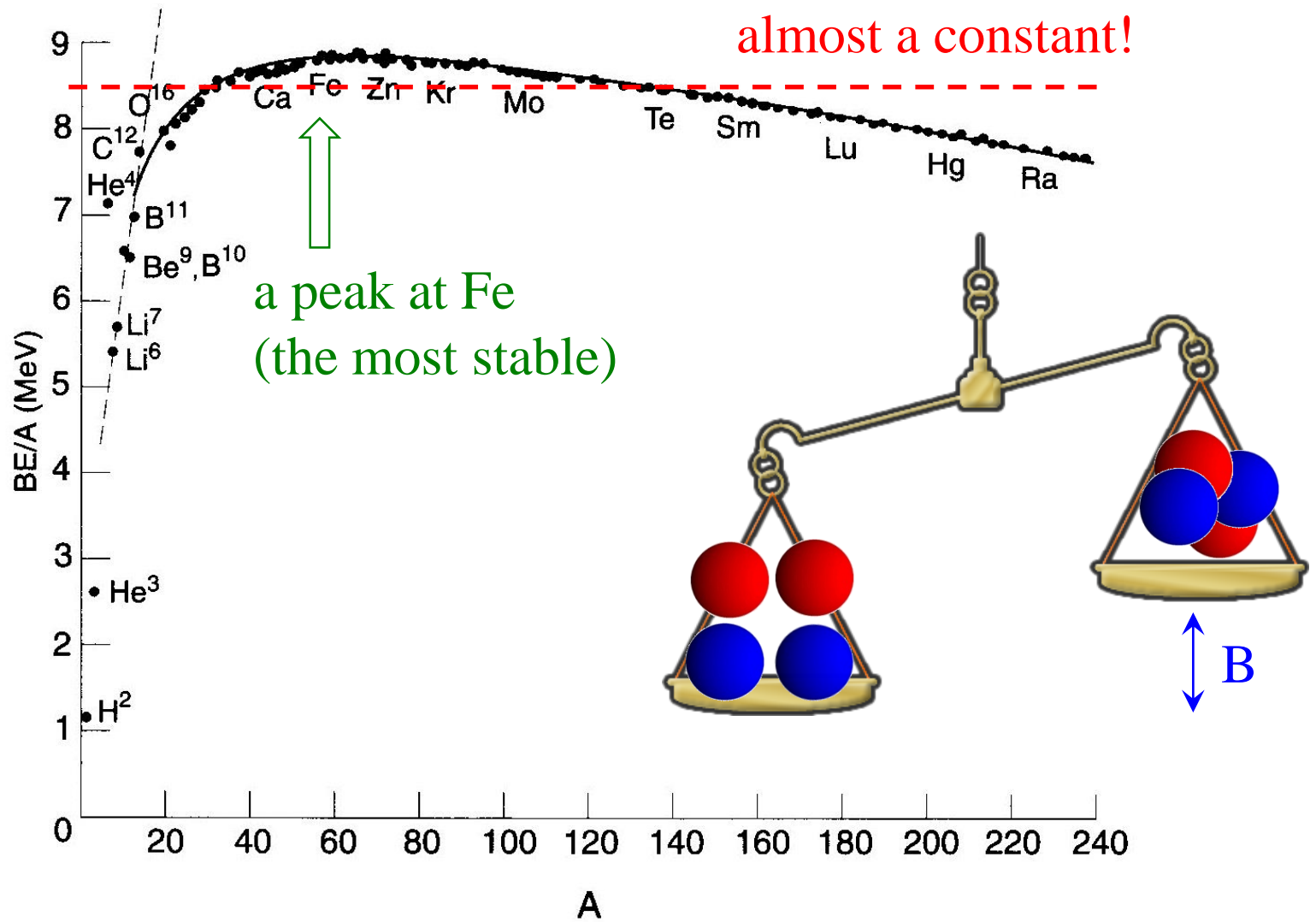
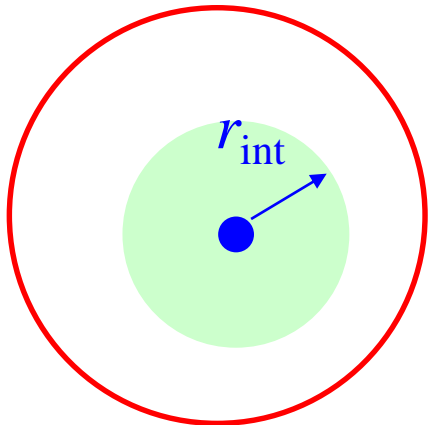


the origin of the peak

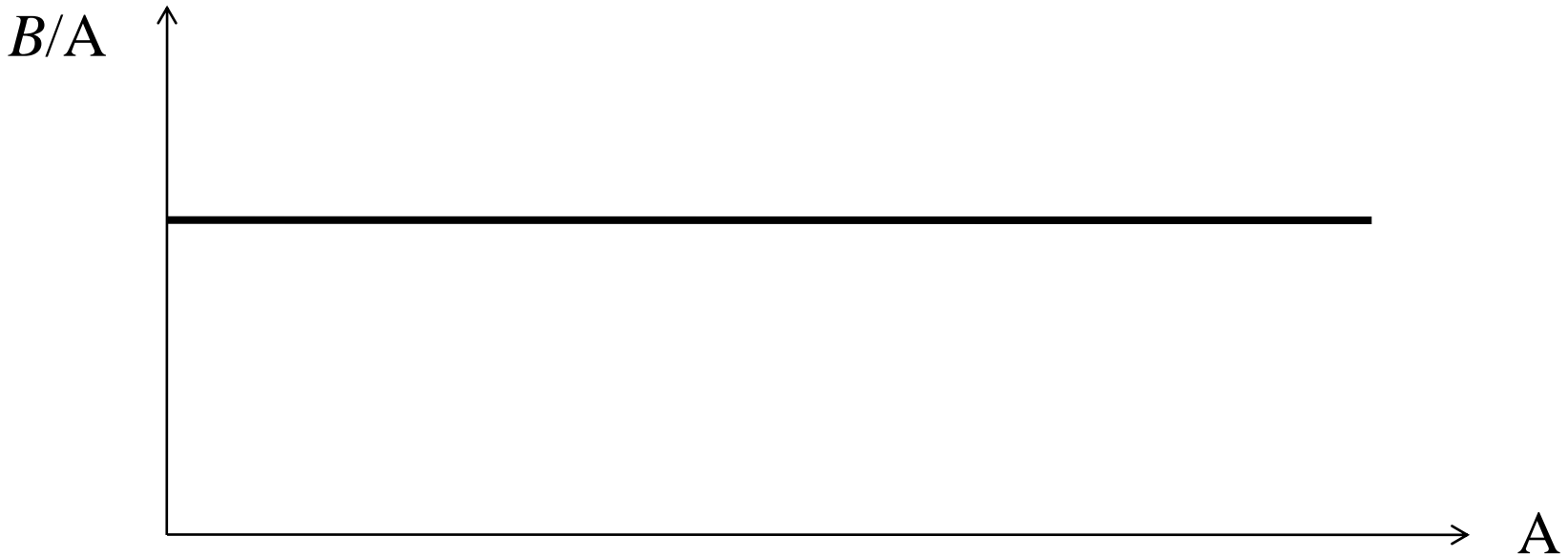


if each nucleon can interact only α -nucleons close by:

$$B \sim \alpha A/2 \longrightarrow B/A \sim \alpha/2 \text{ (const.)}$$

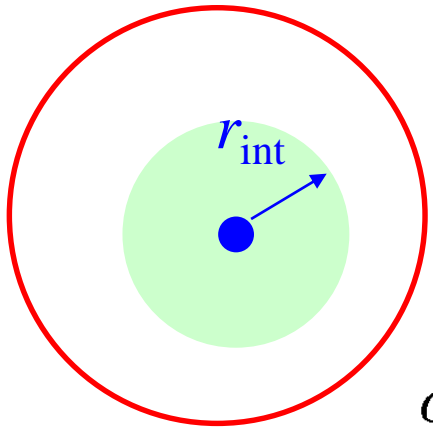


$$\alpha = \frac{4\pi}{3} r_{\text{int}}^3 \cdot \rho$$



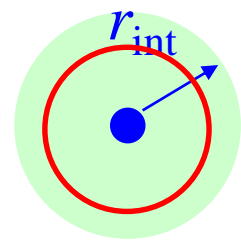
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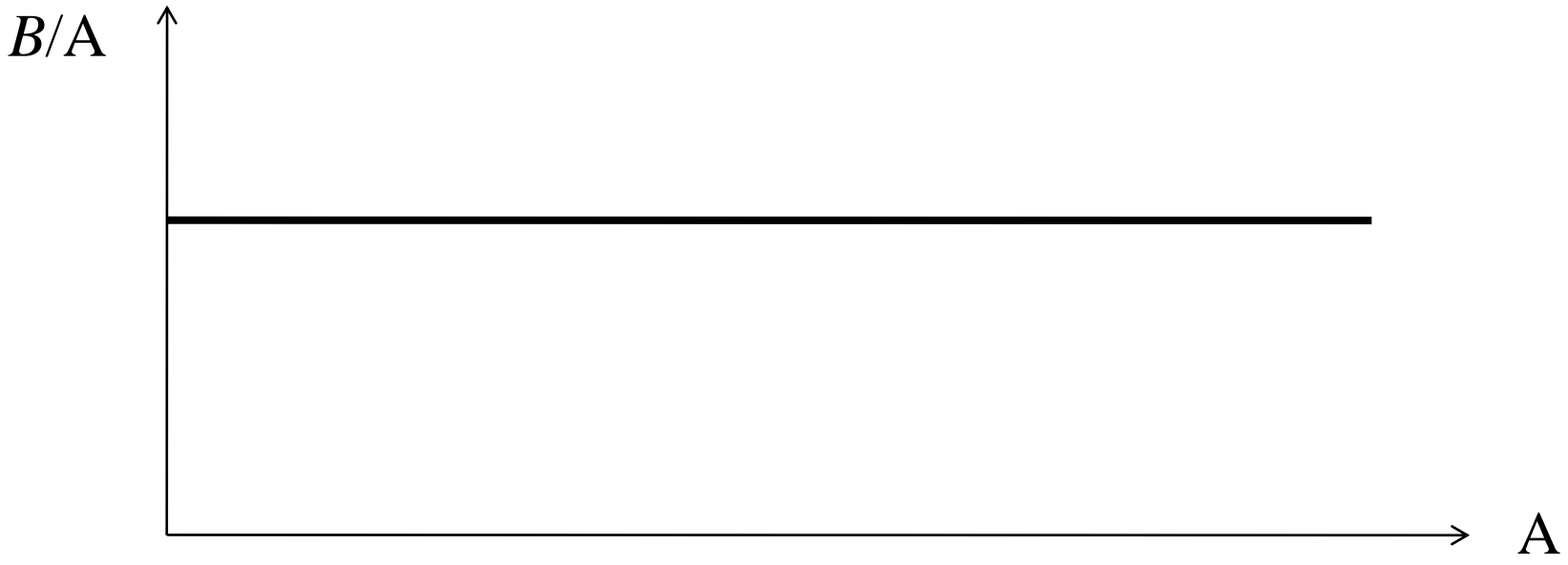


$$\alpha = \frac{4\pi}{3} r_{\text{int}}^3 \cdot \rho$$

a small nucleus

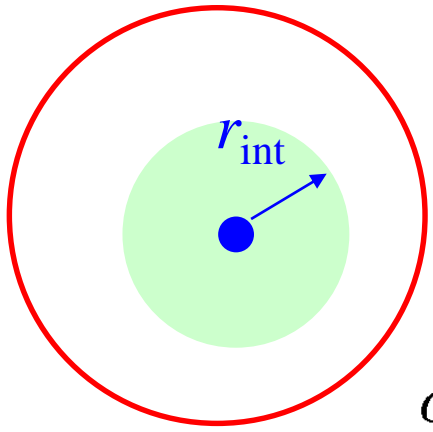


$$\rightarrow B/A \propto A - 1$$



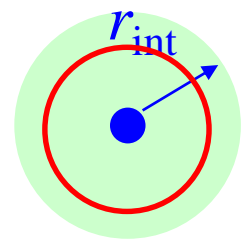
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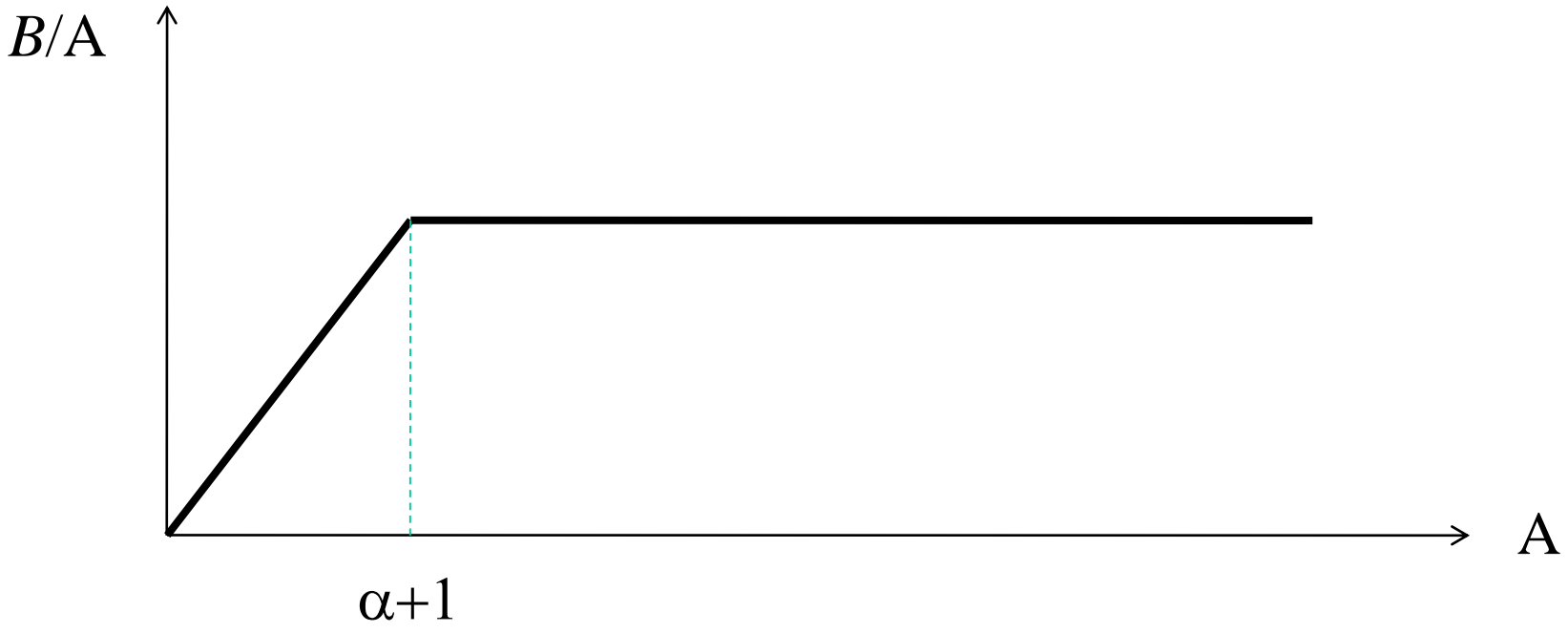


$$\alpha = \frac{4\pi}{3} r_{\text{int}}^3 \cdot \rho$$

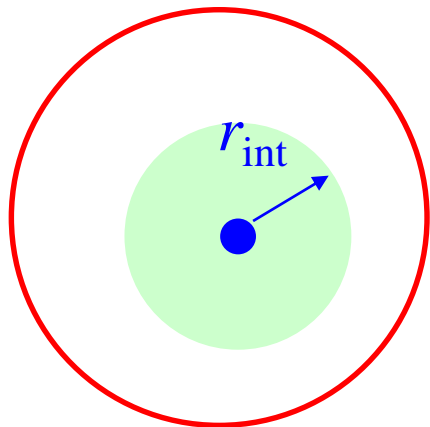
a small nucleus



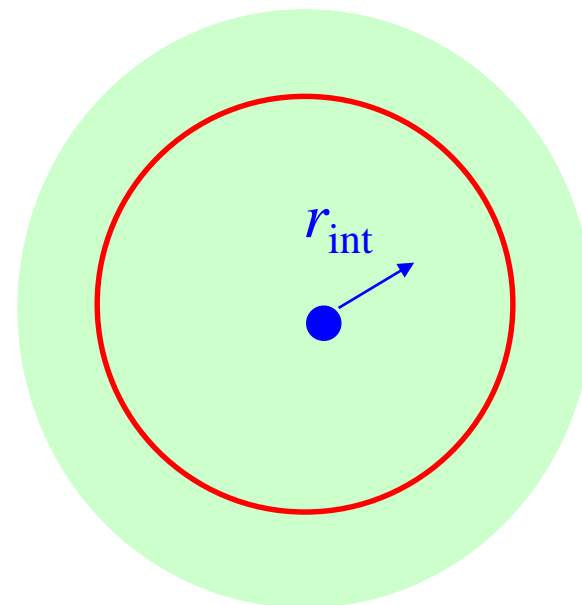
$$\rightarrow B/A \propto A - 1$$



nuclear interaction



Coulomb interaction



B/A

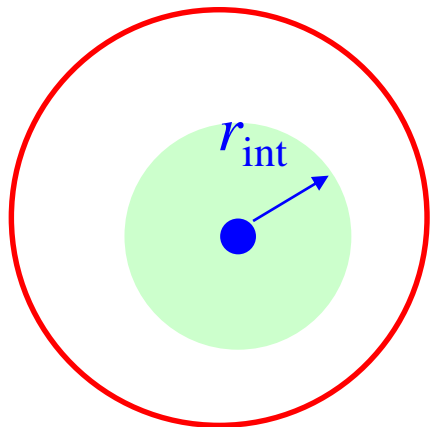


$\alpha+1$

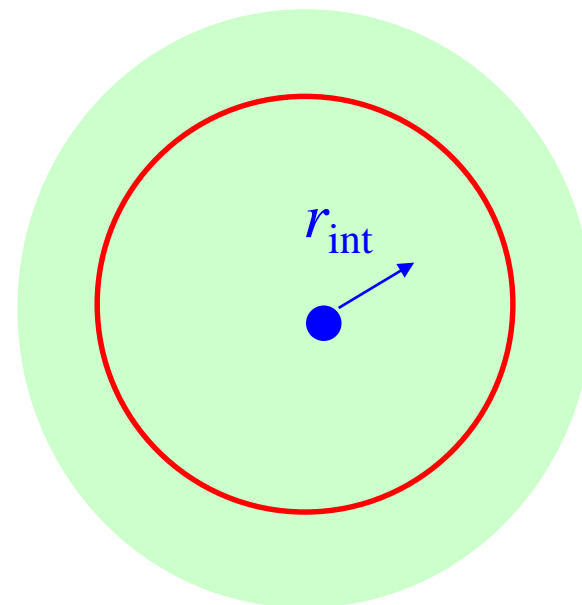
$\rightarrow B/A \propto A - 1$

A

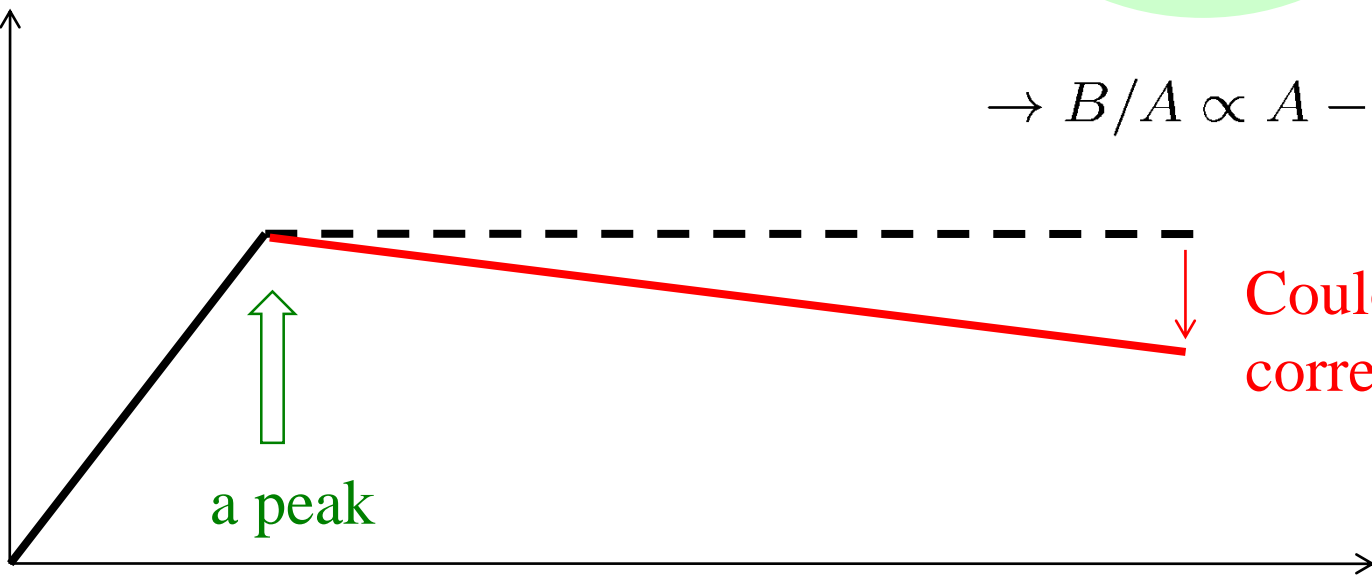
nuclear interaction



Coulomb interaction



B/A



$$\rightarrow B/A \propto A - 1$$

Coulomb correction

a peak

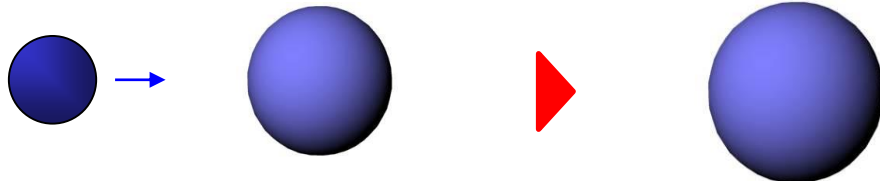
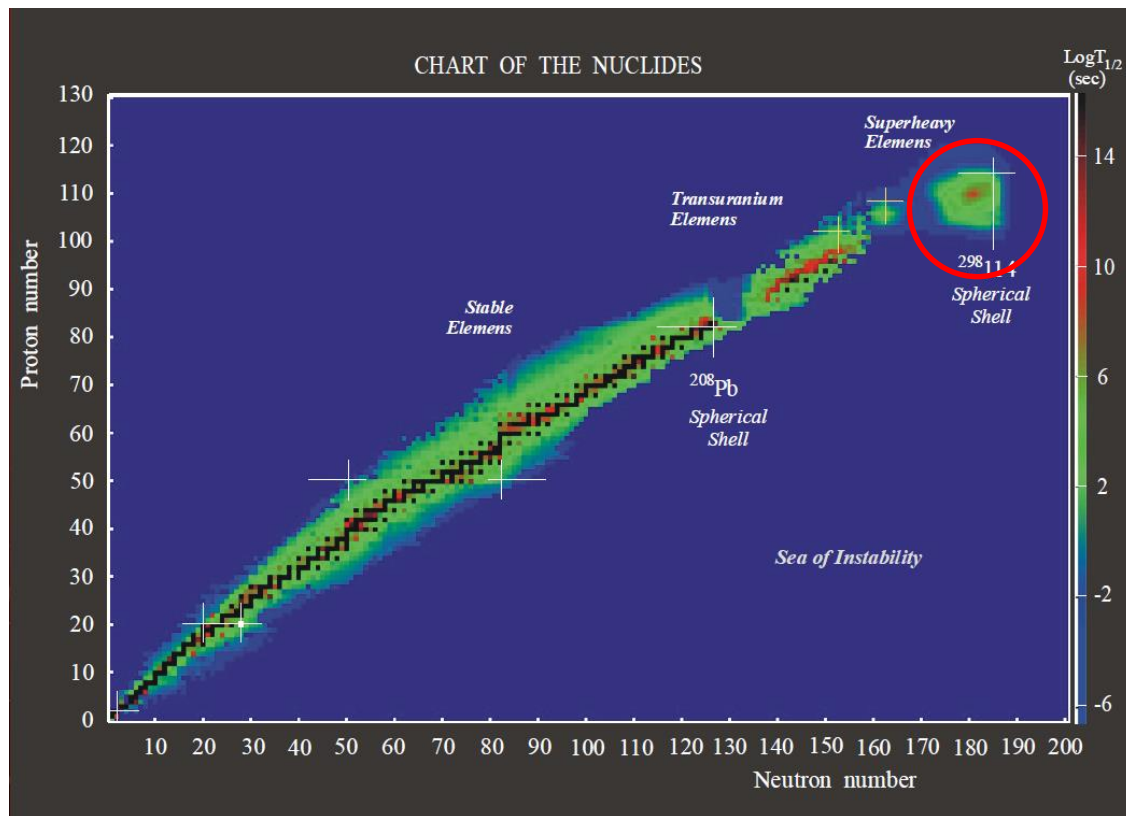
A

Superheavy elements



November, 2016

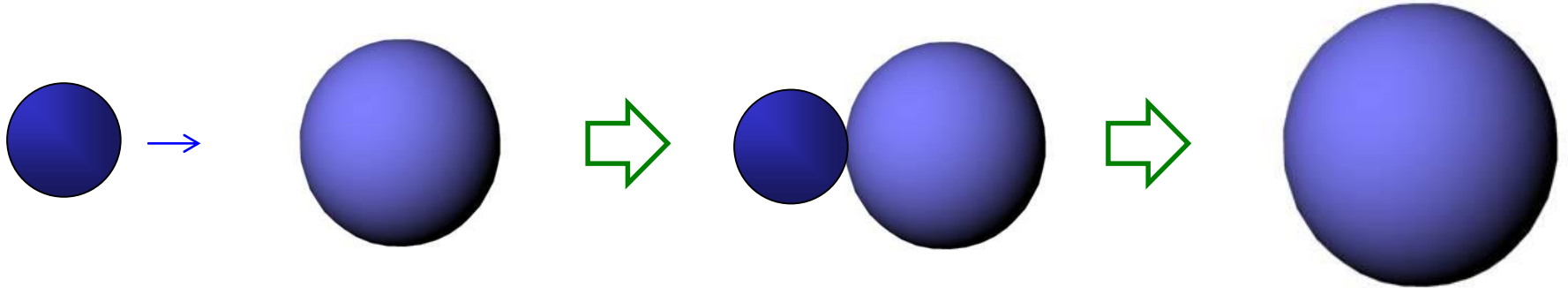
the island of stability (安定的島)



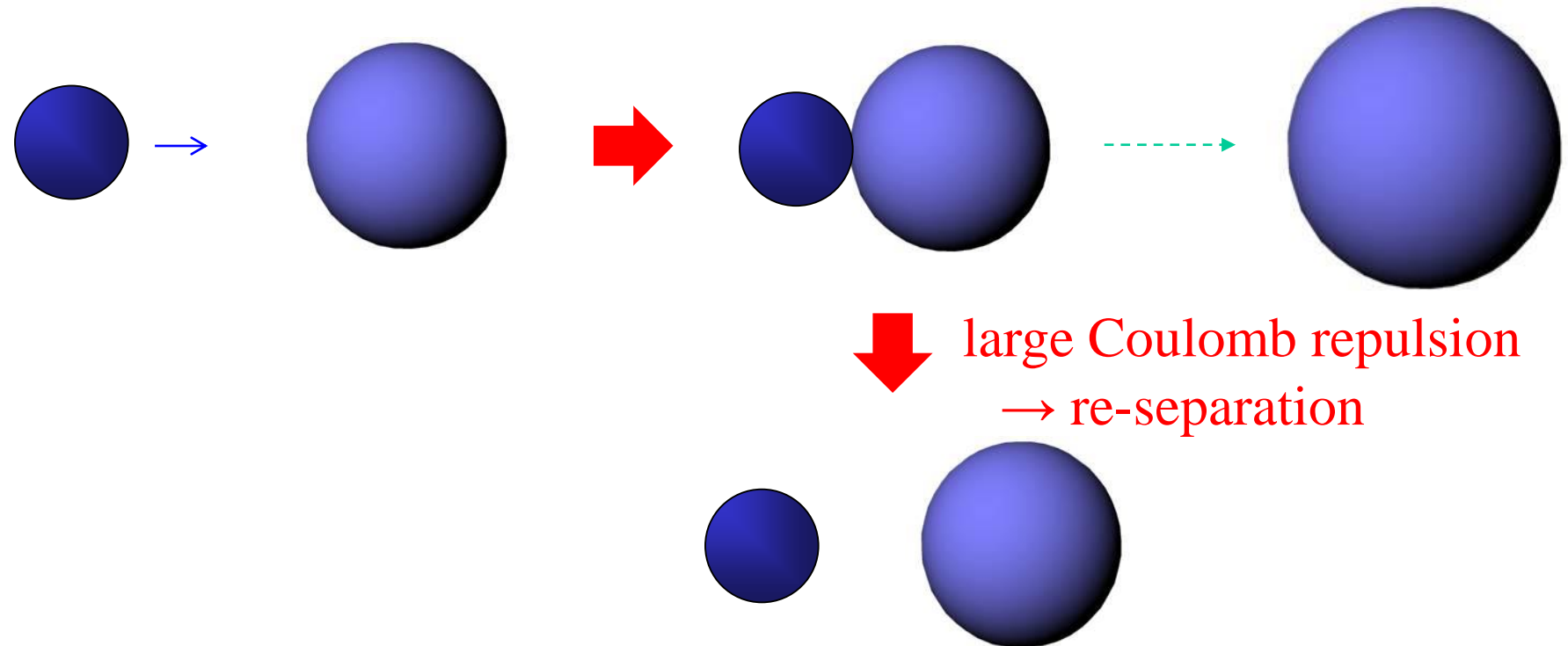
Heavy-ion fusion reaction

A complication

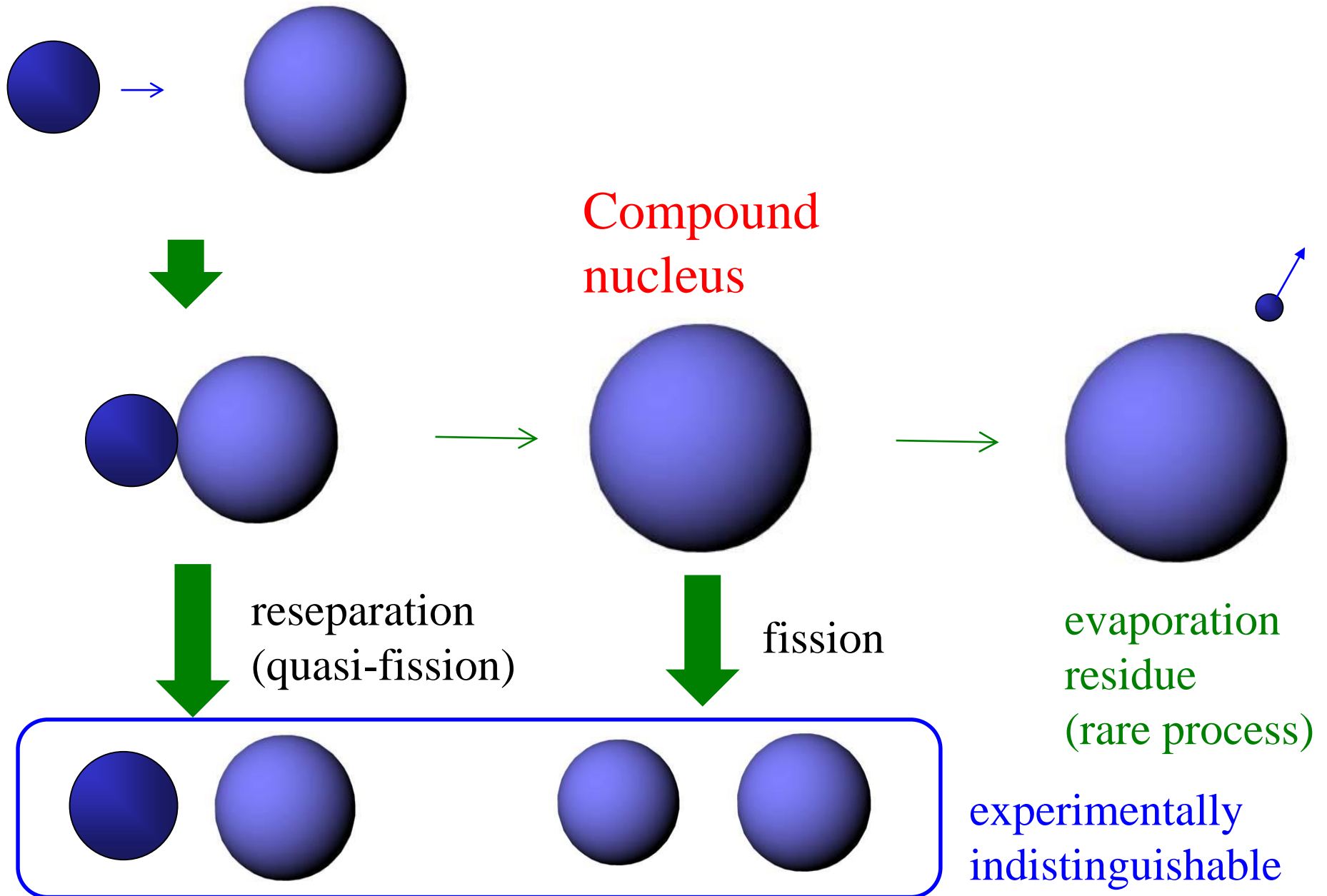
➤ Fusion of medium-heavy systems:



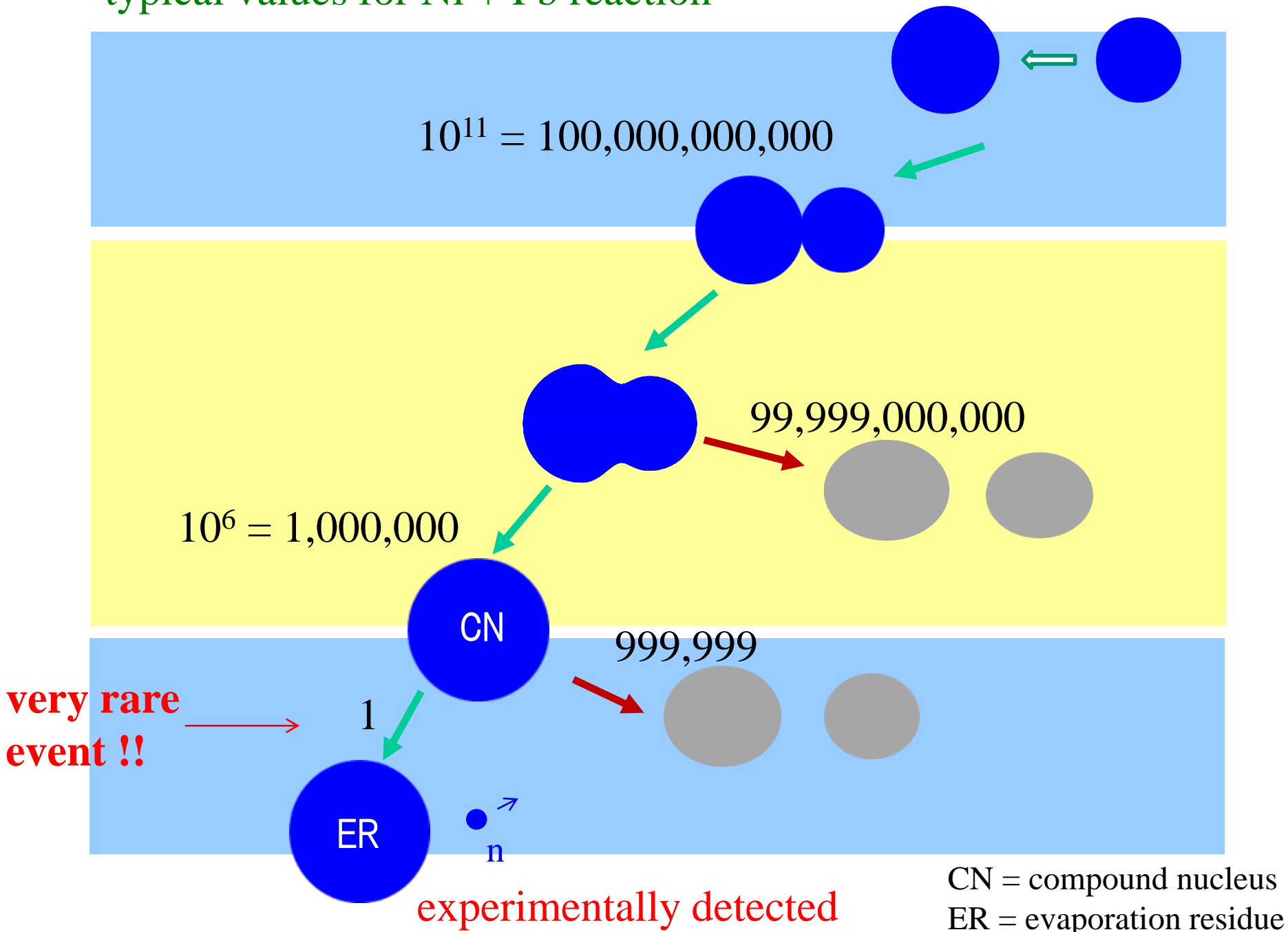
➤ Fusion of heavy and super-heavy systems:



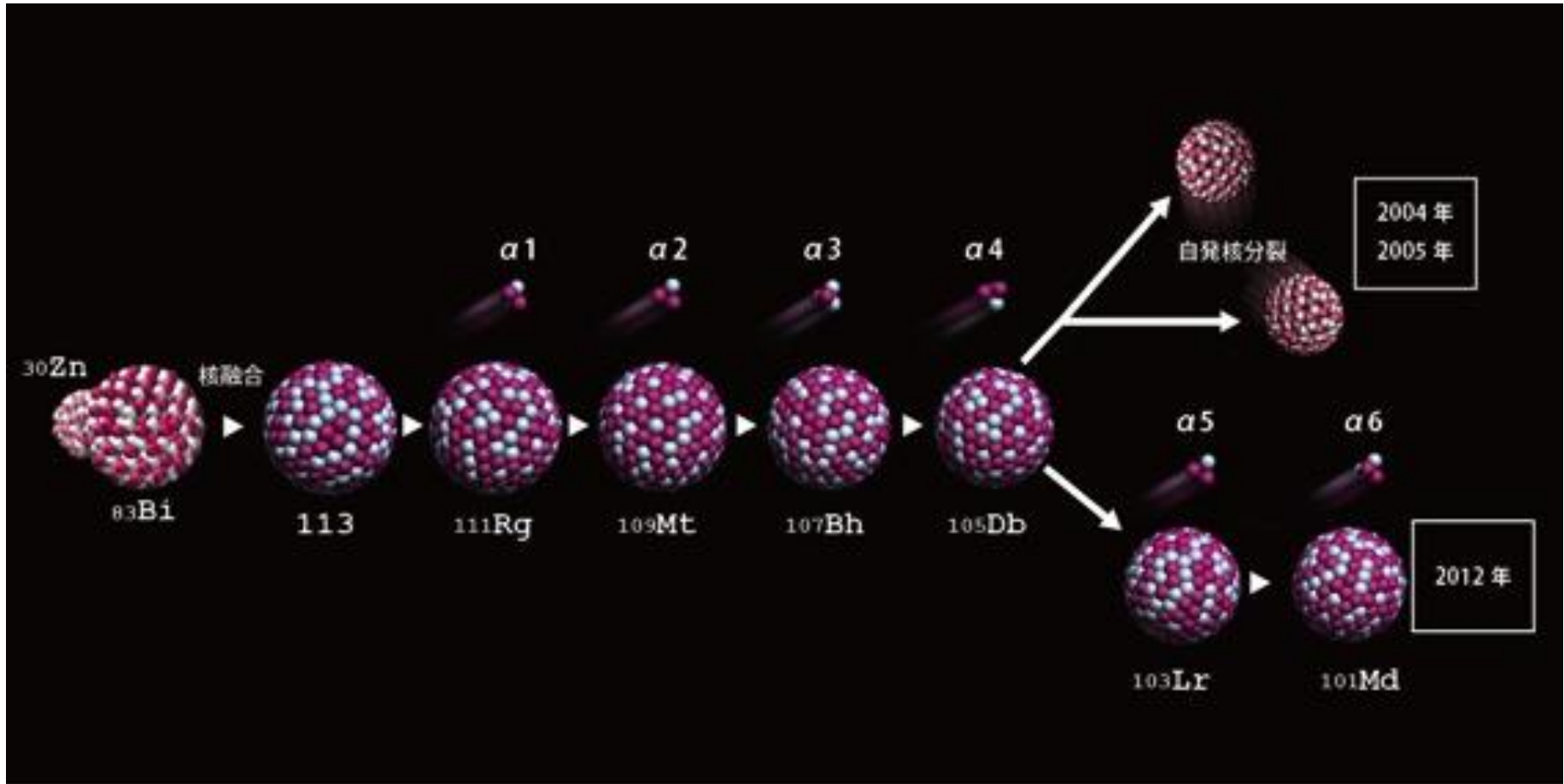
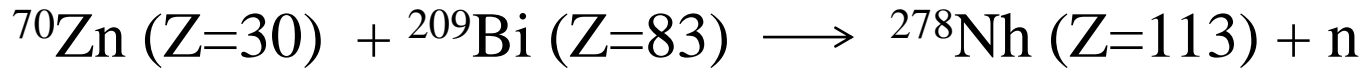
Heavy-ion fusion reactions for superheavy elements



typical values for Ni + Pb reaction



Element 113 (RIKEN, K. Morita et al.)



K. Morita et al., J. Phys. Soc. Jpn. 81('12)103201

only 3 events for 553 days experiment

Theory: Lagenvin approach

multi-dimensional extension of:

$$m \frac{d^2 q}{dt^2} = - \frac{dV(q)}{dq} - \gamma \frac{dq}{dt} + R(t)$$

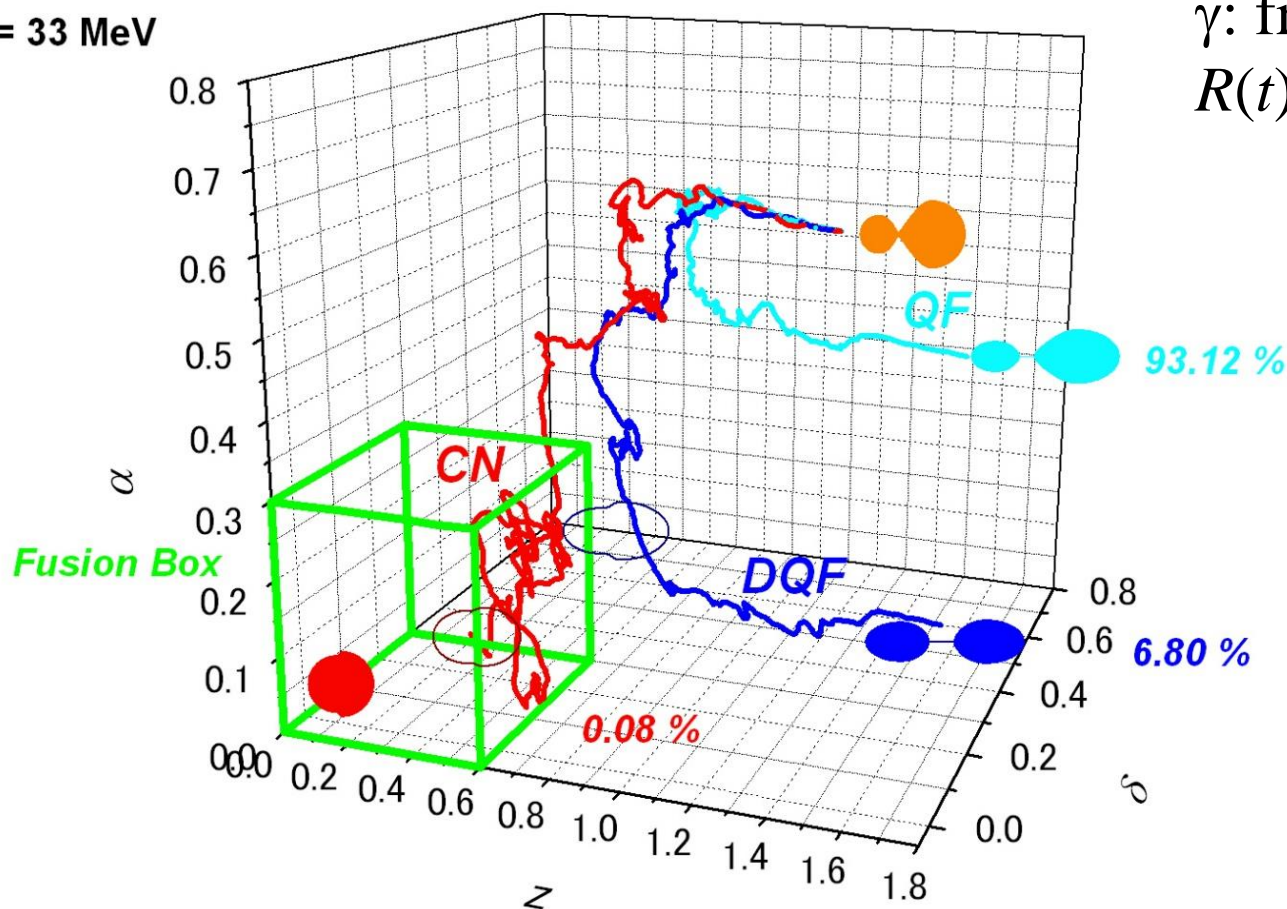


$E^* = 33 \text{ MeV}$

γ : friction coefficient

$R(t)$: random force

摩擦
乱雑力



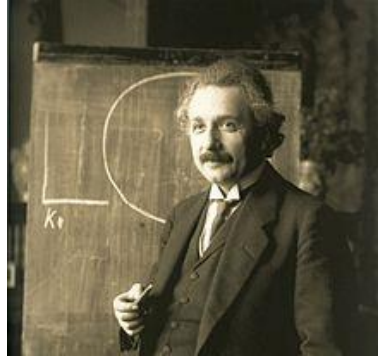
Chemistry of superheavy elements

Group →	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18
↓ Period																		
1	1 H																	2 He
2	3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne
3	11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
4	19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
5	37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
6	55 Cs	56 Ba	57 La	* 72 Hf	* 73 Ta	* 74 W	* 75 Re	* 76 Os	* 77 Ir	* 78 Pt	* 79 Au	* 80 Hg	* 81 Tl	* 82 Pb	* 83 Bi	* 84 Po	* 85 At	* 86 Rn
7	87 Fr	88 Ra	89 Ac	* 104 Rf	* 105 Db	* 106 Sg	* 107 Bh	* 108 Hs	* 109 Mt	* 110 Ds	* 111 Rg	* 112 Cn	* 113 Nh	* 114 Fl	* 115 Mc	* 116 Lv	* 117 Ts	* 118 Og
				* 58 Ce	* 59 Pr	* 60 Nd	* 61 Pm	* 62 Sm	* 63 Eu	* 64 Gd	* 65 Tb	* 66 Dy	* 67 Ho	* 68 Er	* 69 Tm	* 70 Yb	* 71 Lu	
				* 90 Th	* 91 Pa	* 92 U	* 93 Np	* 94 Pu	* 95 Am	* 96 Cm	* 97 Bk	* 98 Cf	* 99 Es	* 100 Fm	* 101 Md	* 102 No	* 103 Lr	

- Are they here in the periodic table?
- Does Nh show the same chemical properties as B, Al, Ga, In, and Tl?

relativistic effect : important for large Z

$$E = mc^2$$



Solution of the Dirac equation (relativistic quantum mechanics)
for a hydrogen-like atom:

$$E_{1S} = mc^2 \sqrt{1 - (Z\alpha)^2} \sim mc^2 \left(1 - \frac{(Z\alpha)^2}{2} - \underbrace{\frac{(Z\alpha)^4}{8} + \dots}_{\text{relativistic effect}} \right)$$

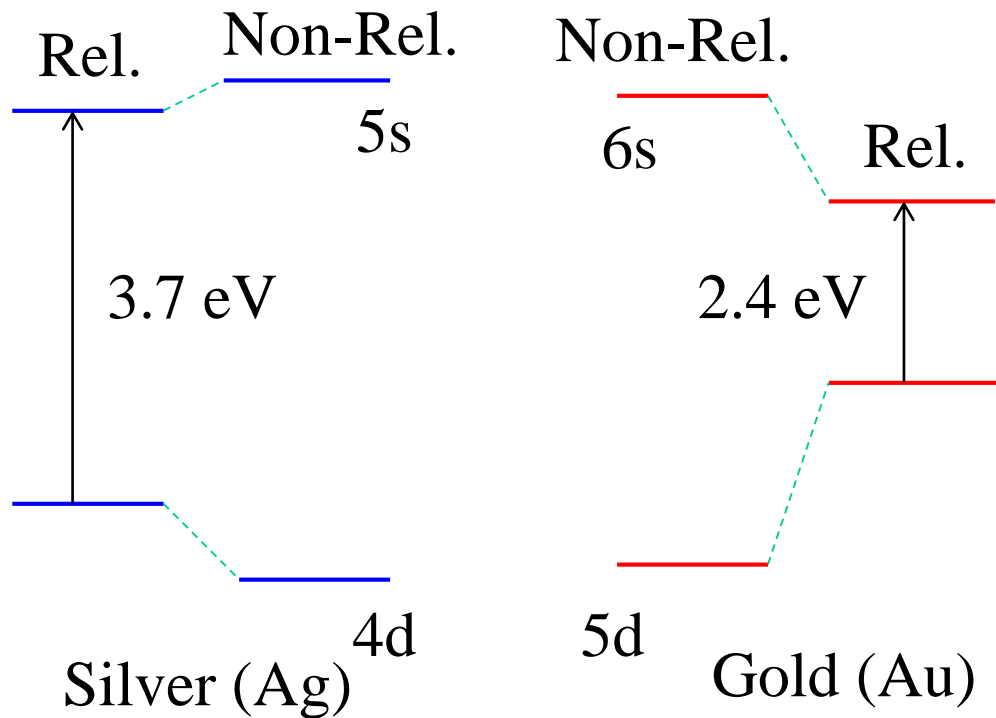
relativistic effect

Famous example of relativistic effects: the color of gold

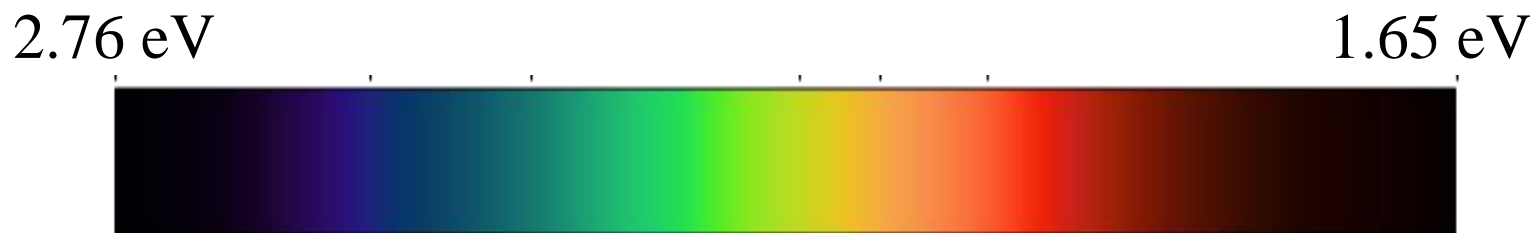
1	1 H																				2 He
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7	87 Fr	88 Ra		104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Ds	111 Rg	112 Cn	113 Uut	114 Fl	115 Uup	116 Lv	117 Uus	118 Uuo			



Gold looked like silver if there was no relativistic effects!



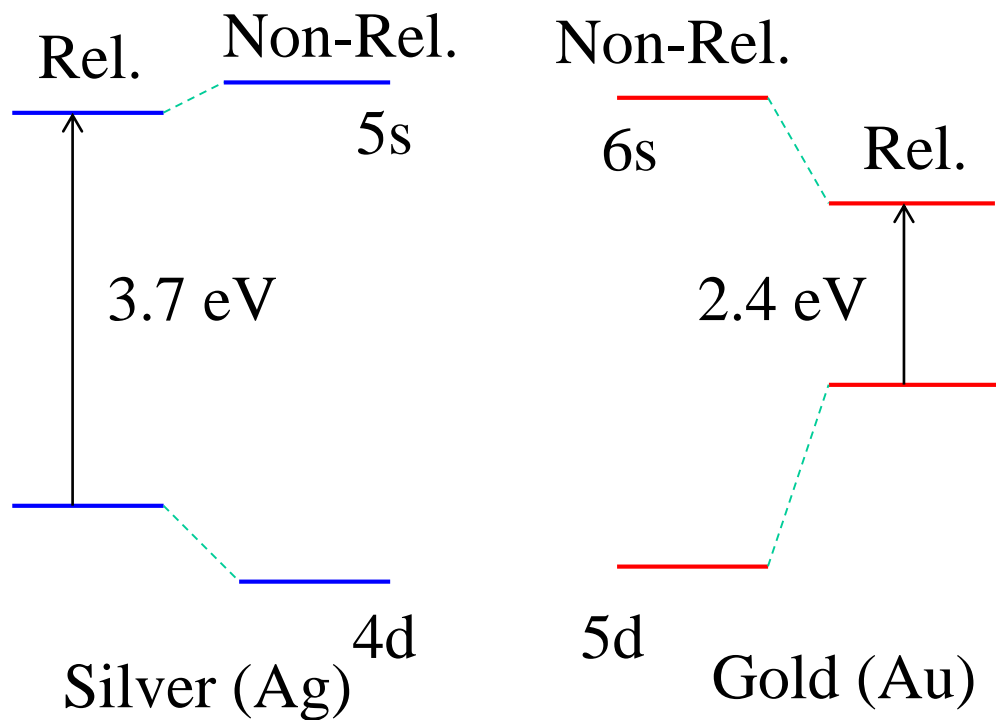
cf. visible spectrum



↑
 3.7 eV

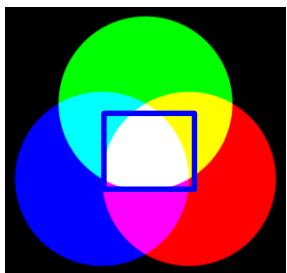


reflected (Ag)

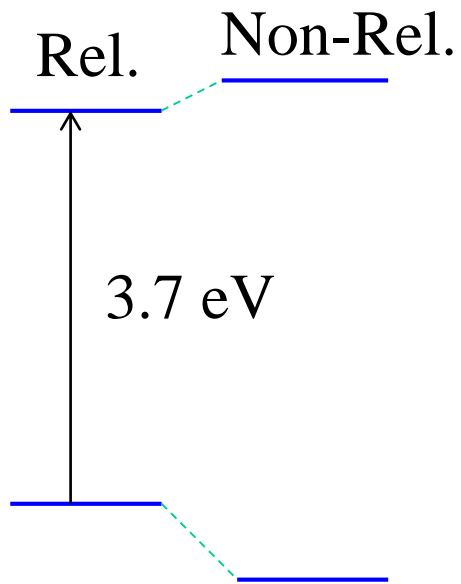


cf. visible spectrum



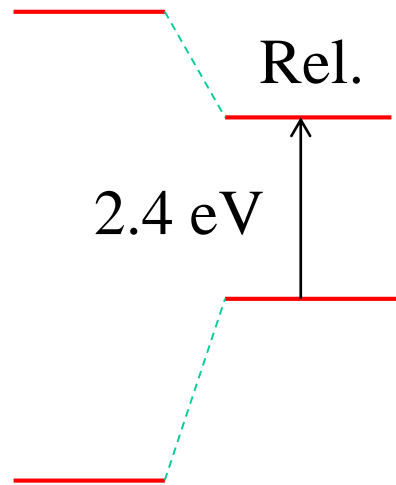


no color
absorbed

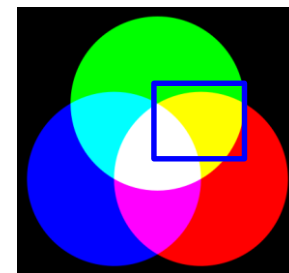


Silver (Ag)

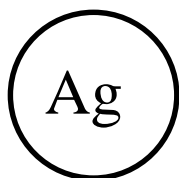
Non-Rel.



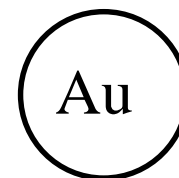
Gold (Au)



blue: absorbed



47th element



79th element

Chemistry of superheavy elements

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1	1 H																		2 He
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	Lanthanides	57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu			
	Actinides	89 Ac	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr			

How do the relativistic effects alter the periodic table for SHE?
 What is the color of superheavy elements?

→ big open questions

Magic Numbers of Atoms and Nuclei : quantum mechanics of many-Fermion systems

Kouichi Hagino

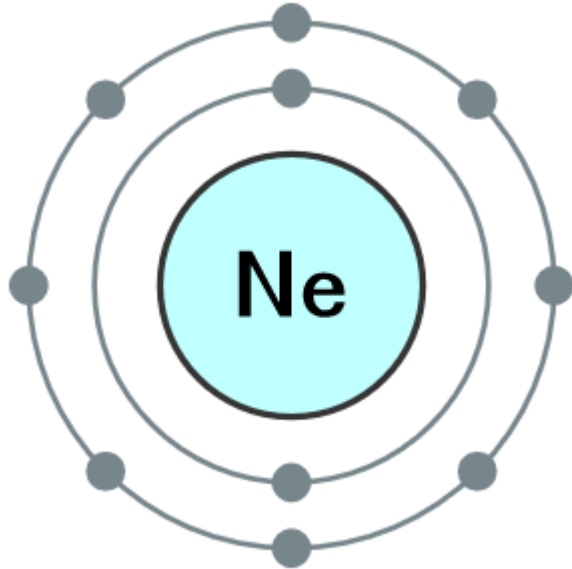
Tohoku University, Sendai, Japan



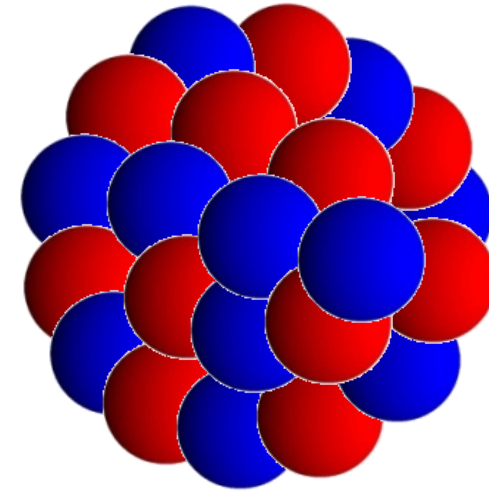
TOHOKU
UNIVERSITY

1. Identical particles: Fermions and Bosons
2. Simple examples: systems with two identical particles
3. Pauli principle
4. Magic numbers

Introduction



atom = nucleus
+ many electrons

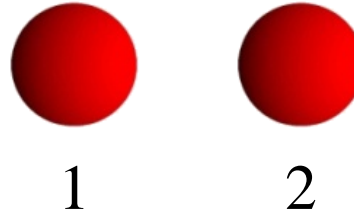


nucleus = many protons
+ many neutrons

Quantum mechanics for those many **Fermion** systems?

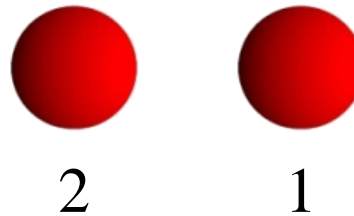
Exchange operator: Fermions and Bosons

a two-particle system



$$H(1, 2) = \frac{p_1^2}{2m} + \frac{p_2^2}{2m} + V(\mathbf{r}_1, \mathbf{r}_2)$$

two particles are identical: particle 1 and 2 cannot be distinguished



$$H(2, 1) = \frac{p_2^2}{2m} + \frac{p_1^2}{2m} + V(\mathbf{r}_2, \mathbf{r}_1) = H(1, 2)$$

Exchange operator: Fermions and Bosons

$$[H, P_{12}] = 0$$

where $P_{12}\Psi(1, 2) = \Psi(2, 1)$ exchange operator

 wave functions have to be simultaneous eigen-states of H and P_{12}

Exchange operator: Fermions and Bosons

$$[H, P_{12}] = 0$$

where $P_{12}\Psi(1, 2) = \Psi(2, 1)$ exchange operator

→ wave functions have to be simultaneous eigen-states of H and P_{12}

Eigen-values of P_{12}

$$P_{12}\Psi(1, 2) = \Psi(2, 1)$$

→ $(P_{12})^2\Psi(1, 2) = P_{12}\Psi(2, 1) = \Psi(1, 2)$

→ $(P_{12})^2 = 1$

→ $P_{12} = \pm 1$

Exchange operator: Fermions and Bosons

$$P_{12} = \pm 1$$

Natural Laws: each particle has a definite value of P_{12}
(independent of e.g., experimental setup and temperature)

◆ particles with a half-integer spin: $P_{12} = -1$ (“Fermion”)

electrons, protons, neutrons,.....

$$\psi^{(-)}(1, 2) = \frac{1}{\sqrt{2}}[\psi(1, 2) - \psi(2, 1)]$$

◆ particles with an integer spin: $P_{12} = +1$ (“Boson”)

photons, pi mesons,.....

$$\psi^{(+)}(1, 2) = \frac{1}{\sqrt{2}}[\psi(1, 2) + \psi(2, 1)]$$

Simple examples: systems with two identical particles

Assume a spin-independent Hamiltonian for a two-particle system:

$$H = \frac{p_1^2}{2m} + \frac{p_2^2}{2m} + V(\mathbf{r}_1, \mathbf{r}_2)$$

→ separable between the space and the spin

$$\Psi(x_1, x_2) = \Psi_{\text{space}}(\mathbf{r}_1, \mathbf{r}_2) \cdot \Psi_{\text{spin}}$$

Simple examples: systems with two identical particles

Assume a spin-independent Hamiltonian for a two-particle system:

$$H = \frac{p_1^2}{2m} + \frac{p_2^2}{2m} + V(\mathbf{r}_1, \mathbf{r}_2)$$

→ separable between the space and the spin

$$\Psi(x_1, x_2) = \Psi_{\text{space}}(\mathbf{r}_1, \mathbf{r}_2) \cdot \Psi_{\text{spin}}$$

◆ spin-zero bosons

no spin → symmetrize the spatial part

$$\begin{aligned}\Psi(\mathbf{r}_1, \mathbf{r}_2) &= \Psi_{\text{space}}(\mathbf{r}_1, \mathbf{r}_2) \\ &= \frac{1}{\sqrt{2}}[\phi(\mathbf{r}_1, \mathbf{r}_2) + \phi(\mathbf{r}_2, \mathbf{r}_1)]\end{aligned}$$

Simple examples: systems with two identical particles

$$\Psi(x_1, x_2) = \Psi_{\text{space}}(\mathbf{r}_1, \mathbf{r}_2) \cdot \Psi_{\text{spin}}$$

◆ spin-1/2 Fermions $\Psi(x_1, x_2) = -\Psi(x_2, x_1)$

Spin part:

$$|S = 1\rangle = |\uparrow\uparrow\rangle, \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle), |\downarrow\downarrow\rangle \quad \text{symmetric}$$

$$|S = 0\rangle = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \quad \text{anti-symmetric}$$

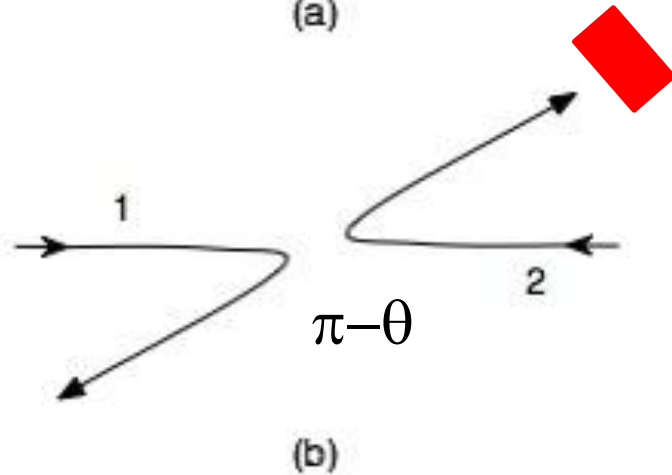
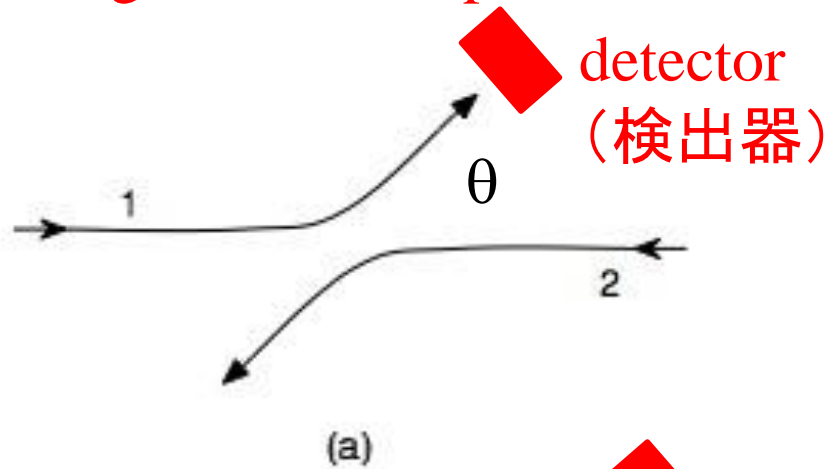
→ spatial part: anti-symmetric for $S = 1$
symmetric for $S = 0$



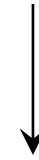
$$\Psi_{S=0}(x_1, x_2) = \frac{1}{\sqrt{2}}[\phi(\mathbf{r}_1, \mathbf{r}_2) + \phi(\mathbf{r}_2, \mathbf{r}_1)]|S = 0\rangle$$

$$\Psi_{S=1}(x_1, x_2) = \frac{1}{\sqrt{2}}[\phi(\mathbf{r}_1, \mathbf{r}_2) - \phi(\mathbf{r}_2, \mathbf{r}_1)]|S = 1\rangle$$

Scattering of identical particles



these two processes cannot be distinguished



add two amplitudes and then take square

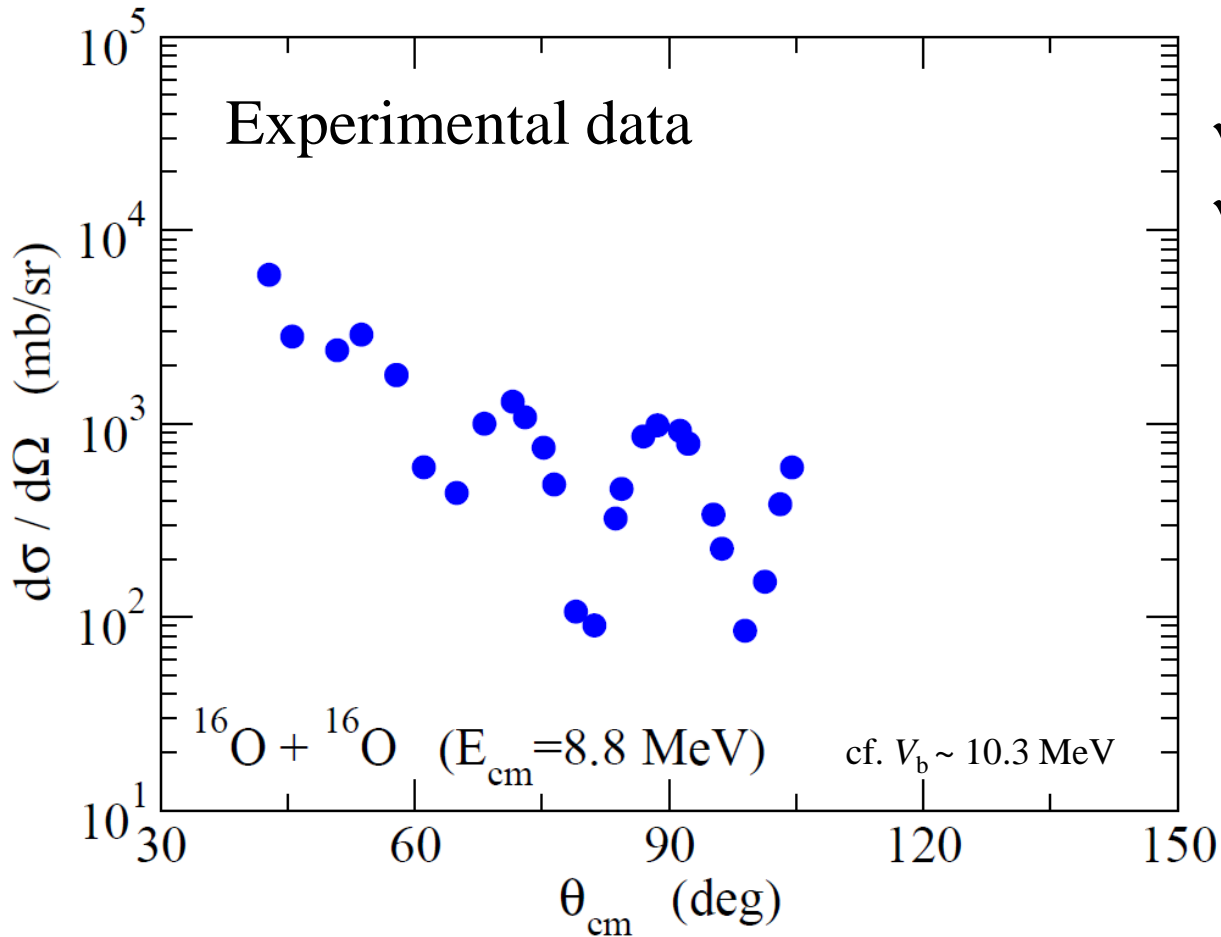
→ interference

干涉

$$\begin{aligned}\frac{d\sigma}{d\Omega} &= |f(\theta) \pm f(\pi - \theta)|^2 \\ &= |f(\theta)|^2 + |f(\pi - \theta)|^2 \pm f^*(\theta)f(\pi - \theta) \pm f(\theta)f^*(\pi - \theta)\end{aligned}$$

+ : for spatially symmetric, and - : for spatially anti-symmetric

$^{16}\text{O} + ^{16}\text{O}$ elastic scattering

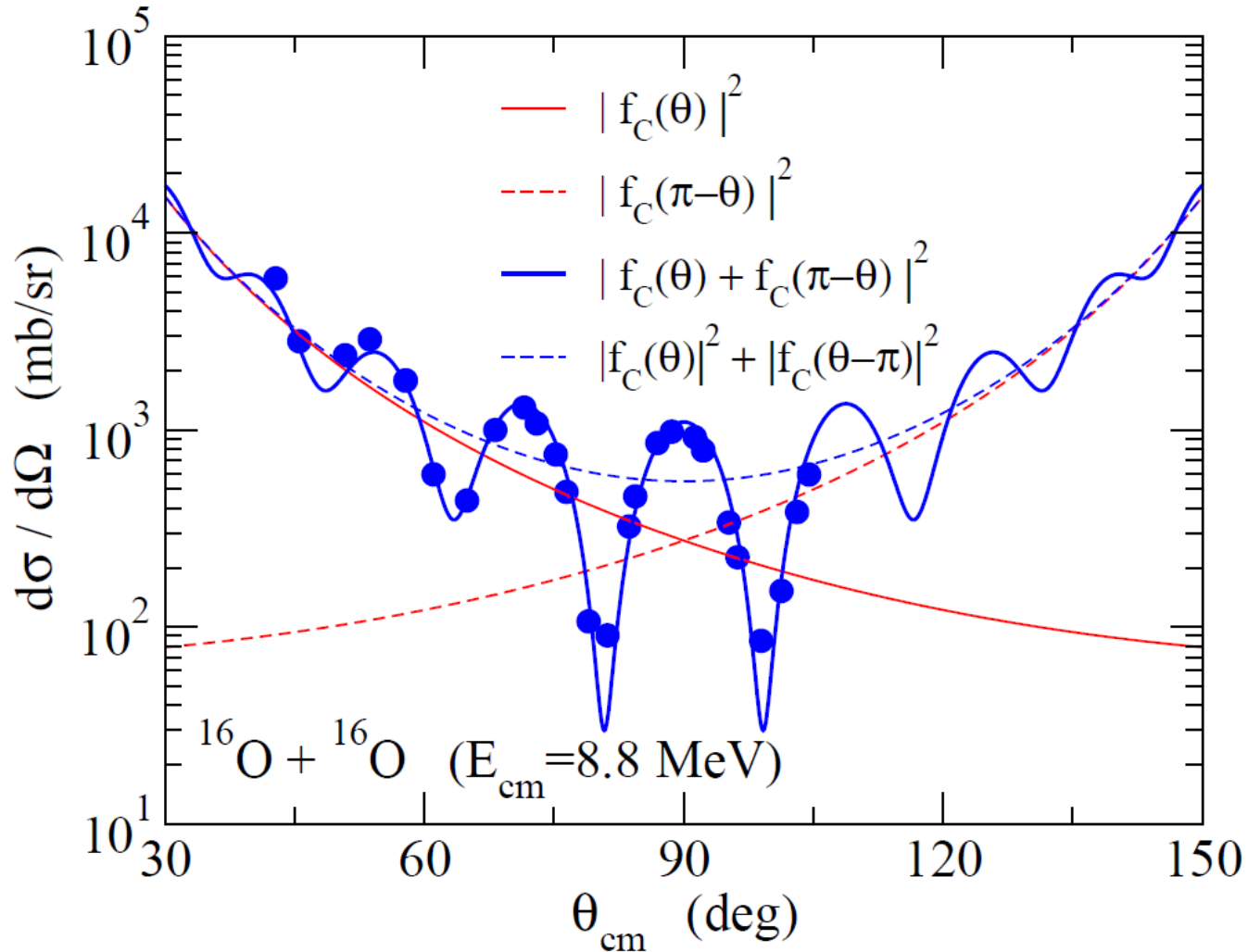


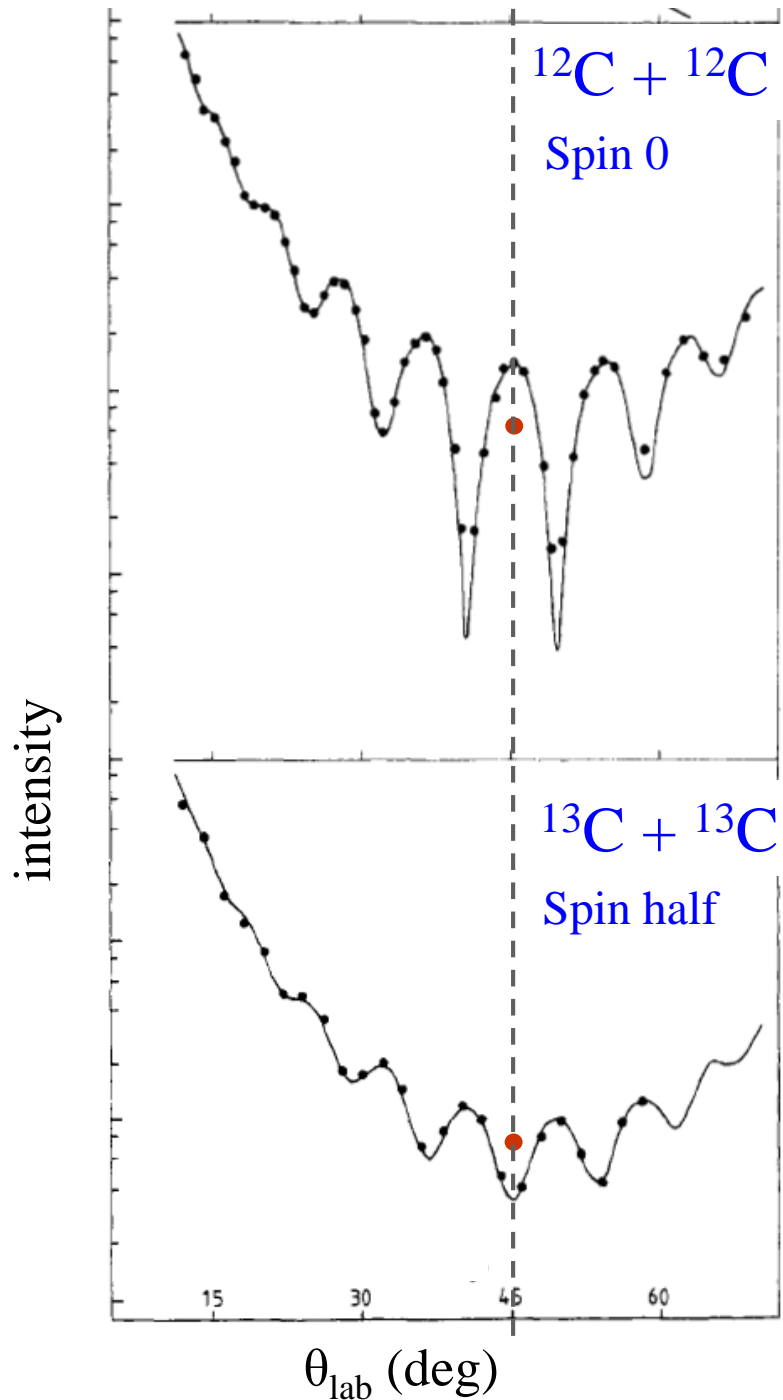
- ✓ Symmetric at 90 deg.
- ✓ clear oscillations

D.A. Bromley et al.,
Phys. Rev. 123('61)878

$^{16}\text{O} + ^{16}\text{O}$ elastic scattering

$$^{16}\text{O}: \text{spin-zero Boson} \rightarrow \frac{d\sigma}{d\Omega} = |f(\theta) + f(\pi - \theta)|^2$$





identical Bosons

→ constructive interference

$$\frac{d\sigma}{d\Omega} = |f(\theta) + f(\pi - \theta)|^2$$

identical Fermions

→ destructive interference

$$\frac{d\sigma}{d\Omega} = \frac{3}{4} |f(\theta) - f(\pi - \theta)|^2 + \frac{1}{4} |f(\theta) + f(\pi - \theta)|^2$$

S=1

S=0

Pauli exclusion principle and Slater determinants

Pauli exclusion principle: two identical Fermion cannot take the same state

Let us assume:

$$H(1, 2) = \underbrace{\frac{p_1^2}{2m} + V(r_1)}_{\equiv h_1} + \underbrace{\frac{p_2^2}{2m} + V(r_2)}_{\equiv h_2}$$

(no interaction between 1 and 2)

separation of variables \rightarrow a product form of wave function

$$\left(\frac{p^2}{2m} + V(r) \right) \phi_n(x) = \epsilon_n \phi_n(x); \quad x \equiv (r, \sigma)$$

$$\longrightarrow \psi^{(-)}(x_1, x_2) = \frac{1}{\sqrt{2}} [\phi_n(x_1) \phi_{n'}(x_2) - \phi_n(x_2) \phi_{n'}(x_1)]$$

Pauli exclusion principle and Slater determinants


Pauli exclusion principle: two identical Fermion cannot take the same state

$$H(1, 2) = \frac{p_1^2}{2m} + V(\mathbf{r}_1) + \frac{p_2^2}{2m} + V(\mathbf{r}_2)$$

separation of variables \rightarrow a product form of wave function

$$\left(\frac{p^2}{2m} + V(\mathbf{r}) \right) \phi_n(x) = \epsilon_n \phi_n(x); \quad x \equiv (\mathbf{r}, \sigma)$$

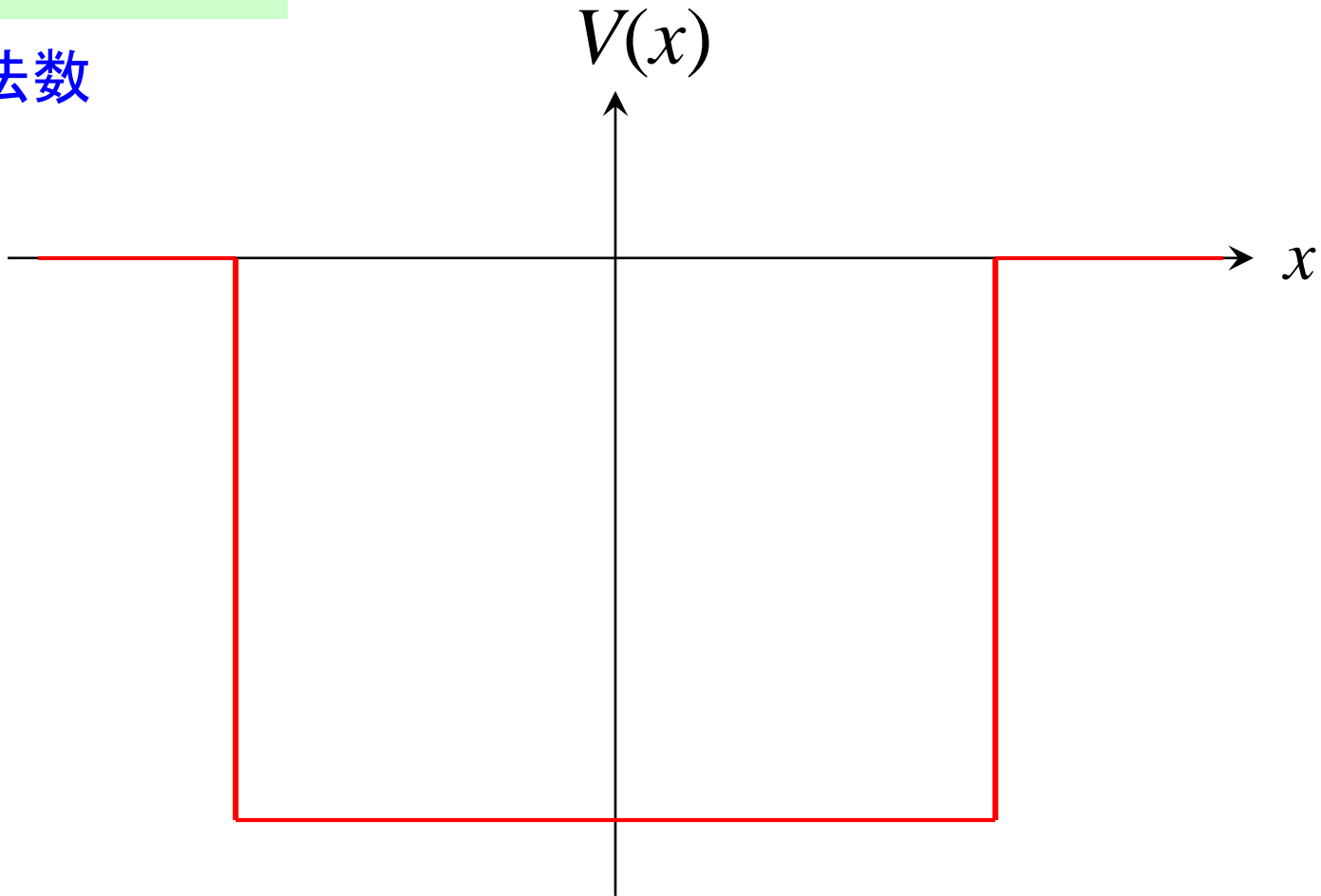
$$\longrightarrow \psi^{(-)}(x_1, x_2) = \frac{1}{\sqrt{2}} [\phi_n(x_1) \phi_{n'}(x_2) - \phi_n(x_2) \phi_{n'}(x_1)]$$


$$\psi^{(-)}(x_1, x_2) = 0 \quad \text{if } n = n'$$

(Pauli principle)

Magic numbers

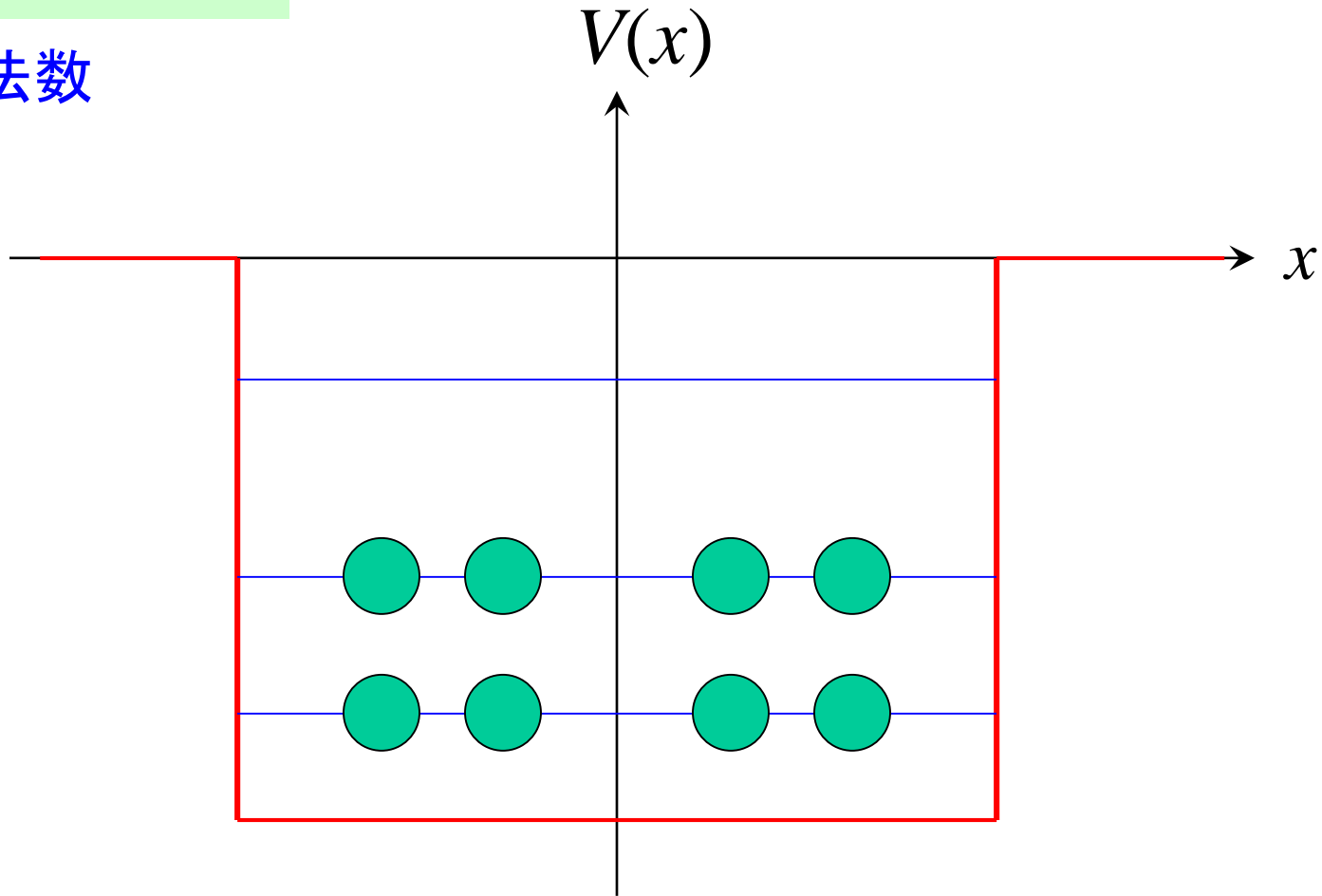
魔法数



$$\left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right) \psi(x) = E\psi(x)$$

Magic numbers

魔法数



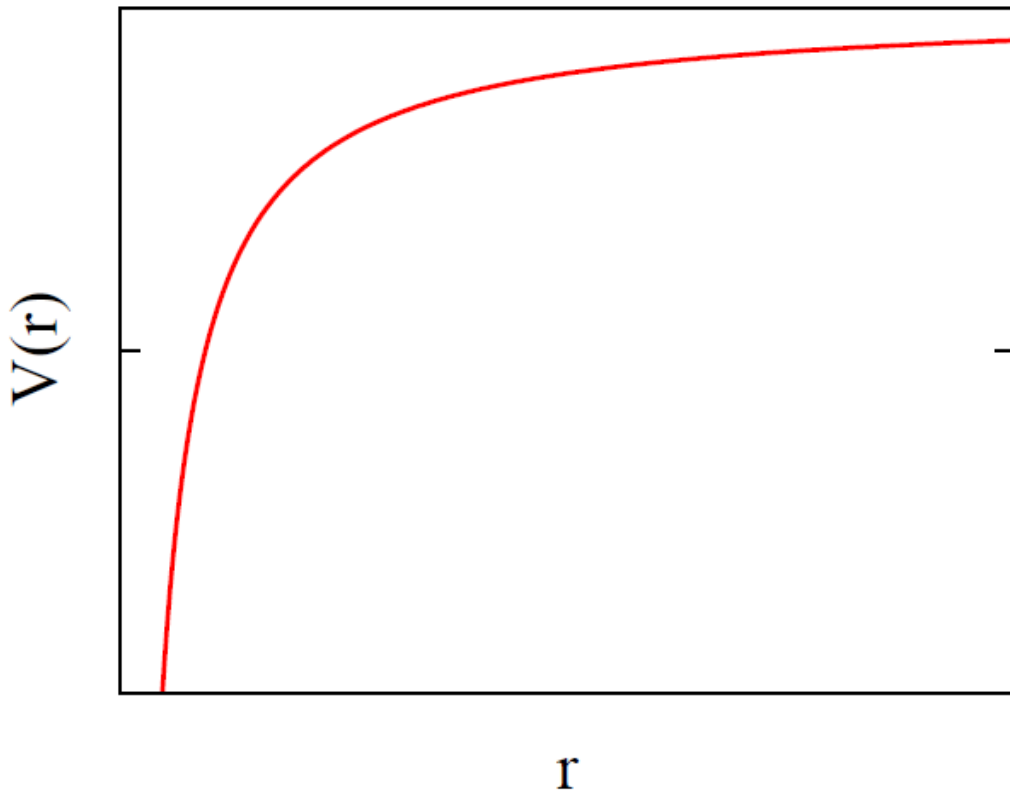
discrete bound states

The lowest state of many-Fermion systems
= put particles from the bottom of the potential well (Pauli principle)

Magic numbers

Hydrogen-like potential:

$$V(r) = -\frac{Ze^2}{r}$$



$$E_n = -\frac{(Z\alpha)^2}{2n^2} mc^2$$

$$\alpha = \frac{e^2}{\hbar c} \sim \frac{1}{137}$$

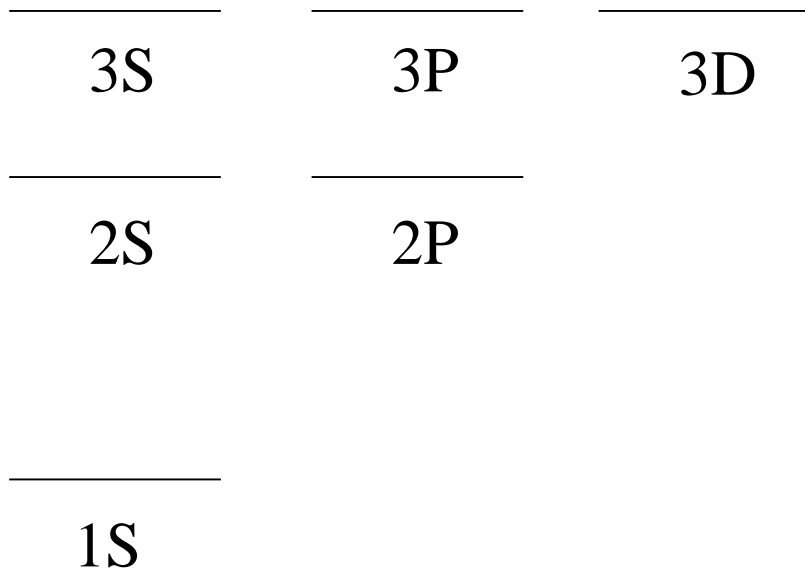
$$n = n_r + l + 1$$

Magic numbers

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Magic numbers

Hydrogen-like potential:

$$V(r) = -\frac{Ze^2}{r}$$

degeneracy = $2 * (2l + 1)$

(spin x l_z)

$$E_n = -\frac{(Z\alpha)^2}{2n^2} mc^2$$

3S [2]

3P [6]

3D [10]

2S [2]

2P [6]

$$\alpha = \frac{e^2}{\hbar c} \sim \frac{1}{137}$$

$$n = n_r + l + 1$$

1S [2]

Magic numbers

Hydrogen-like potential:

$$V(r) = -\frac{Ze^2}{r}$$

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(spin x l_z)

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3S [2]	3P [6]	3D [10]
2S [2]	2P [6]	

$$\alpha = \frac{e^2}{\hbar c} \sim \frac{1}{137}$$

$$n = n_r + l + 1$$



Magic numbers

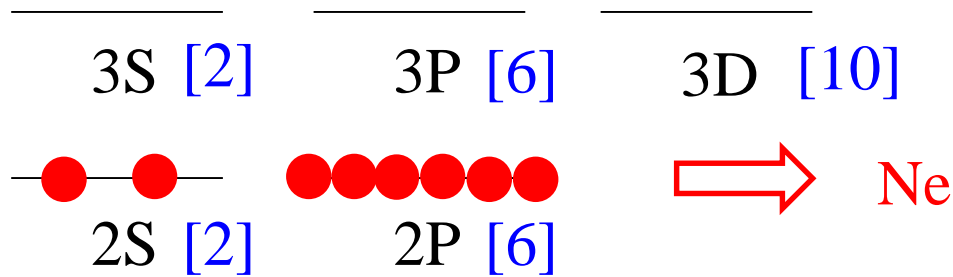
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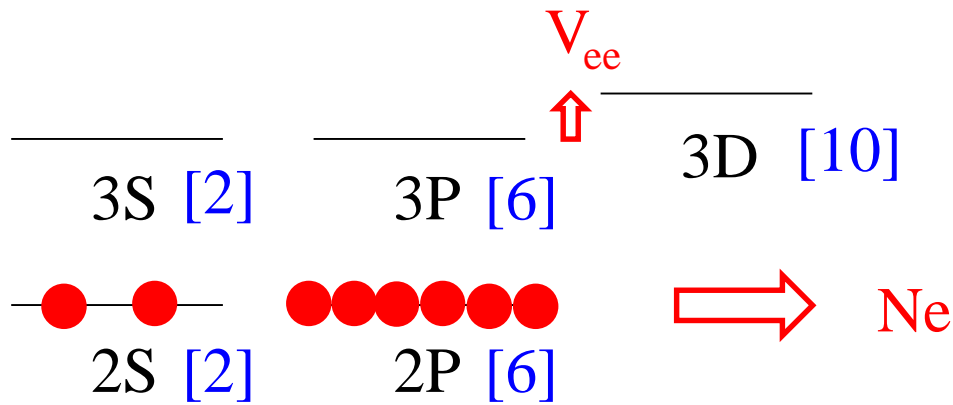
$$n = n_r + l + 1$$

Magic numbers

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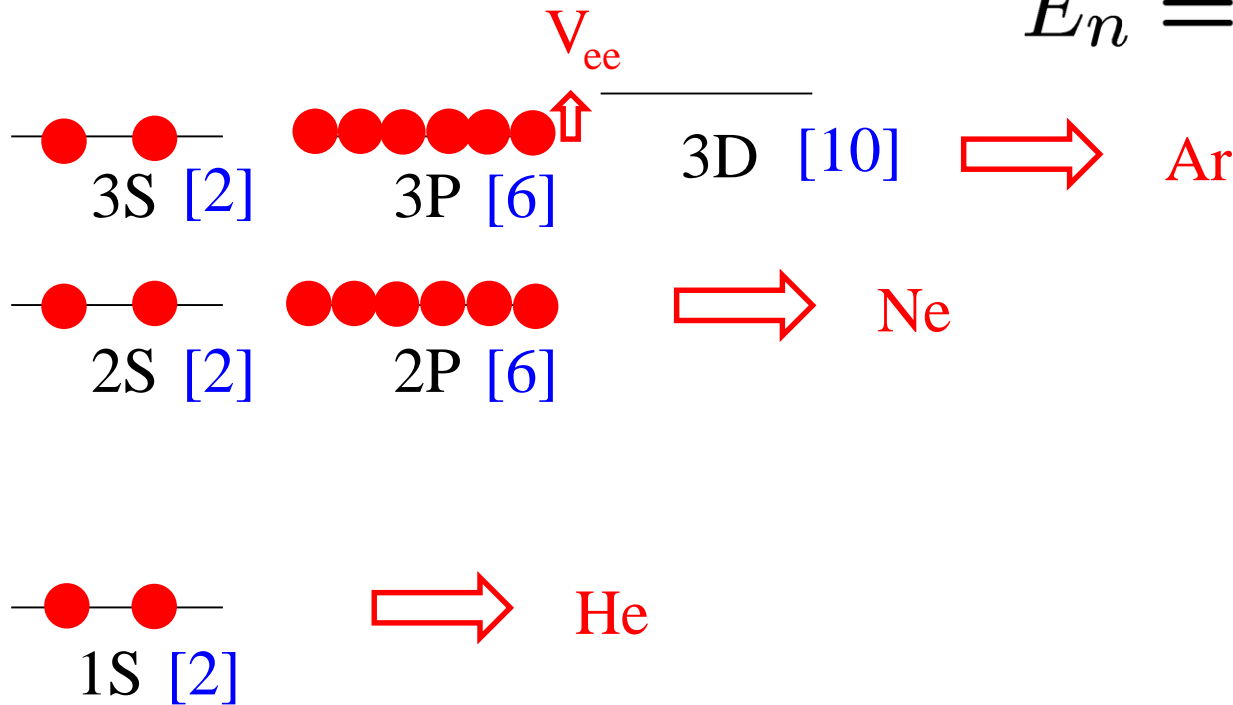
$$n = n_r + l + 1$$

Magic numbers

Hydrogen-like potential:

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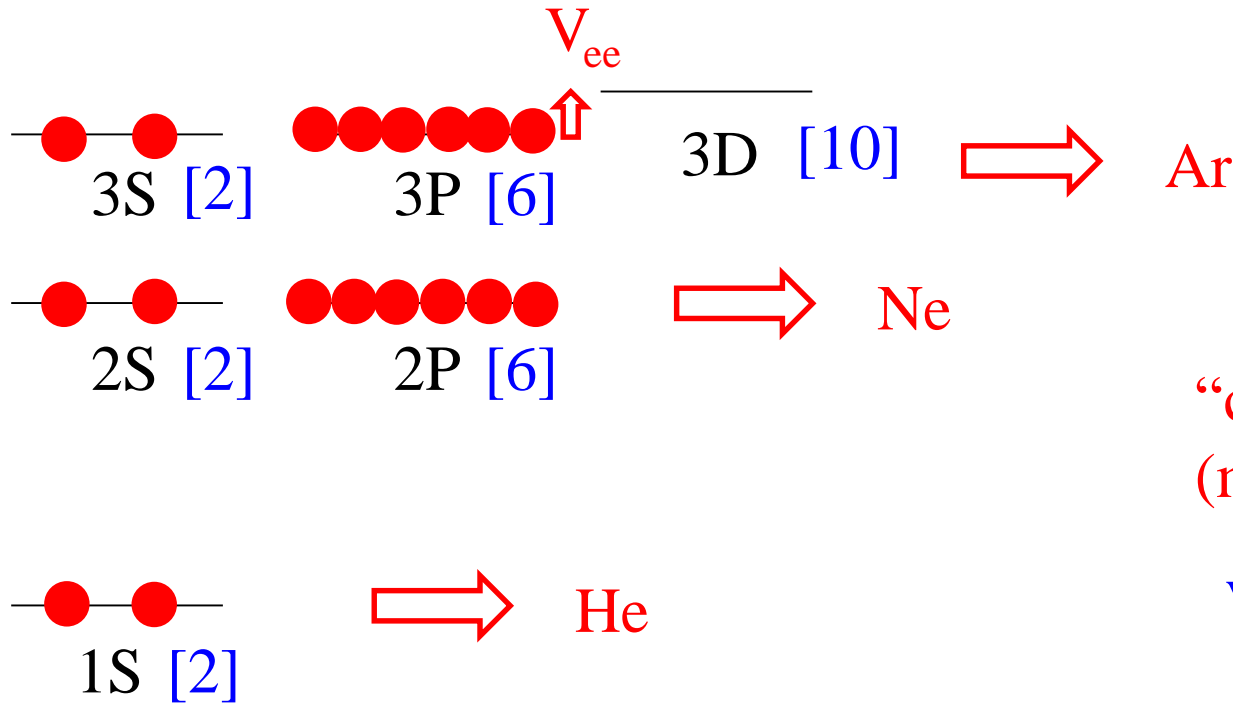
$$n = n_r + l + 1$$

Magic numbers

Hydrogen-like potential:

$$V(r) = -\frac{Ze^2}{r}$$

degeneracy = $2 * (2l + 1)$



“closed shell”
(magic numbers)

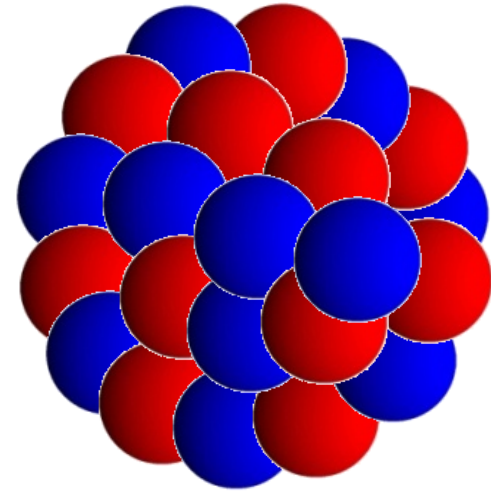
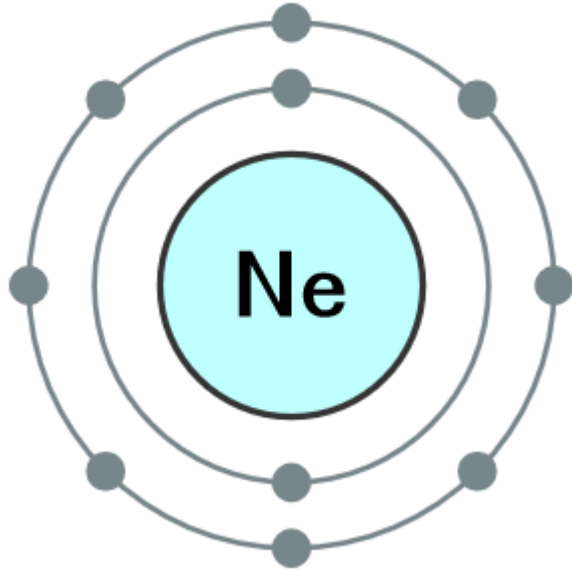
very stable

Periodic Table of elements

noble gas

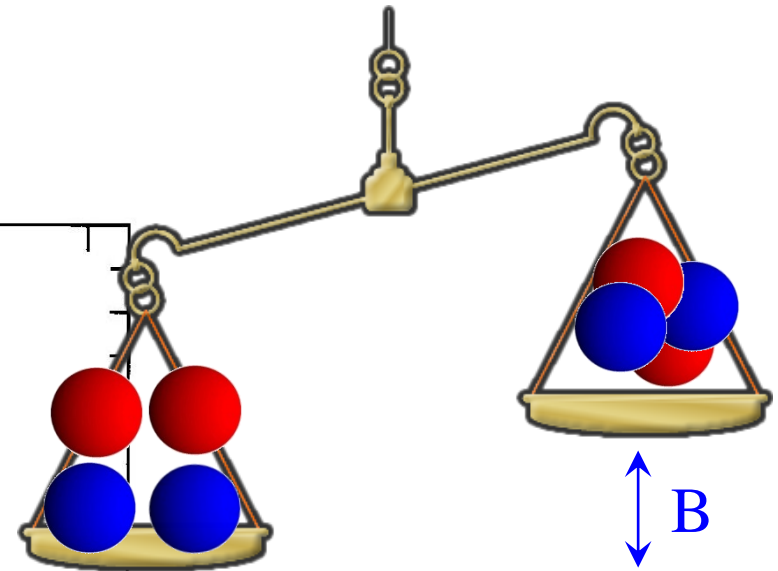
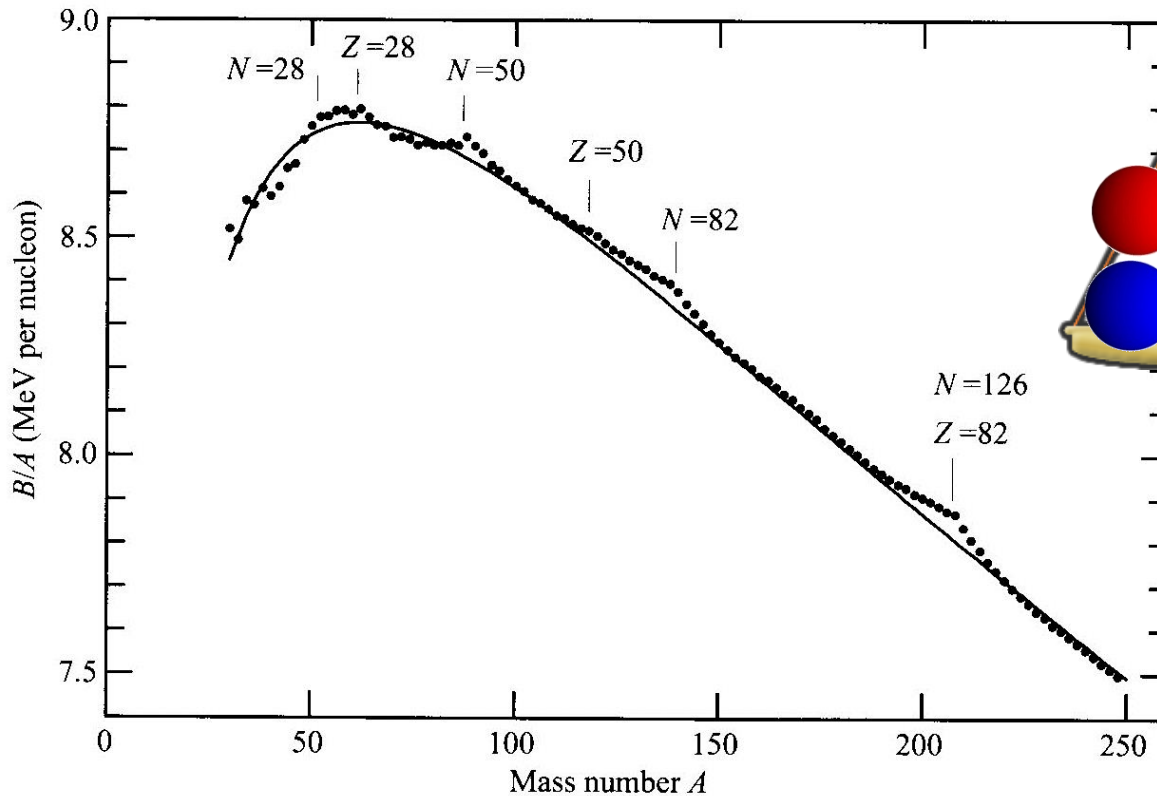
Group →	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	
↓ Period																			
1	1 H																		2 He
2	3 Li	4 Be											5 B	6 C	7 N	8 O	9 F		10 Ne
3	11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl		18 Ar
4	19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br		36 Kr
5	37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I		54 Xe
6	55 Cs	56 Ba	57 La *	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At		86 Rn
7	87 Fr	88 Ra	89 Ac *	104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Ds	111 Rg	112 Cn	113 Nh	114 Fl	115 Mc	116 Lv	117 Ts		118 Og
				* 58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu		
				* 90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr		

Magic numbers



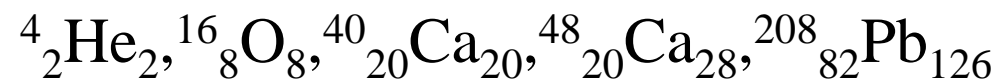
similar magic numbers also in atomic nuclei

Magic numbers

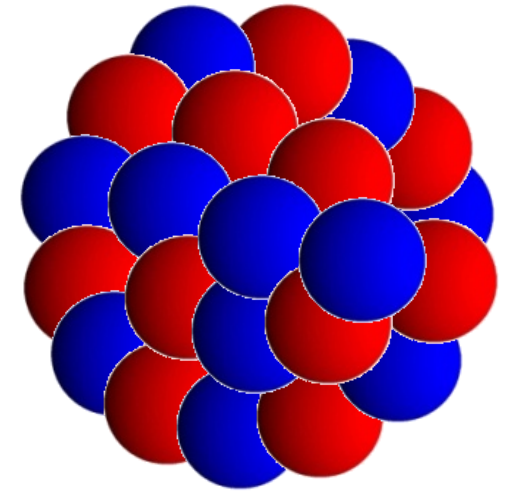


Extra binding for N or $Z = 2, 8, 20, 28, 50, 82, 126$ (magic numbers)

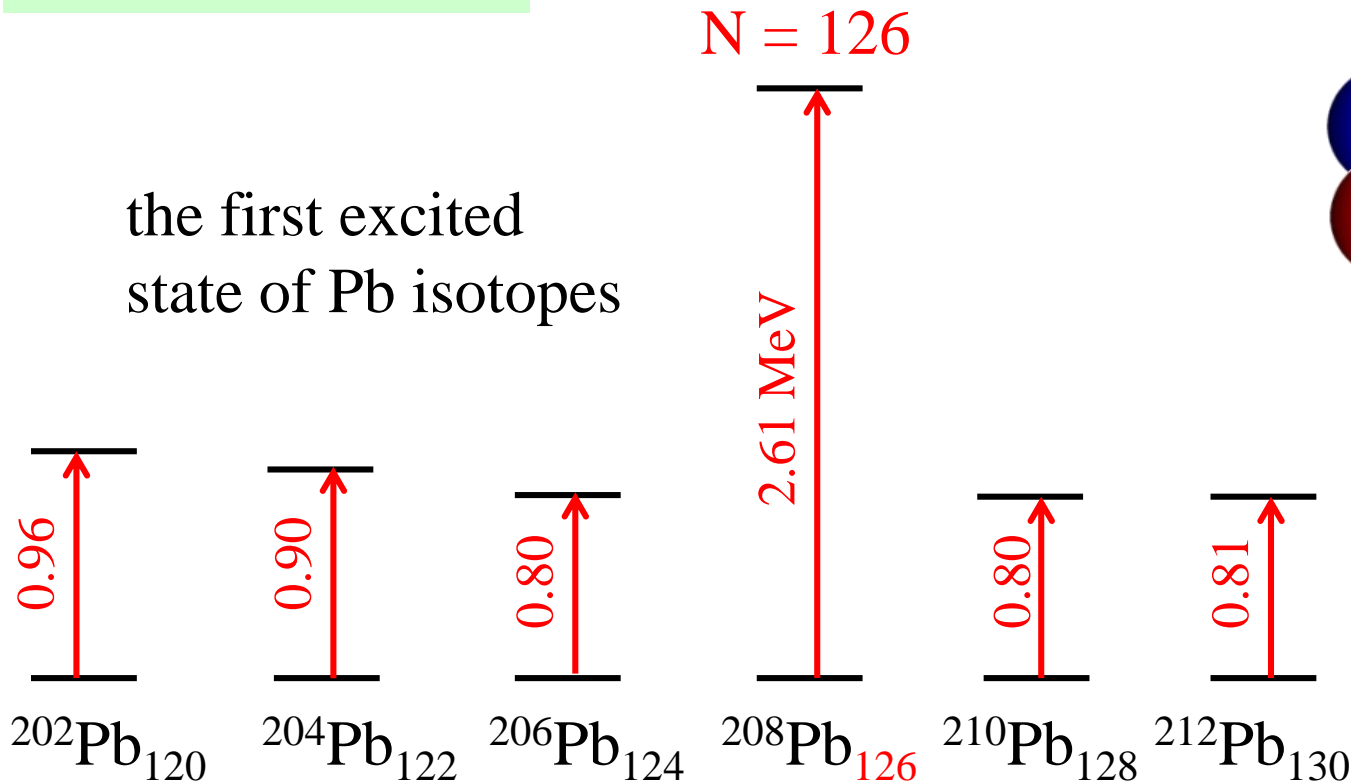
→ Very stable



Magic numbers

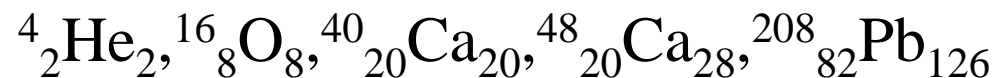


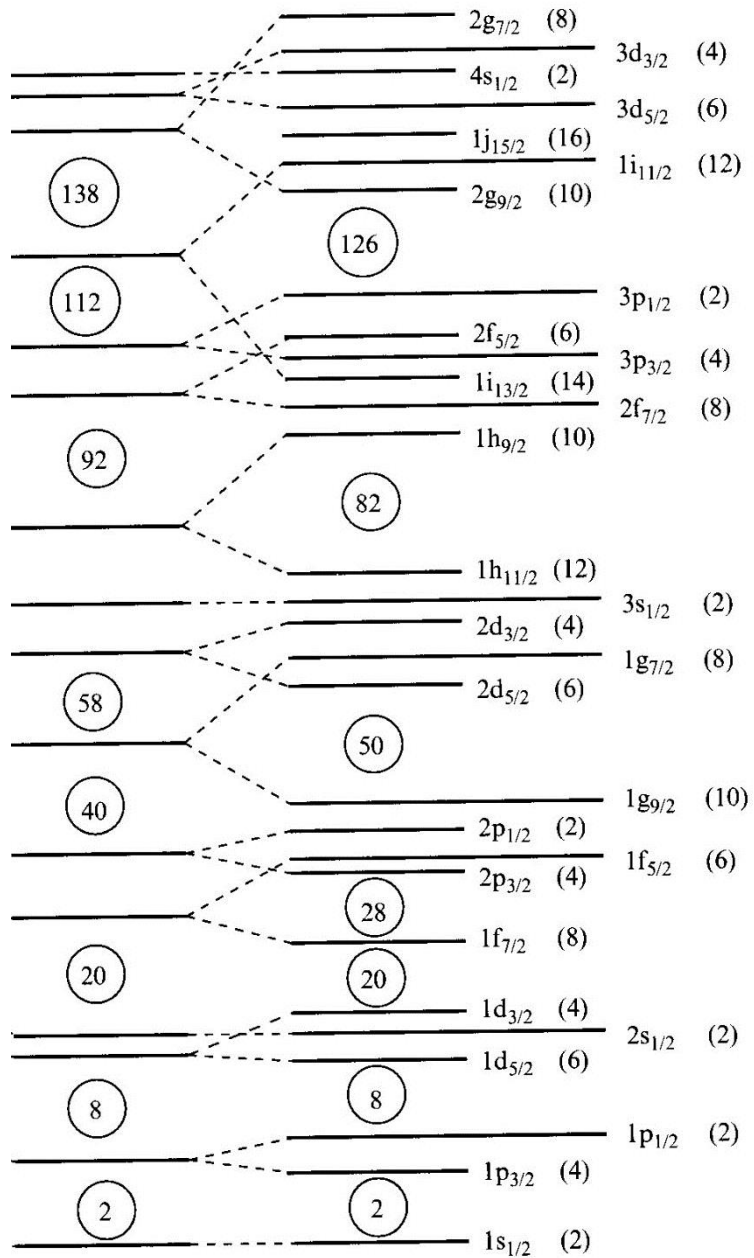
the first excited state of Pb isotopes



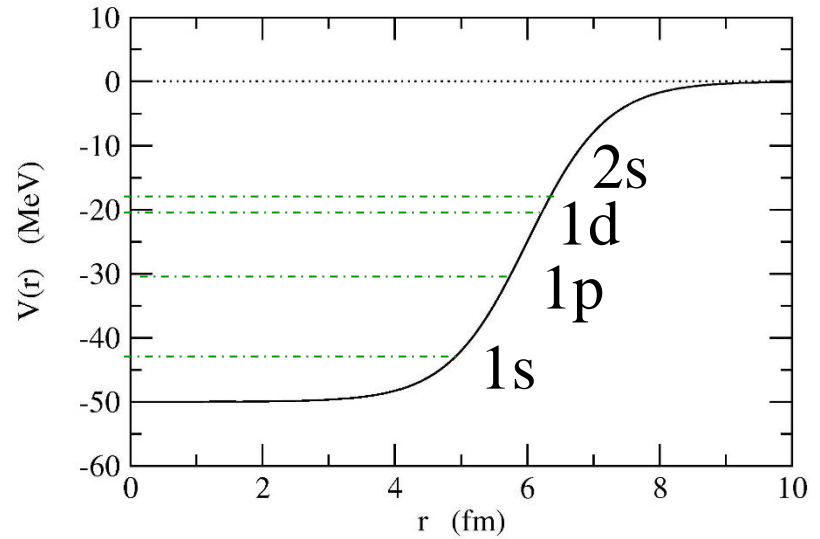
Extra binding for N or $Z = 2, 8, 20, 28, 50, 82, 126$ (magic numbers)

⇒ Very stable





$$V(r) = \frac{-V_0}{1 + \exp((r - R_0)/a)}$$



+ spin-orbit potential

$$V_{ls}(r) \mathbf{l} \cdot \mathbf{s}$$

Woods-Saxon
well

Woods-Saxon plus
spin-orbit coupling

Lucky accident for the origin of life

Atomic magic numbers

electron #: 2, 10, 18, 36, 54, 86

元素の周期表

Double magic

Legend:

- 典型金属元素 (Typical metal element)
- 半金属元素 (Semi-metal element)
- 非金属元素 (Non-metal element)
- 遷移金属元素 (Transition metal element)
- 希ガス (Noble gas)

inert gas : He, Ne, Ar, Kr, Xe, Rn

Nuclear magic numbers

proton # or neutron #

2, 8, 20, 28, 50, 82, 126

→ e.g., $^{16}_8\text{O}_8$ (double magic)

→ many oxygen nuclei:
produced during
nucleosynthesis

→ oxygen: chemically active

→ several complex chemical
reactions, leading to the
birth of life

Summary

Everything is made from atoms.

Nuclear Physics is important for many things.

Nuclei have very rich nature.

Nuclear Physics is interesting!



一清、二白、三紅、四綠、五黃

clear lectures (清)?

謝謝！