

IAS and Skin Constraints on the Symmetry Energy

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Symmetry Energy in Nuclear Mass Formula

Textbook Bethe-Weizsäcker formula:

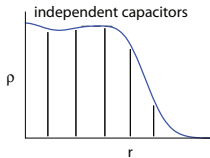
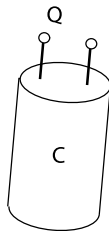
$$E = -a_V A + a_S A^{2/3} + a_C \frac{Z^2}{A^{1/3}} + a_a \frac{(N-Z)^2}{A} + E_{\text{mic}}$$

Symmetry energy: charge $n \leftrightarrow p$ symmetry of interactions

Analogy with capacitor:

$$E_a = a_a \frac{(N-Z)^2}{A} \equiv \frac{(N-Z)^2}{\frac{A}{a_a}} \Leftrightarrow E = \frac{Q^2}{2C}$$

?Volume Capacitance? $E_a = \frac{(N-Z)^2}{\frac{A}{a_a}} \rightarrow \frac{(N-Z)^2}{\frac{A}{a_a^V} + \frac{A^{2/3}}{a_a^S}}$



Thomas-Fermi (local density) approximation:

$$'C' \equiv \frac{A}{a_a(A)} = \int \frac{\rho dr}{S(\rho)} = \frac{A}{a_a^V}, \text{ for } S(\rho) \equiv a_a^V$$

TF breaks in nuclear surface at $\rho < \rho_0/4$

PD&Lee NPA818(2009)36



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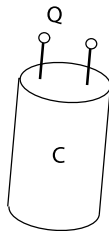
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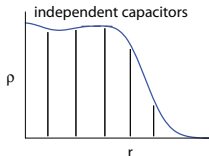
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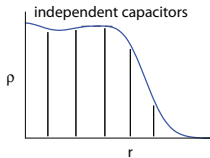
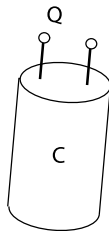
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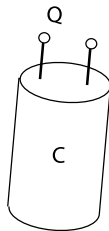
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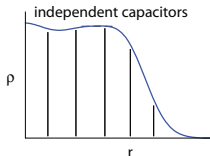
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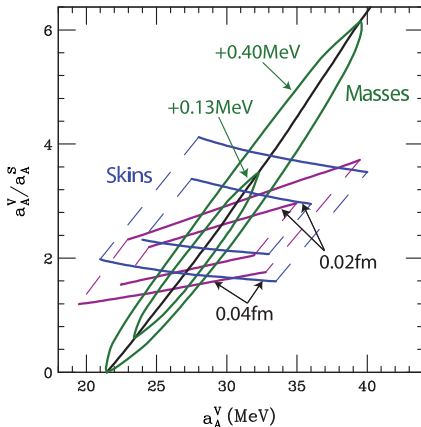
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Mass and Skin Fits



Symmetry Energy:

$$E_a = \frac{a_a^V}{A} \frac{(N-Z)^2}{1 + \frac{a_a^V}{a_a^S A^{1/3}}}$$

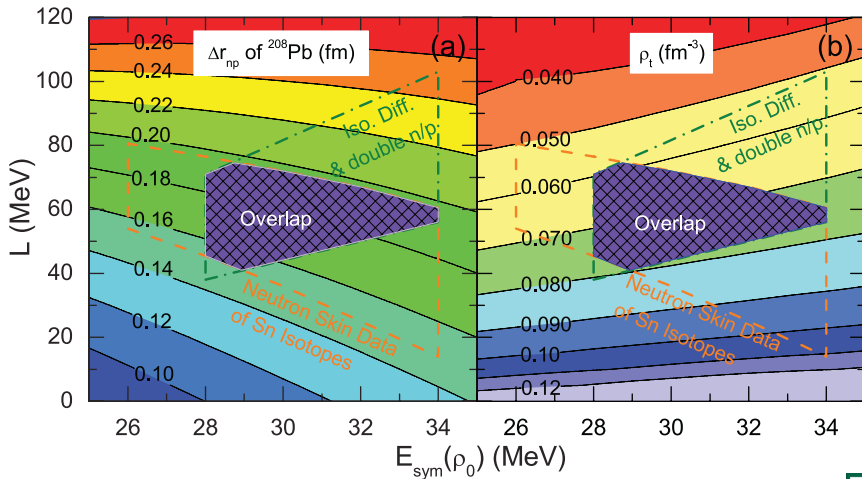
Skin:

$$\Delta r_{np} = \frac{2}{3} \frac{r_{rms}}{A^{1/3}} \frac{a_a}{a_a^S} \left(\frac{N-Z}{A} - Coul \right)$$

PD NPA723(2003)233

$$a_a^S \leftrightarrow L$$



Fits in $L-a_a^V$ PlaneLie-Wen Chen *et al* PRC82(10)024321

Charge Invariance

? $a_a(A)$? Conclusions on sym-energy details, following E -formula fits, interrelated with conclusions on other terms in the formula: asymmetry-dependent Coulomb, Wigner & pairing + asymmetry-independent, due to $(N - Z)/A$ - A correlations along stability line [PD NPA727(03)233]!

Best would be to study the symmetry energy in isolation from the rest of E -formula! Absurd?!

Charge invariance to rescue: lowest nuclear states characterized by different isospin values (T, T_z), $T_z = (Z - N)/2$. Nuclear energy scalar in isospin space:

sym energy
$$E_a = a_a(A) \frac{(N - Z)^2}{A} = 4 a_a(A) \frac{T_z^2}{A}$$

$$\rightarrow E_a = 4 a_a(A) \frac{T^2}{A} = 4 a_a(A) \frac{T(T + 1)}{A}$$



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$a_a(A)$ Nucleus-by-Nucleus

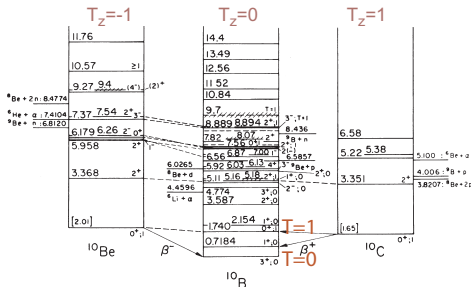
$$\rightarrow E_a = 4 a_a(A) \frac{T(T+1)}{A}$$

In the ground state T takes on the lowest possible value

$T = |T_z| = |N - Z|/2$. Through '+1' most of the Wigner term absorbed.

Formula generalized to the lowest state of a given T (e.g. Jänecke *et al.*, NPA728(03)23).

?Lowest state of a given T : isobaric analogue state (IAS) of some neighboring nucleus ground-state.



Study of changes in the symmetry term possible nucleus by nucleus

PD&Lee arXiv:1307.4130



Queries in the Context of Data

Are expansions valid? Coefficient values??

$$E_{\text{IAS}}^* = E_{\text{IAS}} - E_{\text{gs}} \stackrel{?}{=} \frac{4 a_a(A)}{A} \Delta [T(T+1)] + \Delta E_{\text{mic}}$$

Is the excitation energy linear in the isospin squared??

$$\frac{A}{a_a(A)} \stackrel{?}{=} \frac{A}{a_a^V} + \frac{A^{2/3}}{a_a^S}$$

or

$$a_a^{-1} \stackrel{?}{=} (a_a^V)^{-1} + (a_a^S)^{-1} A^{-1/3}$$

Is the volume-surface separation valid?

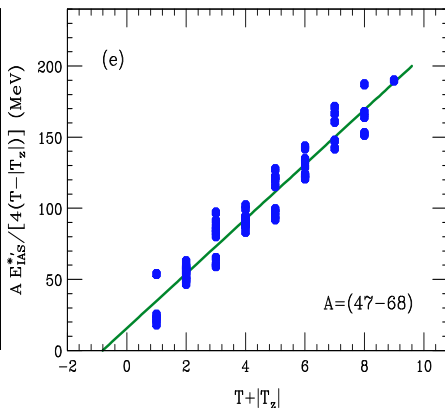
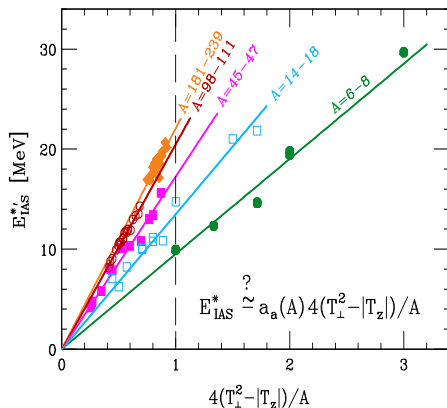
⇒ From an a_a^V - a_a^S fit can one learn about a_a^V and L for uniform matter?



Insight into IAS Analysis

IAS data: Antony *et al.* ADNDT66(97)1

Shell corrections: Koura *et al.* ProTheoPhys113(05)305



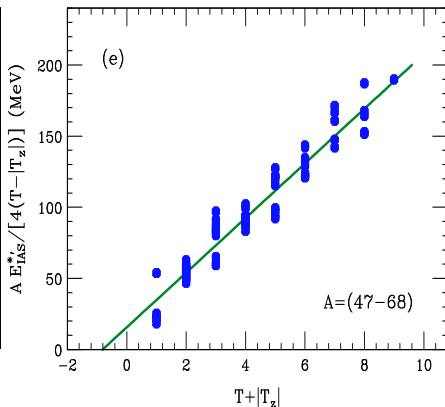
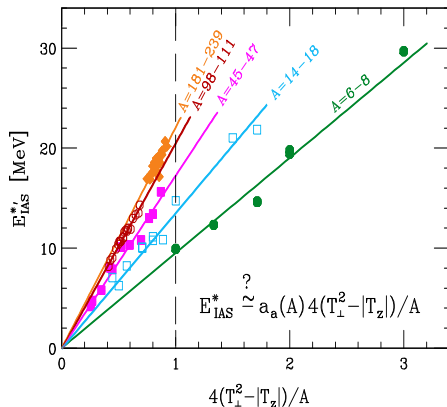
Excitation energies to IAS, E_{IAS}^* , for different A , lumped together in narrow regions of A . Is $E_a \propto T(T + 1)$? YES



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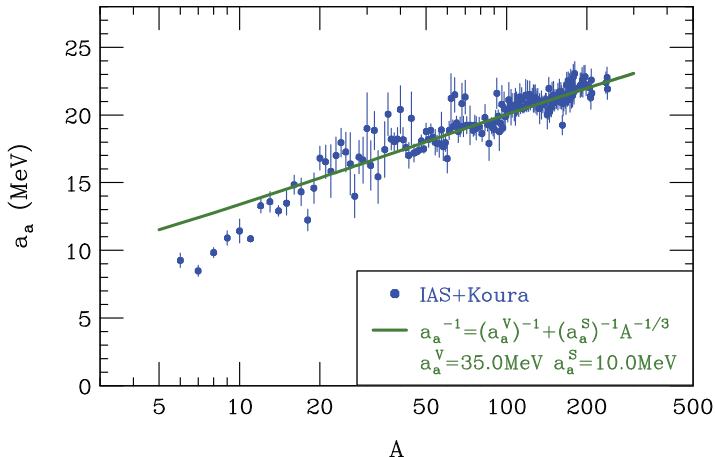
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$a_a(A)$ with Shell Corrections

$$a_a(A) = \frac{A}{4} \frac{E_{\text{IAS}}^* - \Delta E_{\text{mic}}}{\Delta T^2}$$

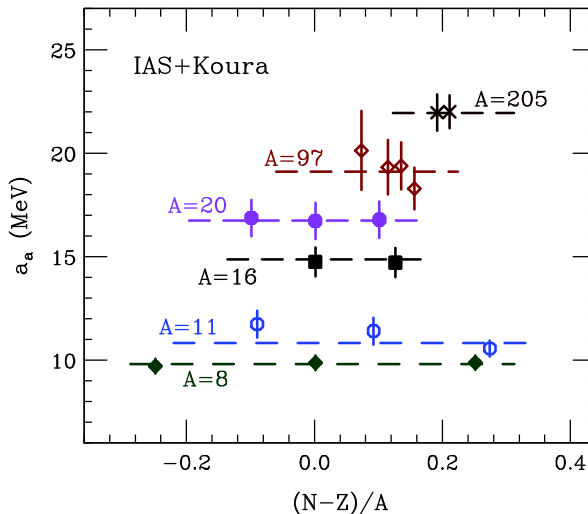
E_{mic} from: Koura *et al.*
ProTheoPhys113(05)305



Heavy nuclei $a_a \sim 22 \text{ MeV}$, light $a_a \sim 10 \text{ MeV}$ (more surface + low ρ)



Z-Dependence of Symmetry Coefficients?



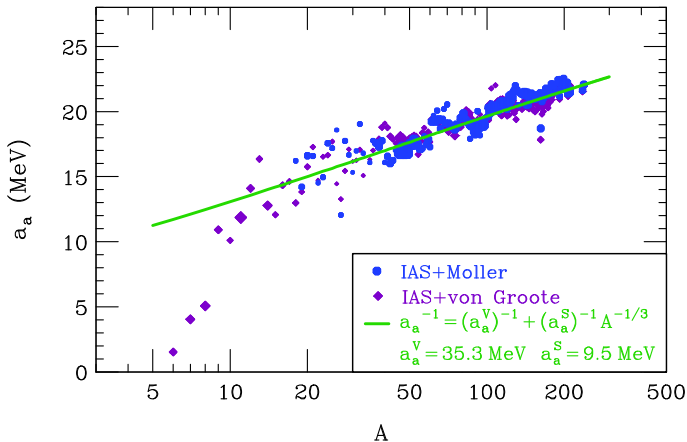
surface
less dominant,
higher
average ρ

surface
more dominant,
lower
average ρ

Symmetry coefficients on a nucleus-by-nucleus basis



Sensitivity to Shell Corrections



Fit to raw data ($A > 30$) in the middle, but:

Moller *et al.* fit: $a_a^V = 39.73 \text{ MeV}$, $a_a^S = 8.48 \text{ MeV}$

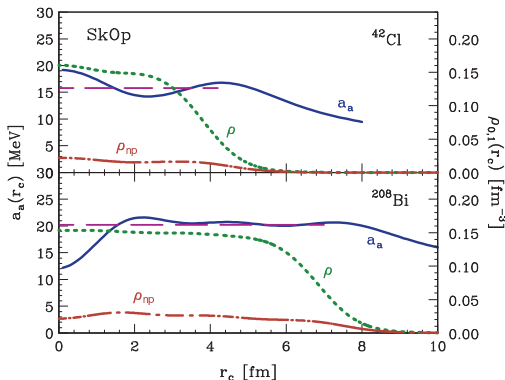
von Groote *et al.*: $a_a^V = 31.74 \text{ MeV}$, $a_a^S = 11.27 \text{ MeV}$



Comparisons to Skyrme-Hartree-Fock

Issues in data-theory comparisons (codes by P.-G. Reinhard):

1. No isospin invariance in SHF - impossible to follow the procedure for data
2. Shell corrections not feasible at such scrutiny as for data
3. Coulomb effects.

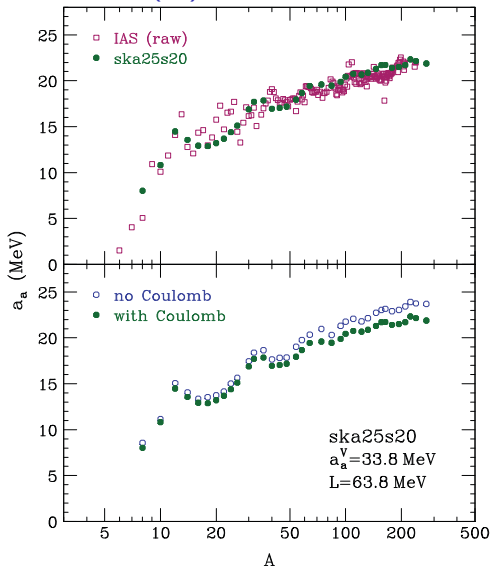


Solution: Procedure that yields the same results as the energy, in the bulk limit, but is weakly affected by shell effects:

$$\begin{aligned} \frac{(N-Z)_{r < r_c}}{N-Z} &= \frac{C_{r < r_c}}{C} \\ &= \frac{a_a}{A a_a^V} \int_{r < r_c} \frac{\rho}{S(\rho)} \end{aligned}$$



$a_a(A)$ from Mean-Field Calculations

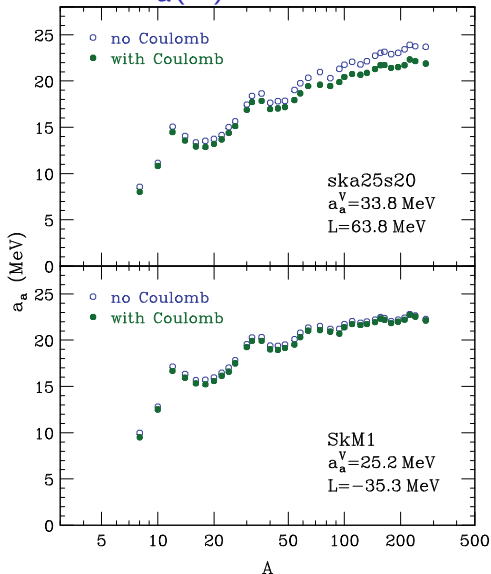


Skyrme-Hartree-Fock theory
(codes by
P.-G. Reinhard)

Similar behavior
with A as for IAS



$a_a(A)$ from Different Mean Fields



?Slope L in ρ
 \Leftrightarrow slope in A ??

Less impact of
the slope L at ρ_0
than expected!

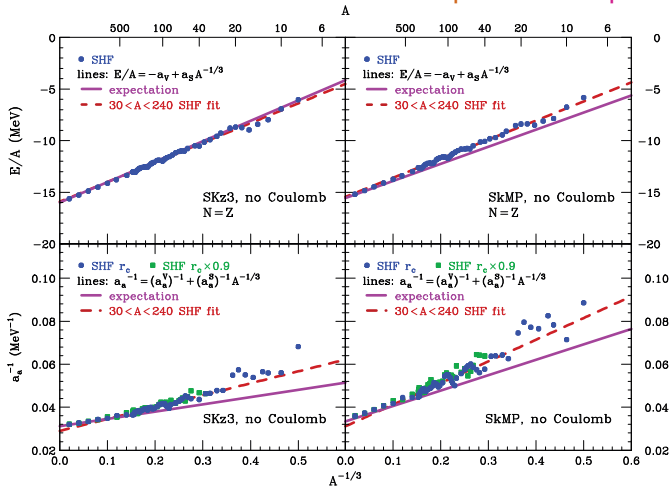
??Difficulty for
 L determination??



Model-Independent Large-A Expansion??

Symbols: results of spherical no-Coulomb SHF calcs

⇒ Lines: volume-surface decomposition - expectation vs fit



→ Symmetric
matter
energy
f/sample
Skyrmes

~ Works

→ Symmetry
coefficient

~ Not...



Expectations from half- ∞ matter.

Can $S(\rho)$ Be Constrained??!

Pearson
correlation
coefficient

$$r_{XY} = \frac{\langle (X - \langle X \rangle)(Y - \langle Y \rangle) \rangle}{\sqrt{\langle (X - \langle X \rangle)^2 \rangle \langle (Y - \langle Y \rangle)^2 \rangle}}$$

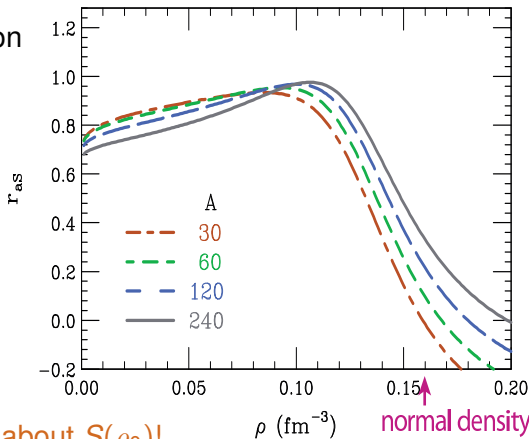
$|r| \sim 1$ - strong correlation

$r \sim 0$ - no correlation

$X \equiv a_a(A)$

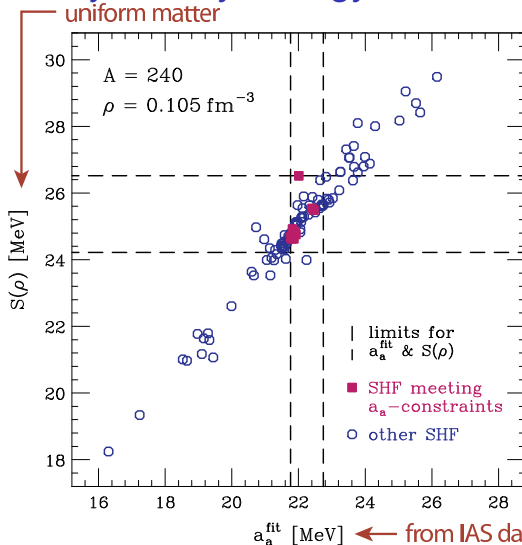
$Y \equiv S(\rho)$

Ensemble of Skyrmes



Nearly no information about $S(\rho_0)$!

Symmetry-Energy Correlations When Strong



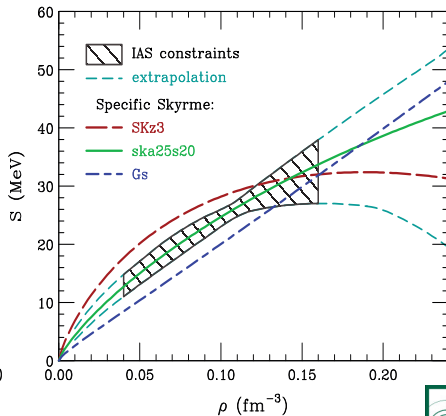
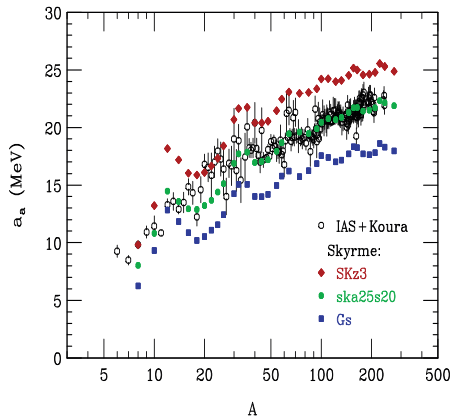
NO $S(\rho) \approx a_a$!

Vehicle: phenomenological mean-field theory



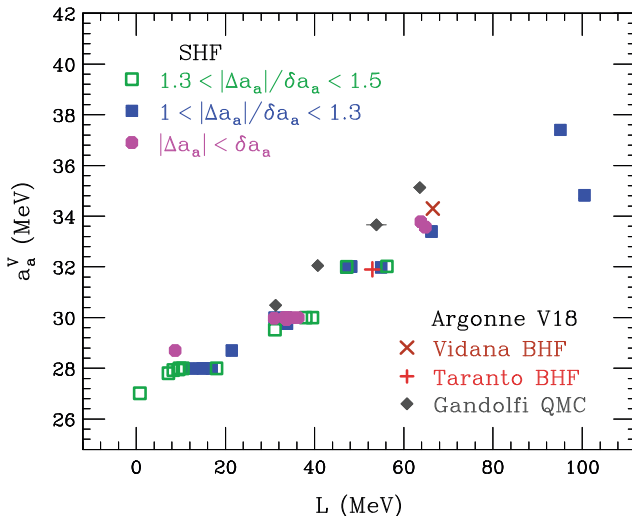
Constraints on Symmetry Energy $S(\rho)$

Demand that Skyrme approximates IAS results at $A > 30$ produces a constraint area for $S(\rho)$:



?Slope constraint??



Correlation at ρ_0 

?Slope constraint??



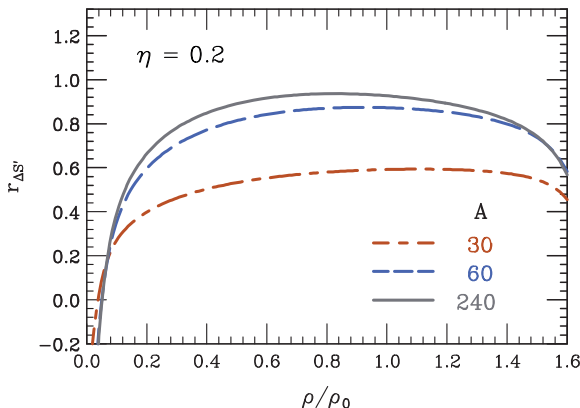
Asymmetry Skin & Energy Stiffness

Pearson coef of
 $\Delta r_{np} = r_n^{\text{rms}} - r_p^{\text{rms}}$
 & stiffness of S

$$\gamma(\rho) := \frac{\rho}{S} \frac{dS}{d\rho}$$

f/different A
 at fixed

$$\eta = (N - Z)/A$$

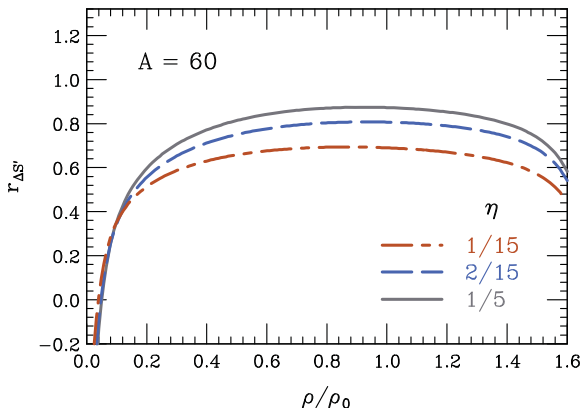


Asymmetry Skin & Energy Stiffness

Pearson coef
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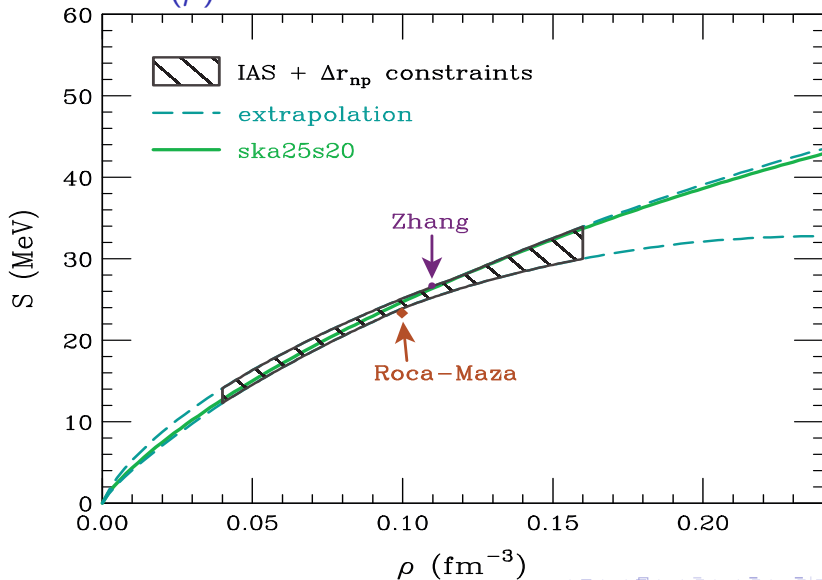
Asymmetry Skins from Measurements

Nucleus	Reference	Data Source	Δr_{np} [fm]	Δr_{np}^{GF} [fm]
^{48}Ca	Friedman [92]	pionic atoms	0.13 ± 0.06	
	Gils et al. [93]	elastic α scattering	0.175 ± 0.050	
	Ray [94]	elastic \vec{p} scattering	0.229 ± 0.050	
	Clark et al. [95]	elastic p scattering	0.103 ± 0.040	
	Shlomo et al. [96]	elastic p scattering	0.10 ± 0.03	
	Gibbs et al. [97]	elastic π scattering	0.11 ± 0.04	
		combined results	$0.129 \pm 0.053^{\boxtimes}$	0.215 ± 0.012

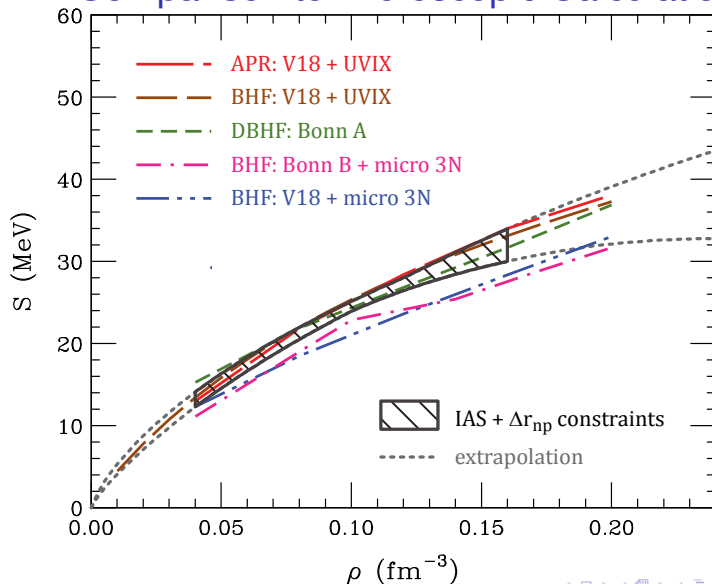


^{207}Pb	Starodubsky et al. [99]	elastic p scattering	0.186 ± 0.041	0.175 ± 0.023
^{208}Pb	Starodubsky et al. [99]	elastic p scattering	0.197 ± 0.042	
	Ray [94]	elastic \vec{p} scattering	0.16 ± 0.05	
	Clark et al. [95]	elastic p scattering	0.119 ± 0.045	
	Zenihiro et al. [98]	elastic p scattering	0.211 ± 0.063	
	Friedman [92]	elastic π^+ scattering	0.11 ± 0.06	
	Friedman [92]	pionic atoms	0.15 ± 0.08	
		combined results	$0.159 \pm 0.041^{\boxtimes}$	0.179 ± 0.023

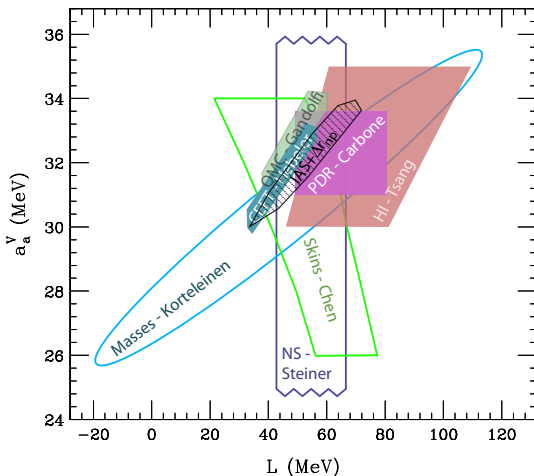
$S(\rho)$ from Combined Constraints



Comparison to Microscopic Calculations



Symmetry Energy at ρ_0 ?



$$S(\rho) = a_a^V + \frac{L}{3} \frac{\rho - \rho_0}{\rho + 0} + \dots$$



Conclusions

- Symmetry-energy term weakens as nuclear mass number decreases: from $a_a \sim 23$ MeV to $a_a \sim 9$ MeV for $A \lesssim 8$.
- For $A \gtrsim 25$, $a_a(A)$ may be fitted with $a_a^{-1} = (a_a^V)^{-1} + (a_a^S)^{-1} A^{-1/3}$, where $a_a^V \approx 35$ MeV and $a_a^S \approx 10$ MeV.
- Weakening of the symmetry term can be tied to the weakening of $S(\rho)$ in uniform matter, with the fall of ρ .
- Including skin sizes, significant, $\lesssim \pm 1.0$ MeV, constraints on $S(\rho)$ at densities $\rho = (0.04-0.13) \text{ fm}^{-3}$.
- Around ρ_0 : *strongly correlated* $a_a^V = (30.2-33.7) \text{ MeV}$ and $L = (35-70) \text{ MeV}$.

To do: Dedicated Skyrme interactions.

PD&Lee arXiv:1307.4130

PHY-1068571



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PD&Lee arXiv:1307.4130

PHY-1068571



Conclusions

- Symmetry-energy term weakens as nuclear mass number decreases: from $a_a \sim 23$ MeV to $a_a \sim 9$ MeV for $A \lesssim 8$.
- For $A \gtrsim 25$, $a_a(A)$ may be fitted with $a_a^{-1} = (a_a^V)^{-1} + (a_a^S)^{-1} A^{-1/3}$, where $a_a^V \approx 35$ MeV and $a_a^S \approx 10$ MeV.
- Weakening of the symmetry term can be tied to the weakening of $S(\rho)$ in uniform matter, with the fall of ρ .
- Including skin sizes, significant, $\lesssim \pm 1.0$ MeV, constraints on $S(\rho)$ at densities $\rho = (0.04-0.13) \text{ fm}^{-3}$.
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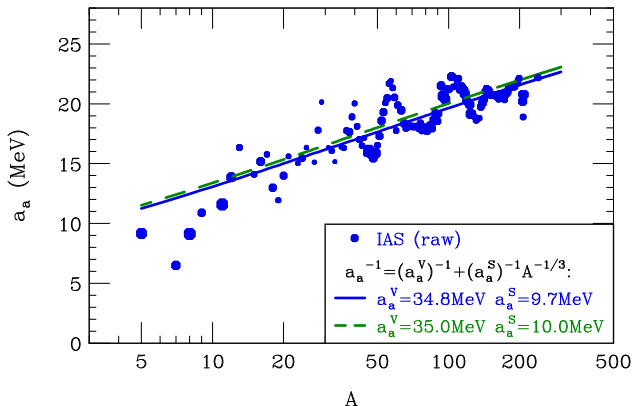
PHY-1068571



$a_a(A)$ without Shell Corrections

$$a_a(A) = \frac{A}{4} \frac{E_{\text{IAS}}^*}{\Delta T^2}$$

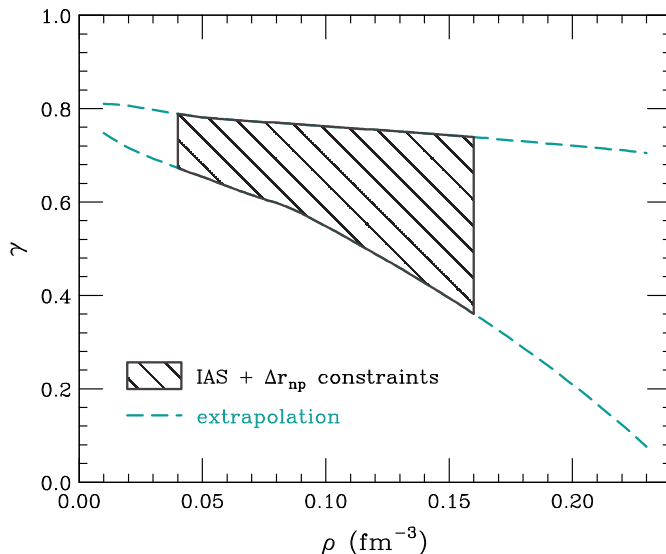
IAS data: Antony *et al.*
ADNDT66(97)1



Lines: fits to $a_a(A)$ assuming *volume-surface competition*
analogous to that for E_1 . ??Fundamental knowledge??



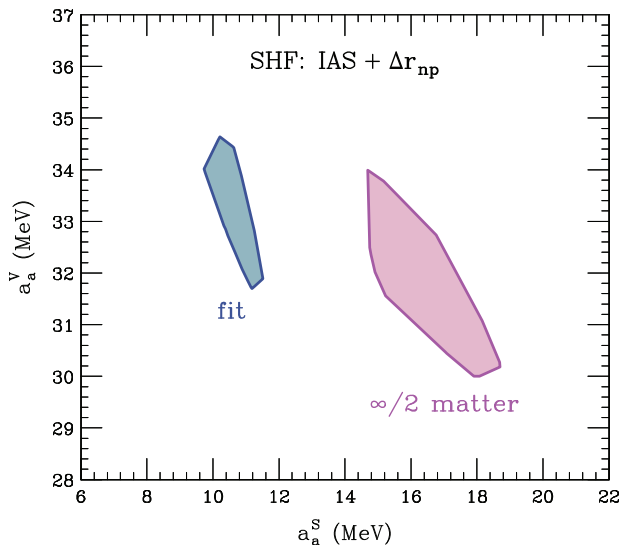
Stiffness of the Symmetry Energy



$$S \propto \rho^\gamma$$



Robustness of Macroscopic Description?



Constraints at ρ_0

