
	NS–NS	1.00	1.4	0.057	0.202			
	NS–NS	0.95	1.4	0.047	0.200			
	NS-NS BH NS	0.88	1.5 5.4	0.057	0.205			
		0.51	5.4	0.000	0.246			
BH-NSq=0.31		NS-NS q=0.88		NS-NS q=0.9	95 NS-N	5 NS-NSq=1.00		
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					40 A	0000000000000000000000000000000000000		

Ejecta Conditions



Nuclear Evolution of the Tails

Dynamical Timescale for the Ejected Material:

 $\tau_{ej} \approx 10 ms$

Ejected Material is neutron rich:

 $Y_{e,ej} \approx 0.05-0.2$

Low initial entropy:

 $s \lesssim 10$

Nuclear Evolution of the Tails

Dynamical Timescale for the Ejected Material:

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 $s \leq 10$ <

Initial distribution will be in NSE, clustered around doubly magic nuclei

Which implies a neutron to seed ratio:

$$\frac{N}{S} \approx \frac{\bar{Z}}{Y_e} - \bar{A} \sim 100$$

Can they make r-process nuclei? easy!

see Lattimer & Schramm '76 and Freiberghaus et al. '99

Nuclear Network

 Density trajectories for particles from SPH simulations taken

 Energy from nuclear reactions self-consistently added back to entropy of material

• Start from NSE distribution

Nuclear Network containing over 6000 isotopes



Heavy nuclei breakup by either neutron induced fission or spontaneous fission, fission fragments distributed over a large portion of (Z,N) space

Final Abundances



- Strong fission cycling
- Some dependence on initial conditions

Broadly consistent with solar system r-process abundances

Isotopic Abundances

- Reasonable agreement with halo stars
- Mostly sensitive to fission fragment distributions



Nuclear Heating Rate



- Power law heating rate (confirms results of Metzger et al. '10)
- Larger number of isotopes involved, sum of numerous individual decays
- Beta-decays and fission

Fairly insensitive to initial conditions (Y_e and entropy)

Optical Signal?

- Model tidal ejecta as decay heated homologously expanding sphere (Li & Paczynski '98)
- General properties of transients only depend on four parameters: heating rate, opacity, velocity, and mass of ejected material
- Reasonable values for these parameters predict

 $t_m \approx 1.5 \beta^{1/2} t_c$ = 0.98 days $\left(\frac{M}{0.01 \ M_{\odot}}\right)^{1/2} \left(\frac{3V}{c}\right)^{-1/2} \left(\frac{\kappa}{\kappa_e}\right)^{1/2}$



$$\times \left(\frac{f}{0.001}\right) \left(\frac{M}{0.01 \ M_{\odot}}\right)^{1/2} \left(\frac{3V}{c}\right)^{1/2} \left(\frac{\kappa}{\kappa_{e}}\right)^{-1/2}$$
$$T_{\text{eff},m} \approx 0.79 \beta^{-1/8} T_{1} = 2.5 \times 10^{4} \text{ K}$$
$$\times \left(\frac{f}{0.001}\right)^{1/4} \left(\frac{M}{0.01 \ M_{\odot}}\right)^{-1/8} \left(\frac{3V}{c}\right)^{-1/8} \left(\frac{\kappa}{\kappa_{e}}\right)^{-3/8}$$

Radiative Transfer Calculations

- Follow radiation transport in homologous ejecta using MC radiation transport code SEDONA
- Directly calculate gammaray thermalization rate
- Gray opacities appropriate to iron assumed for current work





Viewing angle effects larger than effects of mass ratio at peak

Angle effects washed out at late times

Limits from SGRB Observations



Geometry from Photometry?



The Opacity of *r*-Process Material

- Opacities of decay products not well known
- Our past calculations use a gray optical opacity
- Current calculations are for densities four or five orders of magnitude too high
- Maybe spectral signatures can more clearly distinguish tail geometries?

r-Process Mixture Opacity @ T=10 eV



C. Fryer (Private Communication)

What is the opacity?

 Number of possible transitions goes approximately as square of number of permutations of valence electrons

$$C = \Pi_i \frac{g_i!}{n_i!(g_i - n_i)!}$$

 Lanthanides have an open fshell, gives large complexity measure



Kasen et al. '13

r-process produced lanthanides expected to dominate opacity

Redder, Longer Transients



 Increase the opacity by an order of by at least 10, maybe 100

Increase timescale and decrease opacity

Merger Rates

TABLE II: Compact binary coalescence	e rates per M	filky Way	Equivalent (Galaxy per Myr
Source	R_{low}	$R_{\rm re}$	R_{high}	R_{max}
NS-NS (MWEG ⁻¹ Myr ⁻¹)	1 [1] ^a	100 [1] ^b	1000 [1] ^c	4000 [16] ^d
NS-BH ($MWEG^{-1} Myr^{-1}$)	$0.05 [18]^e$	$3 [18]^{f}$	$100 [18]^{g}$	
BH-BH (MWEG ^{-1} Myr ^{-1})	$0.01 [14]^{h}$	$0.4 [14]^i$	$30 [14]^{9}$	
IMRI into IMBH (GC ⁻¹ Gyr ⁻¹)			$3 [19]^k$	$20 [19]^l$
IMBH-IMBH $(GC^{-1} Gyr^{-1})$			0.007 [20]	$n 0.07 [20]^n$

Merger rates from both population synthesis and extrapolation from known NS-NS binary population are very uncertain

Predicted Merger Rates (from Abadie et al. '11)

Table 1. Properties of the observed pulsars in DNS binaries. The table contains names, spin period, spin period derivative, orbital period, mass of observed neutron star, mass of the companion, eccentricity of the orbit and time to merger. All given digits are significant. Errors, where given, are 1σ errors. References: 1 – Stairs (2004), 2 – Jacoby et al. (2006), 3 – Kasian (2008), 4 – Weisberg & Taylor (2005), 5 – Faulkner et al. (2005) and 6 – Janssen et al. (2008).

6 known NS-NS binaries will merge within a Hubble time

Name	P (ms)	$\dot{P} (ss^{-1}/10^{-18})$	$P_{\alpha b}$ (h)	Mobs (M _O)	$M_{\rm cmp}({\rm M}_{\bigodot})$	e	rmg (Gyr)	Reference
J0737-3039A	22.70	1.74	2.454	1.337+0.005	$1.250^{+0.005}_{-0.005}$	0.088	0.085	1
J0737-3039B	2773	8.8×10^2	2.454	$1.250^{+0.005}_{-0.005}$	1.337+0.005	0.088	0.085	1
B2127+11C	30.53	4.99	8.05	1.358+0.01	$1.34^{+0.01}_{-0.01}$	0.681	0.2	2
J1906+0746	144.07	2.028×10^{4}	3.98	$1.248^{+0.018}_{-0.018}$	$1.365^{+0.018}_{-0.018}$	0.085	0.3	3
B1913+16	59.03	8.63	7.752	$1.4414^{+0.0002}_{-0.0002}$	$1.3867^{+0.0002}_{-0.0002}$	0.617	0.3	4
J1756-2251	28.46	1.02	7.67	$1.312^{+0.017}_{-0.017}$	$1.258^{+0.018}_{-0.017}$	0.181	1.7	5
B1534+12	37.90	2.43	10.098	$1.3332^{+0.001}_{-0.001}$	$1.3452^{+0.001}_{-0.001}$	0.274	2.7	1
J1811-1736	104.182	0.91	451.20	$1.62^{+0.22}_{-0.55}$	$1.11^{+0.53}_{-0.15}$	0.828	>10	1
J1518+4904	40.935	0.027	207.216	$0.72^{+0.51}_{-0.58}$	$2.00^{+0.58}_{-0.51}$	0.249	>10	6
J1829+2456	41.0098	0.05	28.0	$1.14^{+0.28}_{-0.48}$	$1.36^{+0.50}_{-0.17}$	0.139	>10	1

Known pulsars in neutron star binaries (from Oslowski et al. '11)

Chemical Evolution Signal?

 $M_{r,MW} \sim 10^4 M_{\odot}$ $r_{NS-NS} \sim 10^{-4} \text{yr}^{-1}$ $M_{eject} \sim 10^{-2} M_{\odot}$

$$\rightarrow M_{r,NS-NS} \sim 10^4 M_{\odot}$$

but...

$$t_{coalesce} \approx 10^{6-8} yr$$



from Argast et al. 2004

SGRB 130603B

- SGRB detected at z=0.356 by the Swift BAT
- Early optical detection of afterglow
- Followed up ~9 days afterward with HST
- Point source seen at the position of the GRB



Berger et al. '13

SGRB 130603B



 Late time emission consistent with standard afterglow, different power laws between different studies

Also consistent with kilonova with M~0.01 Msun and v~0.1c

SGRB 130603B

More to come:

A search for kilonova emission associated with GRB 130603B: the smoking gun signature of a compact binary merger event

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We note that we felt compelled to submit this provisional report of our work, despite our HST DDT program being incomplete, due to other authors having already posted an analysis of the publicly available first epoch data.

- Better background subtraction
- Is the point source really transient?

Conclusions

- Significant amount of neutron rich mass is ejected during NS-NS mergers
- Amount of mass ejected depends on cold NS mass radius relation
- Decay of radioactive isotopes produced in tails can produce observable optical transient, makes possible *in situ* observation of the production of the *r*-process
- Opacities of ejected material biggest question

Potentially observed one in association with SGRB 130603B