### Tidal interactions during neutron star mergers: symmetry energy considerations

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NuSYM13; July 23



### Outline

- Motivation: Symmetry energy at saturation density and higher
- NS-NS tidal interactions
- Equilibrium tides
  - Tidal polarizability/Love number
  - Preparation of Nuclear Matter models
  - Results
- Dynamical Tides I: resonant excitation of g-modes
- Dynamical Tides II: resonant excitation of crust modes; crust shattering
- Summary

#### Symmetry energy



$$E(n,\delta) = E_0(n) + S(n)\delta^2 + \dots \qquad \delta = 1 - 2x$$
  
$$S(n) = J + L\chi + \frac{K_{\text{sym}}}{2}\chi^2 + \dots \qquad \chi = \frac{n - n_0}{3n_0}$$

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Other notations are available

#### **Experimental Constraints: Saturation Density**



- 1. Chen,Ko,Li; PRL94
- 2. Famiano et al; PRL97
- 3. Shetty et al; PRC76
- 4. Klimkiewicz et al; PRC76
- 5. Danielewicz, Lee; NPhys A818
- 6. Tsang et al; PRL102
- 7. Centelles et al; PRL102
- 8. Warda et al; PRC80
- 9. Carbone et al; PRC81
- 10.Chen, Ko, Li, Xu; PRC82
- 11.Zenihiro et al; PRC82
- 12.Xu, Li, Chen; PRC82
- 13.Liu et al; PRC82
- 14.Chen; PRC83
- 15.Möller et al; PRL108
- 16.Lattimer, Lim; arxiv:1203.4286
- 17.Abrahamyan et al, PRL108
- 18.Dong et al; PRC85
- 19.Piekarewicz et al; PRC85
- 20.Zhang, Chen; arxiv:1302.5327
- 21.Roca-Maza et al. PRC87
- 22.Wang, Ou, Liu, PRC87
- 23.Li et al. PLB721
- 24. Agrawal et al, arxiv:1305:5336

#### **Experimental Constraints: Saturation Density**



# Not full story: Saturation constraints on J,L not necessarily constraining at high densities



• Can we find astrophysical observables that are sensitive to high density EOS/symmetry energy? (e.g. neutrino signal from PNS - Roberts et al PRL108 (2012))

**NS-NS** Mergers



#### **Equilibrium tides**



#### Tidal polarizability and Love number

$$Q_{ij} = -k_2 \frac{2R^5}{3G} E_{ij} \equiv -\lambda E_{ij}$$

$$k_{2} = \frac{1}{20} \left(\frac{R_{s}}{R}\right)^{5} \left(1 - \frac{R_{s}}{R}\right)^{2} \left[2 - y_{R} + (y_{R} - 1)\frac{R_{s}}{R}\right] \times \qquad r\frac{dy(r)}{dr} + y(r)^{2} + y(r)F(r) + r^{2}Q(r) = 0 ,$$

$$\times \left\{\frac{R_{s}}{R} \left(6 - 3y_{R} + \frac{3R_{s}}{2R} \left(5y_{R} - 8\right) + \frac{1}{4} \left(\frac{R_{s}}{R}\right)^{2} \left[26 - F(r) - \frac{r - 4\pi r^{3} \left(\mathcal{E}(r) - P(r)\right)}{r - 2M(r)} \right] ,$$

$$- 22y_{R} + \left(\frac{R_{s}}{R}\right) \left(3y_{R} - 2\right) + \left(\frac{R_{s}}{R}\right)^{2} \left(1 + y_{R}\right) \right] \right) +$$

$$+ 3\left(1 - \frac{R_{s}}{R}\right)^{2} \left[2 - y_{R} + (y_{R} - 1)\frac{R_{s}}{R}\right] \times \qquad Q(r) = \frac{4\pi r \left(5\mathcal{E}(r) + 9P(r) + \frac{\mathcal{E}(r) + P(r)}{\partial P(r) / \partial \mathcal{E}(r)} - \frac{6}{4\pi r^{2}}\right)}{r - 2M(r)}$$

$$\times \log\left(1 - \frac{R_{s}}{R}\right)^{-1}, \qquad (1) \qquad - 4\left[\frac{M(r) + 4\pi r^{3}P(r)}{r^{2} \left(1 - 2M(r)/r\right)}\right]^{2}.$$

- Tidal polarizability  $\lambda$  function of global NS properties (M,R) and internal structure (function  $y_R$  obtained by integration out from center of star)
- Clearly discriminates between NSs and Strange stars

- Postnikov, Prakash, Lattimer, PRD82, 024012 (2010)

• Sensitivity to high density symmetry energy?



- Famous correlation between fiducial pressure P and NS radius R
- Scatter of order 1km due largely to differences super saturation symmetry energy behavior
- Differences from different parameterizations within same EDF, different EDFs; would like to disentangle

• Skyrme-Hartree-Fock (SHF) model of nuclear matter:

$$\begin{aligned} \mathscr{H} &= \frac{\hbar^2}{2M} \tau + t_0 \left[ (2 + x_0) \rho^2 - (2x_0 + 1) \left( \rho_n^2 + \rho_p^2 \right) \right] / 4 \\ &+ t_3 \rho^{\sigma} \left[ (2 + x_3) \rho^2 - (2x_3 + 1) \left( \rho_n^2 + \rho_p^2 \right) \right] / 24 \\ &+ \left[ t_2 \left( 2x_2 + 1 \right) - t_1 \left( 2x_1 + 1 \right) \right] \left( \tau_n \rho_n + \tau_p \rho_p \right) / 8 + \left[ t_1 \left( 2 + x_1 \right) + t_2 \left( 2 + x_2 \right) \right] \tau \rho / 8 \\ &+ \left[ 3t_1 \left( 2 + x_1 \right) - t_2 \left( 2 + x_2 \right) \right] \left( \nabla \rho \right)^2 / 32 - \left[ 3t_1 \left( 2x_1 + 1 \right) + t_2 \left( 2x_2 + 1 \right) \right] \left[ \left( \nabla \rho_n \right)^2 + \left( \nabla \rho_p \right)^2 \right] / 32 \\ &+ W_0 \left[ \vec{J} \cdot \nabla \rho + \vec{J_n} \cdot \nabla \rho_n + \vec{J_p} \cdot \nabla \rho_p \right] / 2 + \left( t_1 - t_2 \right) \left[ J_n^2 + J_p^2 \right] / 16 - \left( t_1 x_1 + t_2 x_2 \right) J^2 / 16 . \end{aligned}$$

-9 parameters  $\{t_0, t_1, t_2, t_3, x_0, x_1, x_2, x_3, \sigma\}$ -2 purely isovector parameters:  $x_0, x_3$ 

• Relativistic Mean Field (RMF) model of nuclear matter:

$$\begin{aligned} \mathscr{L} &= \bar{\psi} \left[ \gamma^{\mu} \left( i \partial_{\mu} - g_{v} V_{\mu} - \frac{g_{\rho}}{2} \tau \cdot \mathbf{b}_{\mu} - \frac{e}{2} (1 + \tau_{3}) A_{\mu} \right) - (M - g_{s} \phi) \right] \psi + \frac{1}{2} \partial_{\mu} \phi \, \partial^{\mu} \phi - \frac{1}{2} m_{s}^{2} \phi^{2} \\ &- \frac{1}{4} V^{\mu\nu} V_{\mu\nu} + \frac{1}{2} m_{v}^{2} V^{\mu} V_{\mu} - \frac{1}{4} \mathbf{b}^{\mu\nu} \cdot \mathbf{b}_{\mu\nu} + \frac{1}{2} m_{\rho}^{2} \mathbf{b}^{\mu} \cdot \mathbf{b}_{\mu} - \frac{1}{4} F^{\mu\nu} F_{\mu\nu} - U(\phi, V_{\mu}, \mathbf{b}_{\mu}) , \\ &U(\phi, V^{\mu}, \mathbf{b}^{\mu}) = \frac{\kappa}{3!} (g_{s} \phi)^{3} + \frac{\lambda}{4!} (g_{s} \phi)^{4} - \frac{\zeta}{4!} g_{v}^{4} (V_{\mu} V^{\mu})^{2} - \Lambda_{v} g_{\rho}^{2} \mathbf{b}_{\mu} \cdot \mathbf{b}^{\mu} g_{v}^{2} V_{\nu} V^{\nu} \\ &- 7 \text{ parameters} \qquad \left\{ g_{s}, g_{v}, g_{\rho}, \kappa, \lambda, \zeta, \Lambda_{v} \right\} \\ &- 2 \text{ purely isovector parameters} \qquad g_{\rho}, \Lambda_{v} \end{aligned}$$

- Take 2 reference RMF models: NL3, IU-FSU
- Create 2 reference Skyrme models with identical saturation NM properties: SkNL3, SkIU-FSU (by writing Skyrme parameters in terms of NM parameters – Chen et al PRC80, 014322 (2009); PRC82, 024321 (2010))
- Adjust  $\rho_0$  and  $E_0$  in SHF models to reproduce double magic nuclei BE,  $r_{ch}$

	$ ho_0 ~({\rm fm}^{-3})$	$E_0$ (MeV)	$K_0$ (MeV)	$M^*_{\rm D}$ $(M)$	$M^*_{\rm L}(M)$	$M^*_{\rm S}(M)$	$M_{\rm V}^*$ (M)	J (MeV)	L (MeV)	$K_{\rm sym}$ (MeV)	$R_{\rm skin}$ (fm)
NL3*	0.1500	-16.32	258.49	0.594	0.671	-	-	38.7	122.7	105.7	0.29
SkNL3*	0.1527	-15.76	258.49	-	-	0.671	0.671	38.7	122.7	62.7	0.27
IU-FSU	0.1546	-16.40	231.33	0.609	0.687	-	-	31.3	47.2	28.5	0.16
SkIU-FSU	0.1575	-15.70	231.33	-	-	0.687	0.687	31.3	47.2	-132.0	0.16

• Re-fit the 2 purely isovector parameters in all models to the results of microscopic calculations

$$\chi^{2}(\mathbf{p}) \equiv \sum_{n=1}^{N} \left( \frac{\mathcal{O}_{n}^{(\mathrm{th})}(\mathbf{p}) - \mathcal{O}_{n}^{(\mathrm{exp})}}{\Delta \mathcal{O}_{n}} \right)^{2}$$

Model parameters:  $\mathbf{p} = (p_1, \dots, p_F)$ 

(numerical "experimental" data:

 $E_{\rm PNM}: 0.04 \le \rho \le 0.16 \ {\rm fm}^{-3}$ 

Akmal et al PRC58, 1802 (1998) Schwenk and Pethick PRL79, 160401 (2005) Gandolfi et al PRC79, 054005 (2009) Hebeler and Schwenk, PRC82, 014314 (2010))



Fattoyev, Newton, Li, PRC86, 025804

**RESULTING 1-σ CONFIDENCE ELLIPSES: SYMMETRY ENERGY PARAMETERS** 



Experiment:  $-760 < K_{\tau} < -372$  MeV e.g. Dutra et al, PRC85, 035201 (2012)



Fattoyev, Newton, Li, PRC86, 025804

• When constrained by PNM calculations, RMF models are systematically stiffer at high density than SHF models.

### Tidal polarizability and Love number: sensitivity to high-p symm. energy



Fattoyev, Carvajal, Newton, Li, PRC87, 15806 (2013)

#### Tidal polarizability and Love number: sensitivity to high-p symm. energy



Variation of saturation density properties/high density SNM



Fattoyev, Carvajal, Newton, Li, PRC87, 15806 (2013)

- Detector sensitivities assuming
   Optimally oriented, equal mass binary at D=100 Mpc
- Damour, Nagar, PRD81, 084016 (2010)
- Damour, Nagar, Villain, PRD85, 123007 (2012)

$$\Delta \tilde{\lambda} \approx \alpha \left(\frac{M}{M_{\odot}}\right)^{2.5} \left(\frac{m_2}{m_1}\right)^{0.1} \left(\frac{f_{\rm end}}{\rm Hz}\right)^{-2.2} \left(\frac{D}{100 \rm Mpc}\right)$$

Hinderer et al, PRD 81, 123016 (2010)

 At 1.4M<sub>SUN</sub>, high density behavior of symmetry energy at limit of AdvLIGOs ability to constrain

### Dynamical tides I





• Tidal field E<sub>ij</sub> resonates low frequency g-modes, inertial modes ≈100Hz

 $\omega_{lpha} = m\Omega_{
m orb}, \quad m = 2, 3, \cdots$ 

- Resulting energy transfer appears as phase shift in gravitational waveform
- Estimated to produce negligible phase shift dN < 0.1, but only estimated for R=10km
- But...  $\Delta N \propto R^4$
- Radius/symmetry energy measurements determine whether we should worry

Lai, D., MNRAS 270, (1994) Ho, W.C.G., Lai, D., MNRAS 308 (1999) Lai, D., Wu, Y., PRD 74, 024007 (2006)

- NS-NS mergers strong candidates for sGRBs
- Precursor flares observed 1-10s before 4 GRBs
- Possible interpretation: crust shattering by tidal excitation of crustal interface mode ≈100Hz (Tsang et al PRL108, 2012)



### Summary

- Explored sensitivity of tidal polarizability to high density symmetry energy
  - Sample models: RMFs/Skyrmes with same saturation properties, fit to PNM
  - RMFs give systematically stiffer symmetry energy at high density
- High density symmetry energy behavior of models can be distinguished (just) by Adv. LIGO
- Tidal field/g-mode resonance: change in GW waveform assumed negligible, but symmetry energy/radius measurements needed to bolster confidence
- Tidal field/crustal interface mode resonance could shatter crust
  - EM signature: precursor flares to sGRBs?
  - If so, favors mid-high saturation stiffness