

Heavy ion collisions における quark number scaling

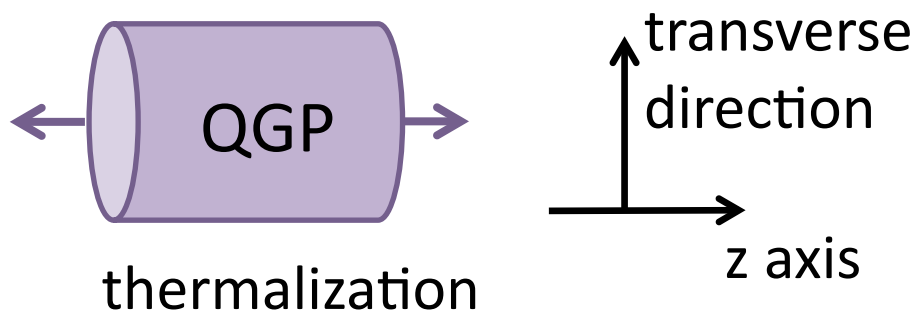
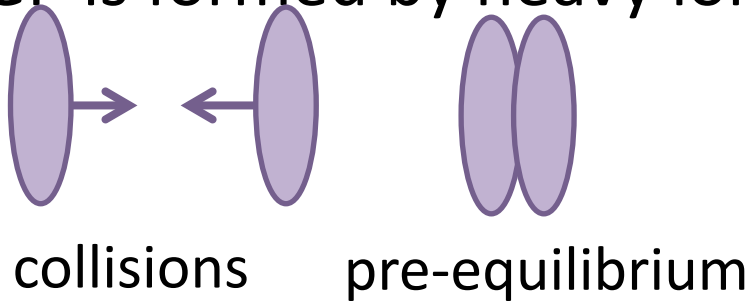
R.J.Fries, B.Muller, S.A.Bass and C.Nonaka
PHYSICAL REVIEW C 68, 044902 (2003)

京都大学 原子核理論

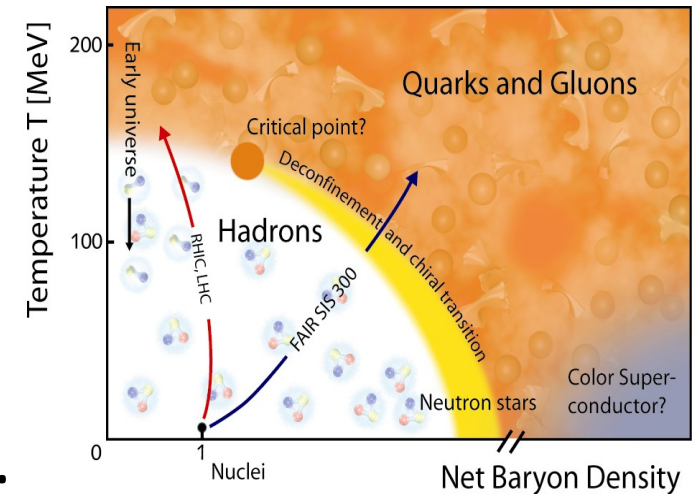
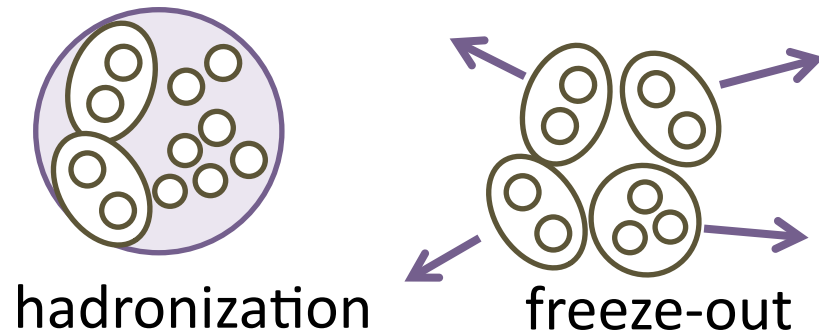
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QGP

- QGP: quark gluon plasma at RHIC.
- QGP is formed by heavy ion collisions.



thermalization
Hydrodynamical expansion



The Physics of High Baryon Densities
International Workshop at [ECT*](#) in
Trento May 29 - June 2, 2006

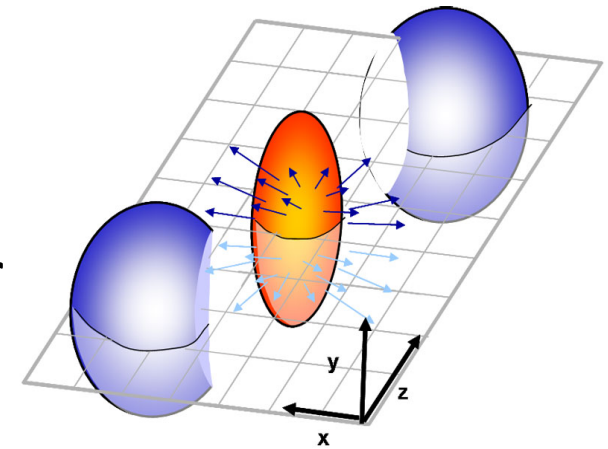
- One of the important measurement is **elliptic flow**.

Elliptic flow v_2

- To know particle emission dependence on azimuthal angle, expand in a Fourier cosine series.

$$E \frac{d^3 N}{d^3 P} = \frac{d^2 N}{2\pi P_T dP_T dy} \left(1 + \sum_n 2v_n \cos n\phi \right)$$

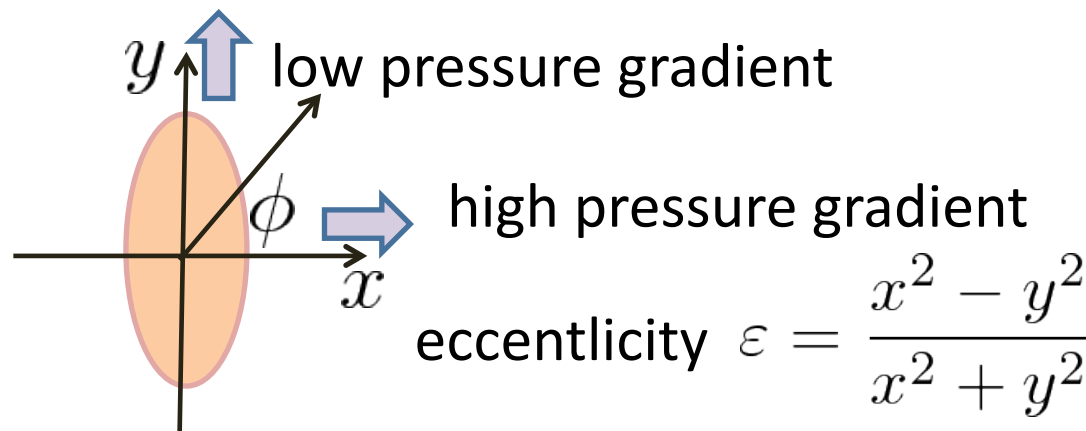
– elliptic flow v_2



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October 25, 2006

y : rapidity P_T : transverse momentum

- Elliptic flow reflects initial geometrical anisotropy.
QGP: hydrodynamics \Rightarrow difference of pressure gradient

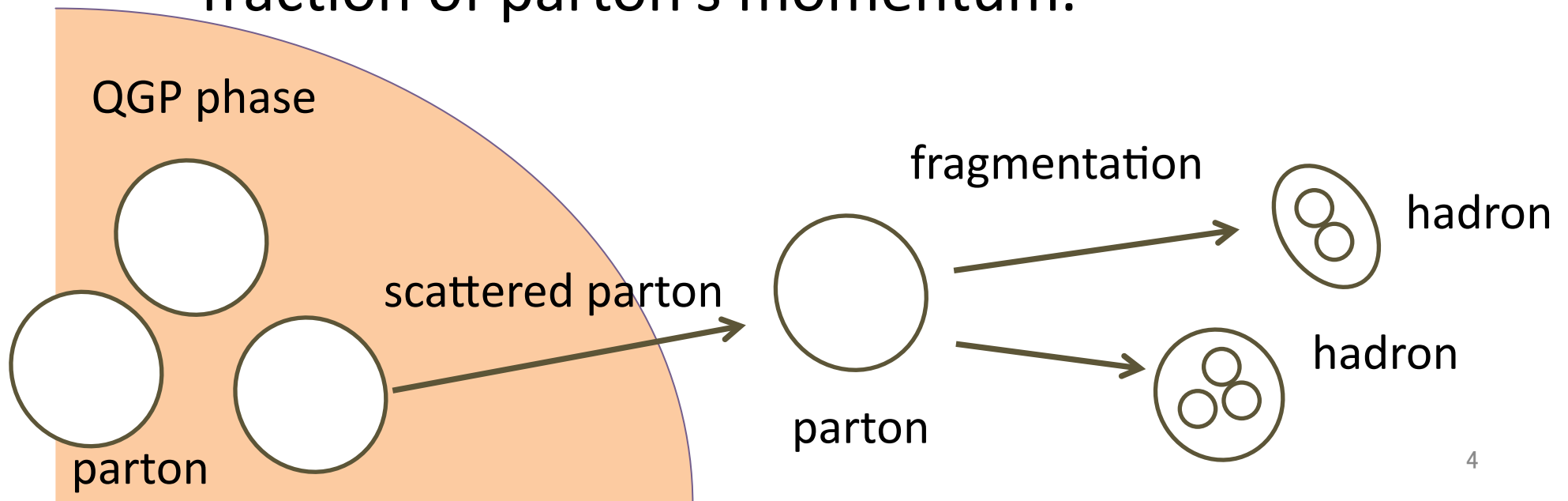


$$v_2 = \langle \cos 2\phi \rangle = \left\langle \frac{p_x^2 - p_y^2}{p_x^2 + p_y^2} \right\rangle$$

Existing mechanism

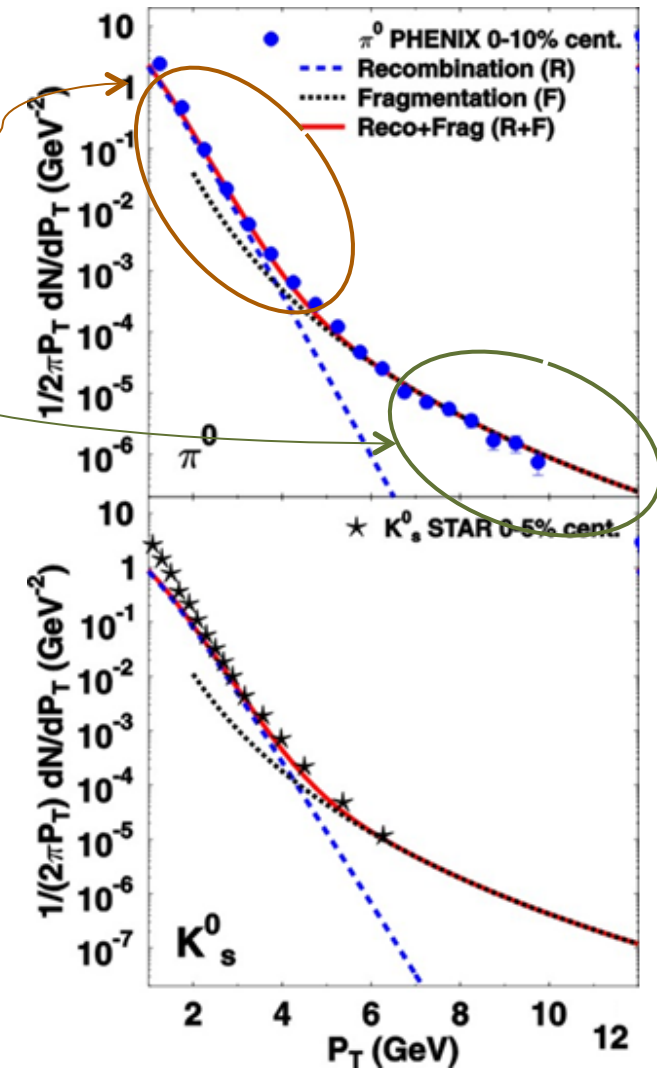
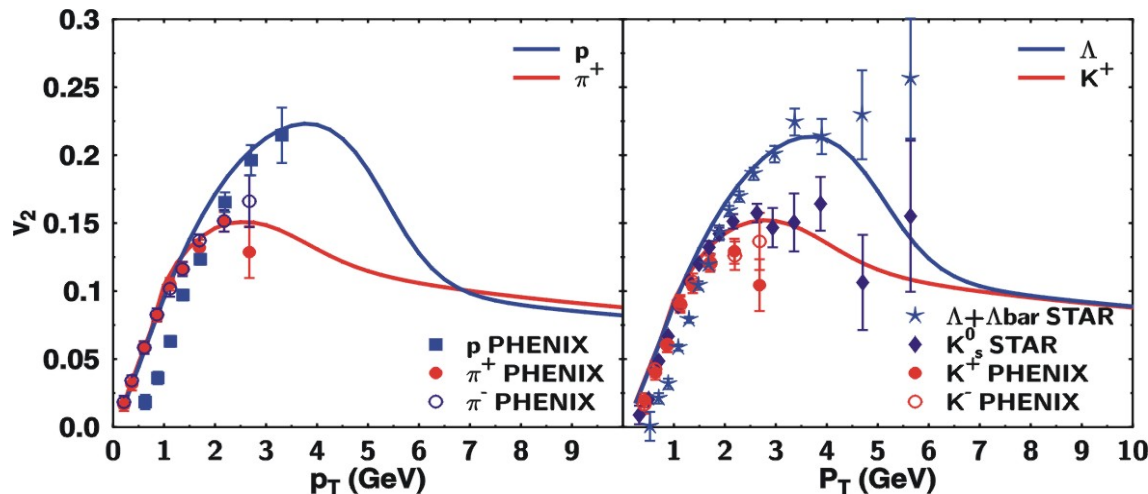
~ fragmentation

- **Hadron production** at sufficient large momentum can be described by **pQCD**.
- **Fragmentation** mechanism is a parton fragments into hadron carrying with a certain fraction of parton's momentum.



Fragmentation to Recombination

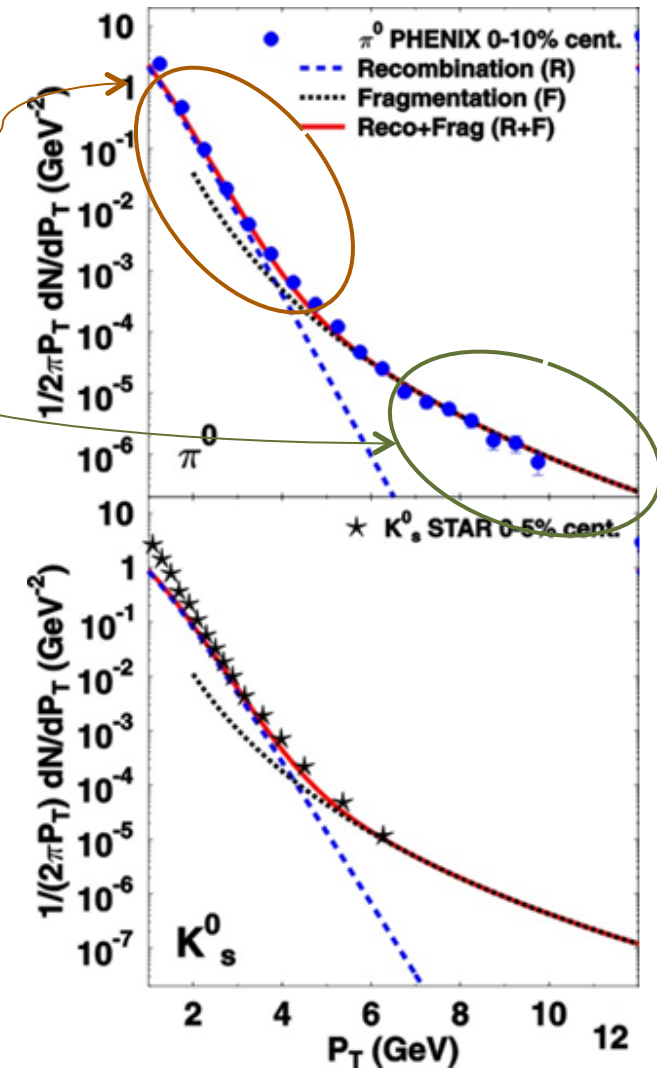
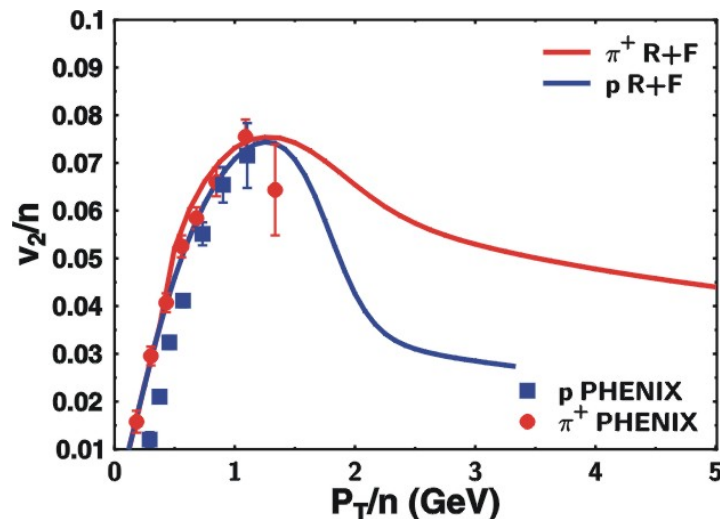
- The two component form of hadron transverse spectra, including an **exponential part (thermalized phase)** and a **power law tail (pQCD)**.
- The particle dependence of the elliptic flow v_2 .



Fragmentation to Recombination

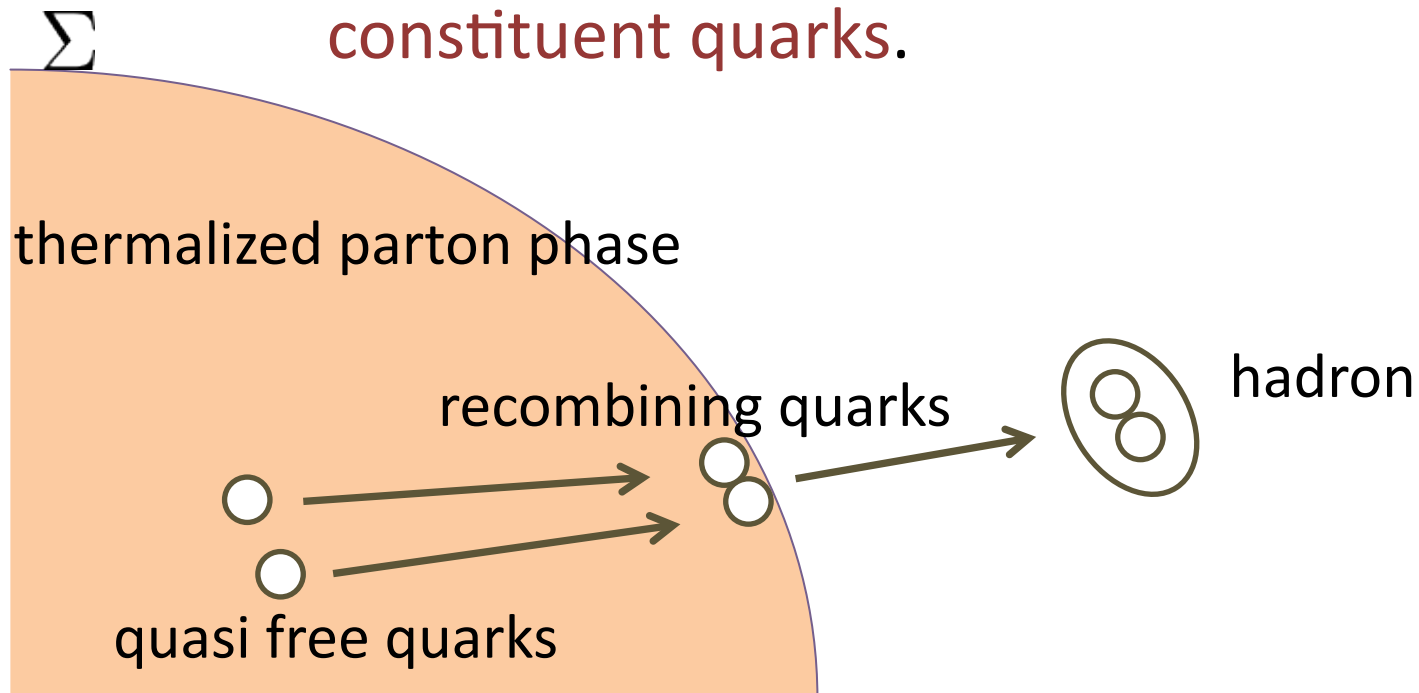
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- The particle dependence of the elliptic flow v_2 .

➔ Quark number scaling



Recombination model

- Assumption
 - Recombination takes place in **thermal phase**.
 - **An instantaneous recombination**, corresponding to an infinitely thin hypersurface Σ .
 - **No dynamical gluons**. Taking into account **only constituent quarks**.



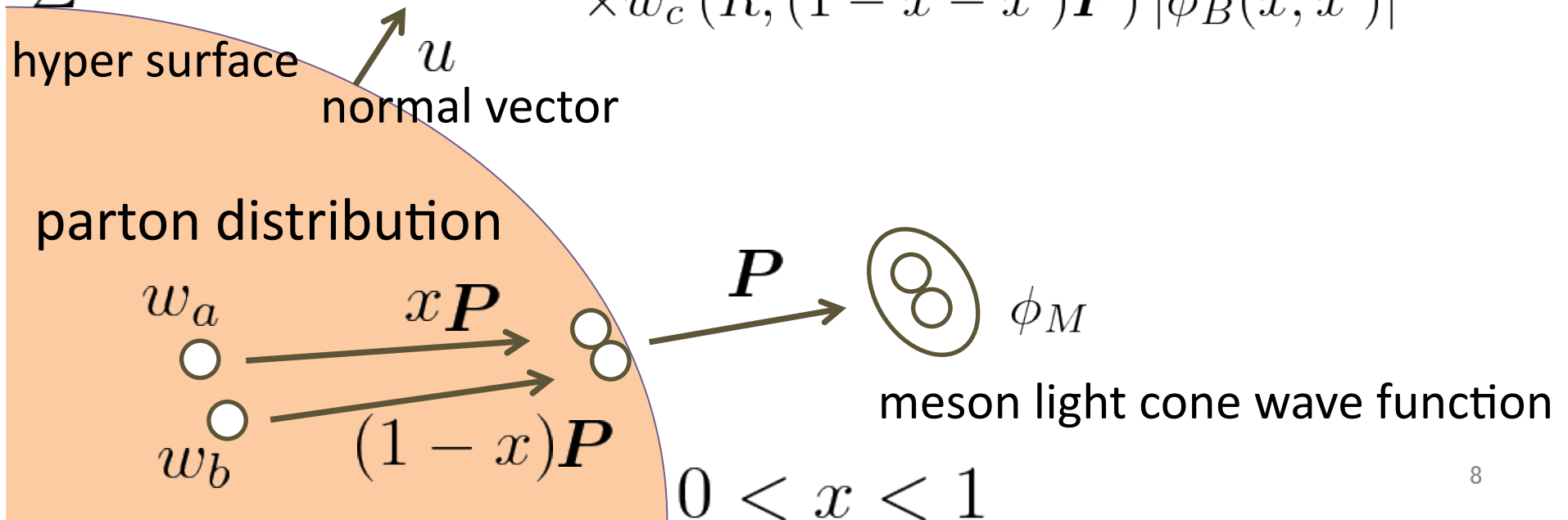
Recombination formalism

- Meson and Baryon spectrum is given by

$$E \frac{dN_M}{d^3 P} = \sum_{a,b} \int_{\Sigma} \frac{d\sigma P \cdot u}{(2\pi)^3} \int_0^1 dx w_a(R, x\mathbf{P}) w_b(R, (1-x)\mathbf{P}) |\phi_M(x)|^2$$

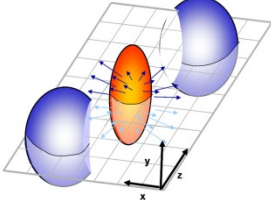
$$E \frac{dN_B}{d^3 P} = \sum_{a,b,c} \int_{\Sigma} \frac{d\sigma P \cdot u}{(2\pi)^3} \int_0^1 dx dx' w_a(R, x\mathbf{P}) w_b(R, x'\mathbf{P})$$

$$\Sigma \times w_c(R, (1-x-x')\mathbf{P}) |\phi_B(x, x')|^2$$



scaling law of elliptic flow $v_2 = \langle \cos 2\phi \rangle$

- Assuming quarks and anti-quarks have a pure elliptic flow.

$$w_a(R, P) \rightarrow w_a(R, P) [1 + 2v_{2,q}(P_T) \cos 2\phi]$$


- In the thermalized region, elliptic flow of mesons and baryons are

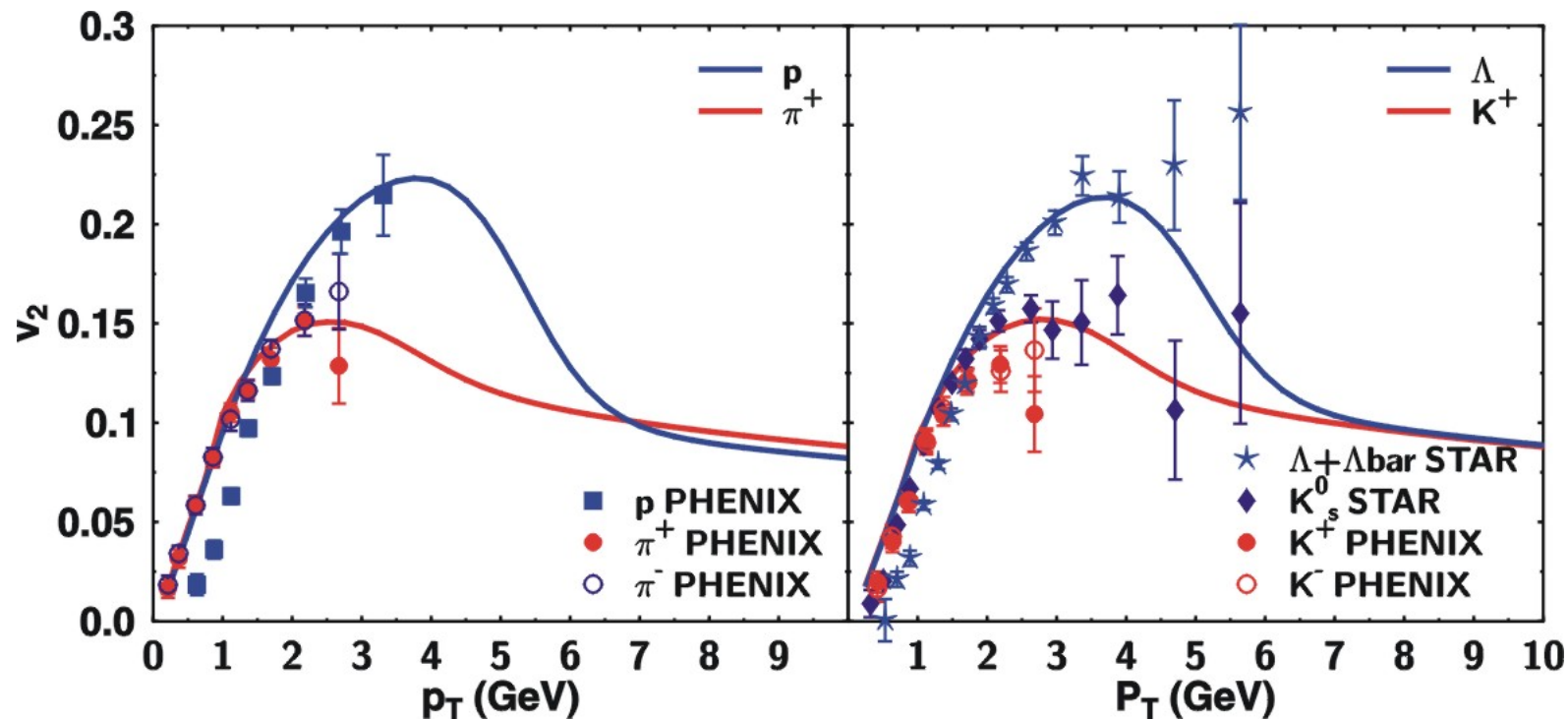
$$v_2^M(P_T) = \frac{2v_{2,q}\left(\frac{P_T}{2}\right)}{1 + 2v_{2,q}\left(\frac{P_T}{2}\right)^2}, \quad v_2^B(P_T) = \frac{3v_{2,q}\left(\frac{P_T}{3}\right) + 3v_{2,q}\left(\frac{P_T}{3}\right)^3}{1 + 6v_{2,q}\left(\frac{P_T}{3}\right)^2}$$

- Ignoring quadratic and cubic terms, we can arrive at following simple **scaling law**:

$$v_2^H = nv_{2,q} \left(\frac{1}{n} P_T \right), \quad \text{where } n \text{ is \# of constituent quarks.}$$

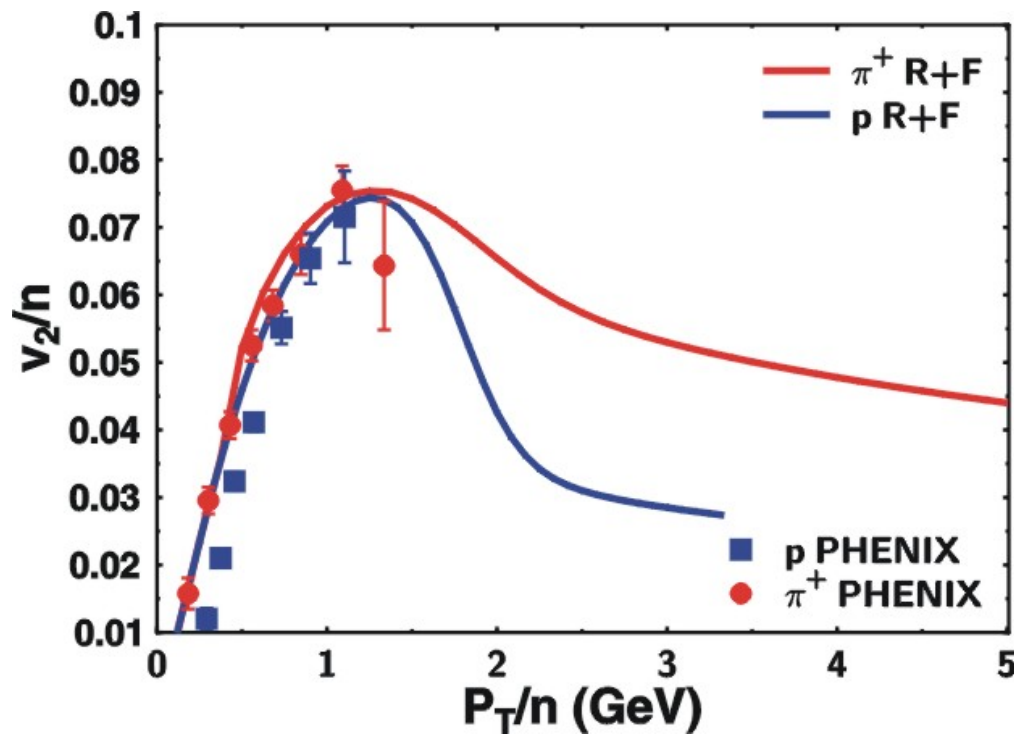
Results

- The particle dependence of elliptic flow.
- The different behavior of mesons.
 - protons and pions. kaons and Λ 's.



Scaling law

- Follow one universal curve.
- Scaling law breaks down in the pQCD domain.



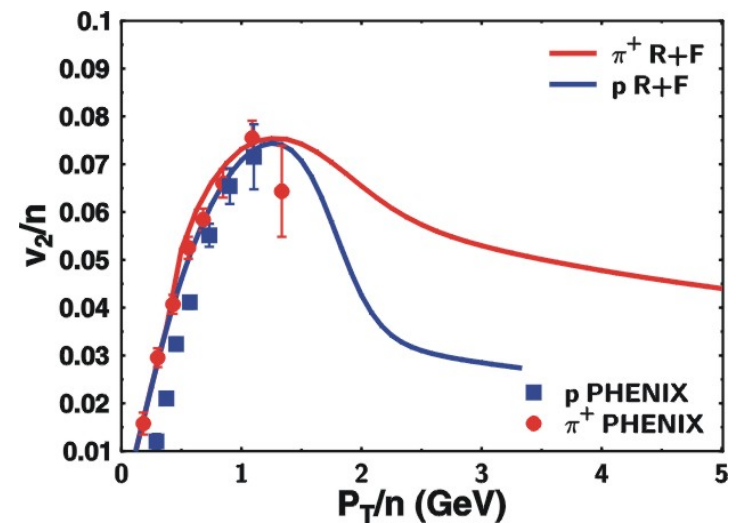
$$v_2^H = n v_{2,q} \left(\frac{1}{n} P_T \right)$$

Conclusions

- Recombination model explains scaling law.



- Recombination model implies ...
 - the existence of **thermalized parton phase (QGP)**.
 - hydrodynamical evolution on a **quark level!**



Back up

Modeling parton phase

- Low P_T domain

- The thermal region is parameterized as :

$$w_a = \gamma_a e^{-p \cdot v/T} e^{-\eta^2/2\Delta^2} f(\rho, \phi)$$

γ_a : a flavor fugacity factor

T : temperature

$f(\rho, \phi)$: transverse distribution

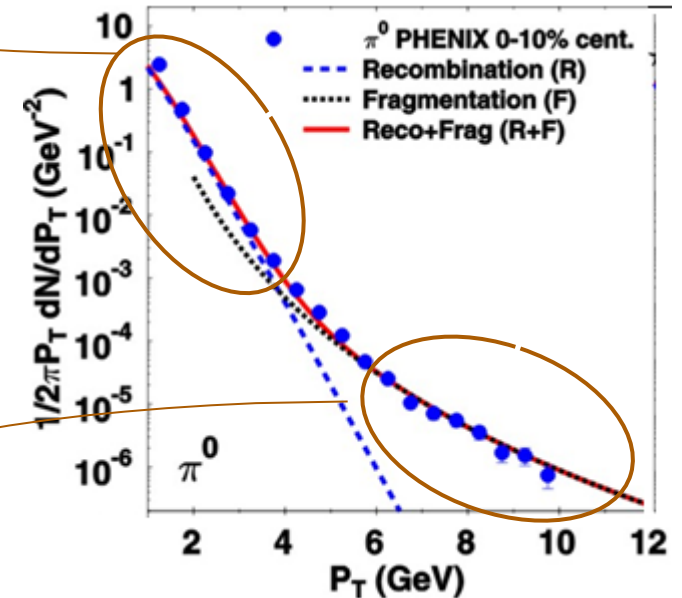
Δ : rapidity distribution width

- High P_T domain

- The pQCD component is parameterized as:

$$\left. \frac{dN_a}{P_T dP_T dy} \right|_{y=0} = K \frac{C}{(1 + P_T/B)^\beta}$$

C , B and β is taken from a pQCD.



Recombination formalism 1

- By introducing the **density matrix** $\hat{\rho}$ the number of quark-antiquark states that we interpret as mesons is given by

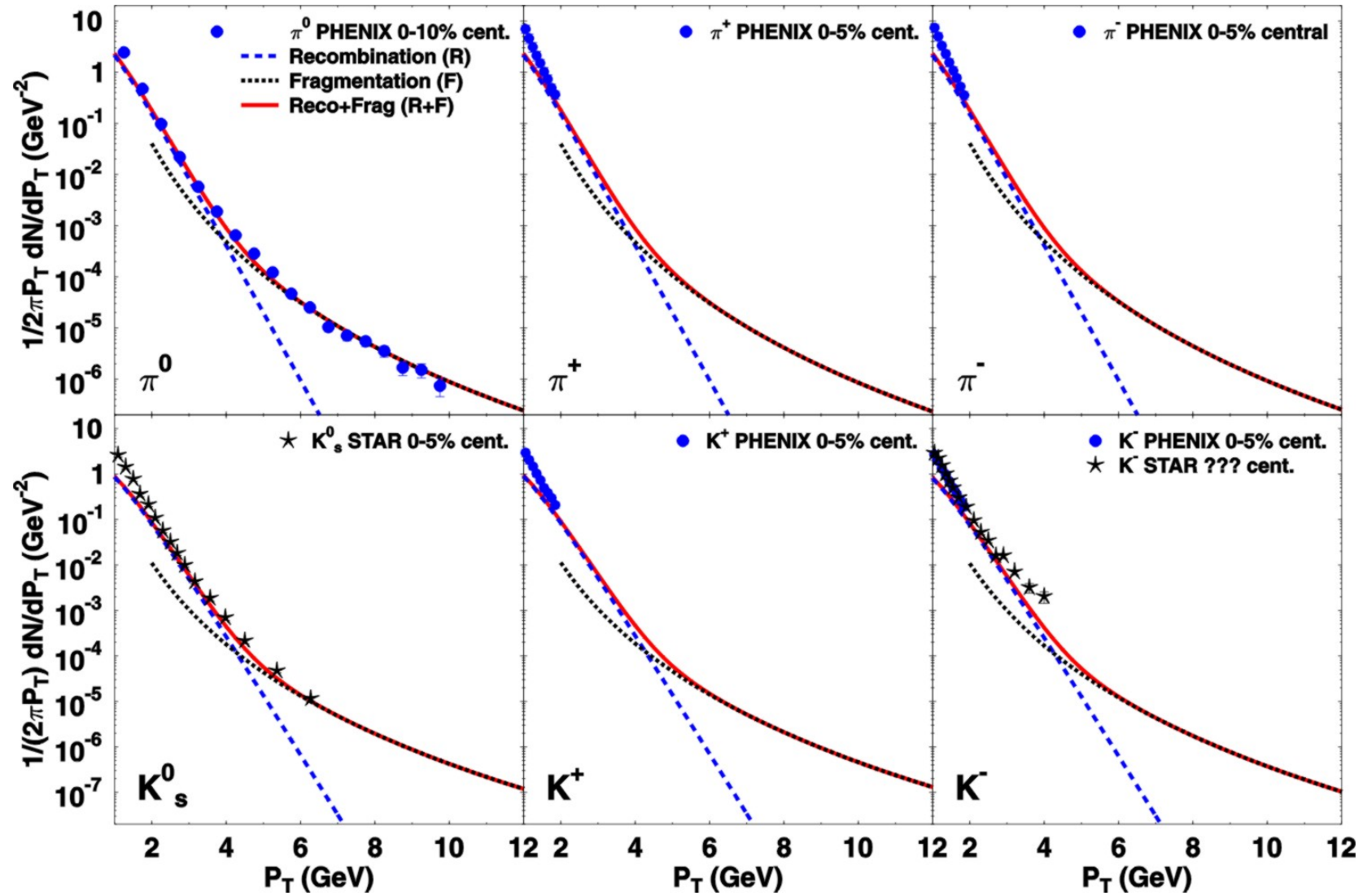
$$N_M = \sum_{a,b} \int \frac{d^3 P}{(2\pi)^3} \langle M; \mathbf{P} | \hat{\rho}_{ab} | M; \mathbf{P} \rangle$$

Sum is over all combination of quantum numbers
– flavor, helicity and color.

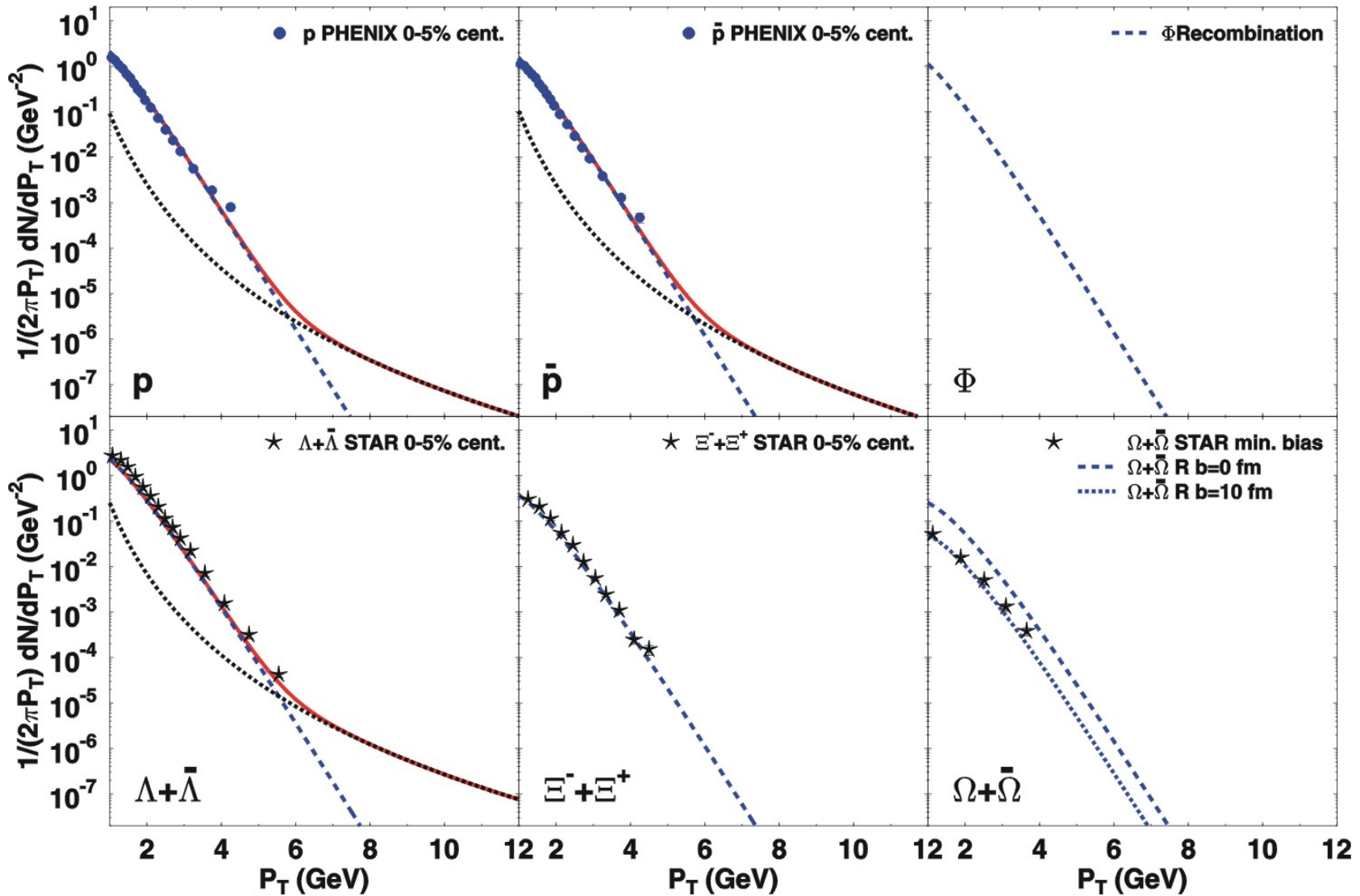
- Introducing **Wigner functions** W, Φ for two-quarks and mesons respectively.

$$\frac{dN_M}{d^3 P} = \sum_{a,b} \int \frac{d^3 R}{(2\pi)^3} \int \frac{d^3 q d^3 r}{(2\pi)^3} W_{ab} \left(\mathbf{R} + \frac{\mathbf{r}}{2}, \mathbf{R} - \frac{\mathbf{r}}{2}, \frac{\mathbf{P}}{2} + \mathbf{q}, \frac{\mathbf{P}}{2} - \mathbf{q} \right) \Phi_M^W(\mathbf{r}, \mathbf{q})$$

Hadron spectra 1



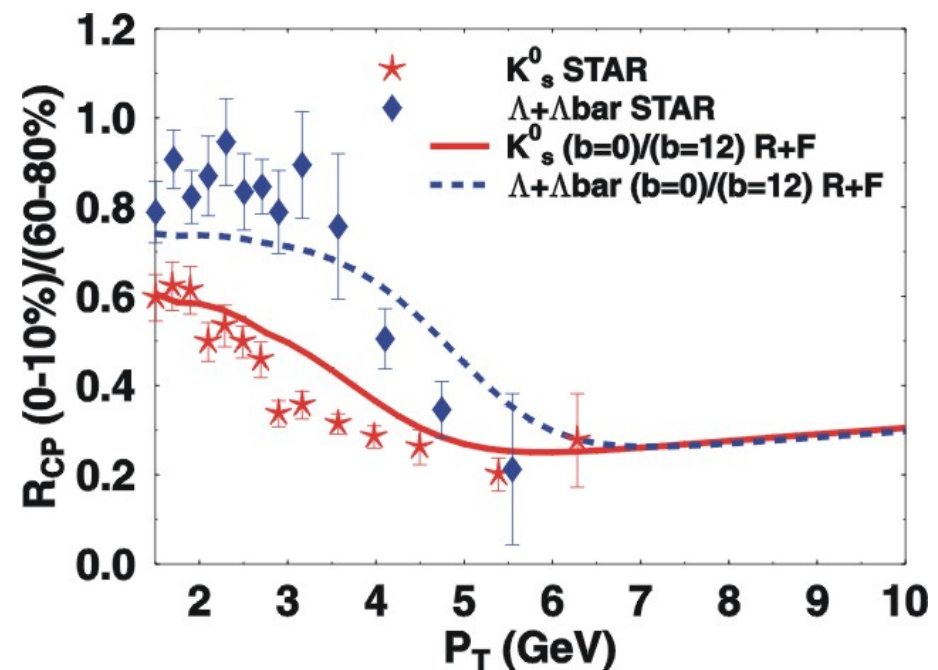
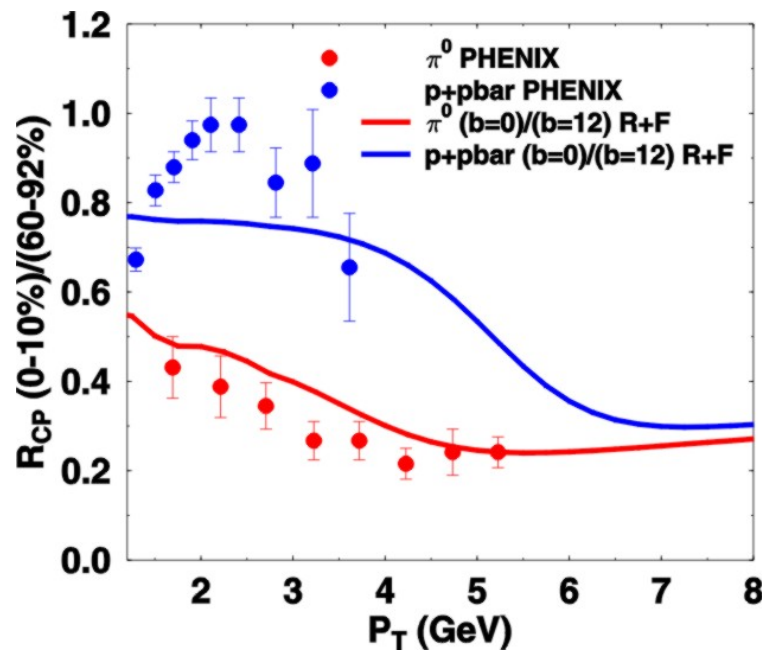
Hadron spectra 2



R_{CP}

- The ratio of particle production in central and peripheral collisions normalized by the number of binary collisions.

$$R_{CP} = \frac{\frac{1}{N_{\text{binary}}^{\text{central}}} \frac{dN^{\text{central}}}{dP_T}}{\frac{1}{N_{\text{binary}}^{\text{peripheral}}} \frac{dN^{\text{peripheral}}}{dP_T}}$$



Rapidity

- Boost along the z axis,

$$\beta = \frac{\beta_1 + \beta_2}{1 + \beta_1\beta_2}$$

- Rapidity : $y = \arctan \beta$

$$\beta = \frac{\beta_1 + \beta_2}{1 + \beta_1\beta_2} \Rightarrow y = y_1 + y_2$$

- other expressions

$$y = \frac{1}{2} \ln \frac{1 + \beta}{1 - \beta} = \frac{1}{2} \ln \frac{E + P_z}{E - P_z}$$

- For small β ,
 $y \approx \beta$

Light cone coordinates

- Light cone coordinates

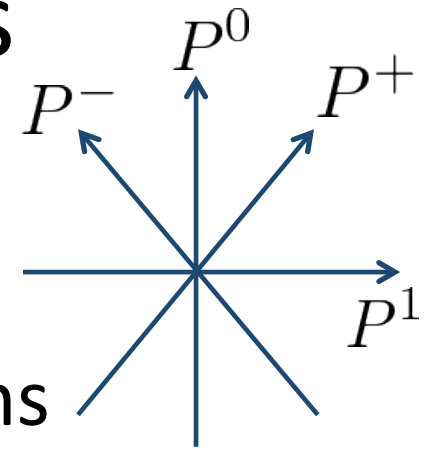
$$x^+ = \frac{1}{\sqrt{2}} (x^0 + x^1) \quad x^- = \frac{1}{\sqrt{2}} (x^0 - x^1)$$

- The line elements in $D=d+1$ dimensions

$$ds^2 = -2dx^+ dx^- + (dx^2)^2 + \dots + (dx^d)^2$$

- Momentum on the light cone

$$P^+ = \frac{1}{\sqrt{2}} (P^0 + P^1) \quad P^- = \frac{1}{\sqrt{2}} (P^0 - P^1)$$



fragmentation

- The invariant cross section for a hadron h with momentum P can be given in factorized form,

$$E \frac{d\sigma_h}{d^3 P} = \sum_{a,b} \int_0^1 \frac{dx}{x^2} D_{a \rightarrow h}(x) \frac{d\sigma_a}{d^3 (P/x)}$$

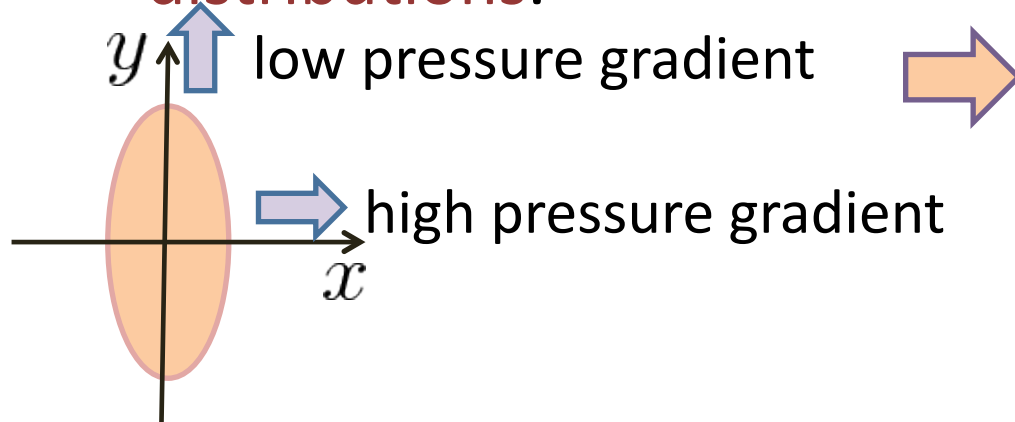
- $D_{a \rightarrow h}$: fragmentation function
- x : momentum fraction carried by each quark

QGP

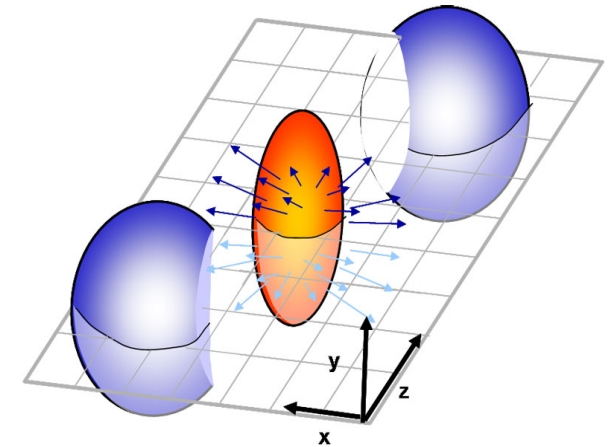
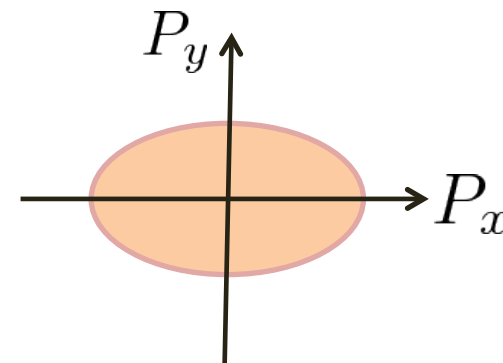
- QGP is formed by heavy ion collisions.
- QGP is described by **hydrodynamics**, which leads to the difference of pressure gradient.



- The anisotropy of hadron momentum distributions.



Azimuthal distribution



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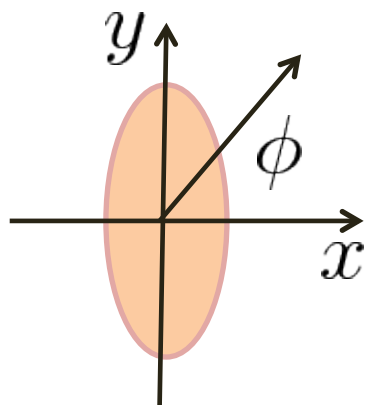
Elliptic flow v_2

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$$E \frac{d^3 N}{d^3 P} = \frac{d^2 N}{2\pi P_T dP_T dy} \left(1 + \sum_n 2v_n \cos n\phi \right) \quad \begin{array}{l} y : \text{rapidity} \\ P_T : \text{transverse} \\ \text{momentum} \end{array}$$

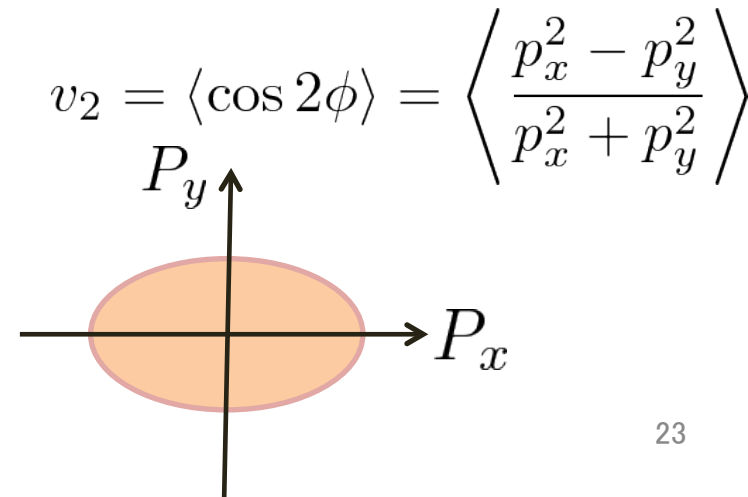
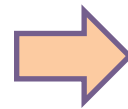
- Elliptic flow reflects initial geometrical anisotropy.

$$v_2 = \langle \cos 2\phi \rangle = \left\langle \frac{p_x^2 - p_y^2}{p_x^2 + p_y^2} \right\rangle$$



eccentricity

$$\varepsilon = \frac{x^2 - y^2}{x^2 + y^2}$$



scaling law of elliptic flow v_2

$$v_2 = \langle \cos 2\phi \rangle$$

- In the thermalized region, elliptic flow of mesons and baryons

$$v_2^M(P_T) = \frac{2v_2 \left(\frac{P_T}{2}\right)}{1 + 2v_2 \left(\frac{P_T}{2}\right)^2} \quad v_2^B(P_T) = \frac{3v_2 \left(\frac{P_T}{3}\right) + 3v_2 \left(\frac{P_T}{3}\right)^3}{1 + 6v_2 \left(\frac{P_T}{3}\right)^2}$$

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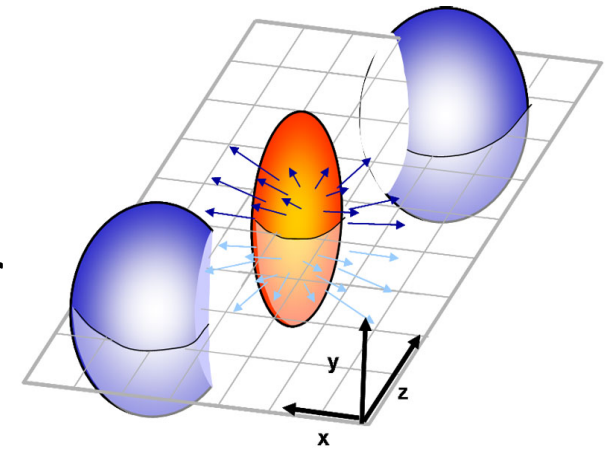
$$E \frac{dN_B}{d^3 P} = \sum_{a,b,c} \int_{\Sigma} \frac{d\sigma P \cdot u}{(2\pi)^3} \int_0^1 dx dx' w_a(R, x\mathbf{P}) w_b(R, x'\mathbf{P}) \\ \times w_c(R, (1-x-x')\mathbf{P}) |\phi_B(x, x')|^2$$

- Choose a hypersurface Σ for hadronization.
- w_a : single particle distribution function for quarks at hadronization.
- ϕ_M and ϕ_B : light cone wave function for the meson and baryon.
- x, x' : momentum fraction carried by each quark.
- u : the future oriented unit vector orthogonal to the hypersurface.

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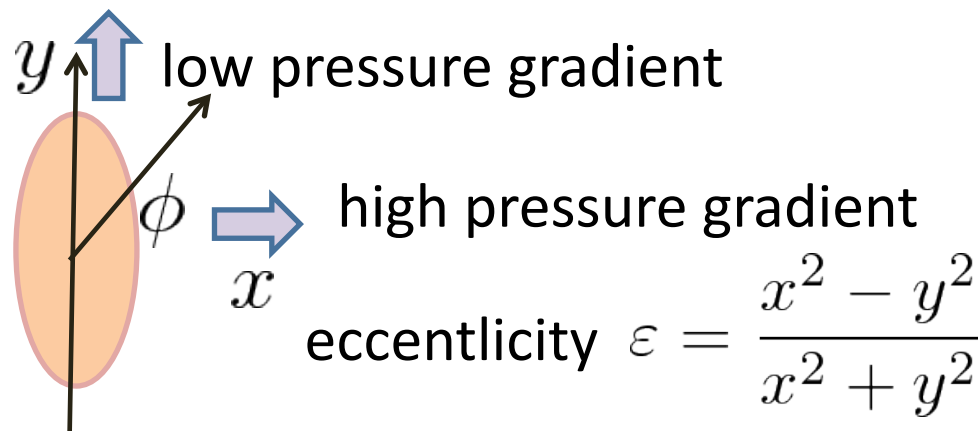
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