

Nuclear Reactions

Shape, interaction, and excitation structures of nuclei ← scattering expt.
cf. Experiment by Rutherford (α scatt.)

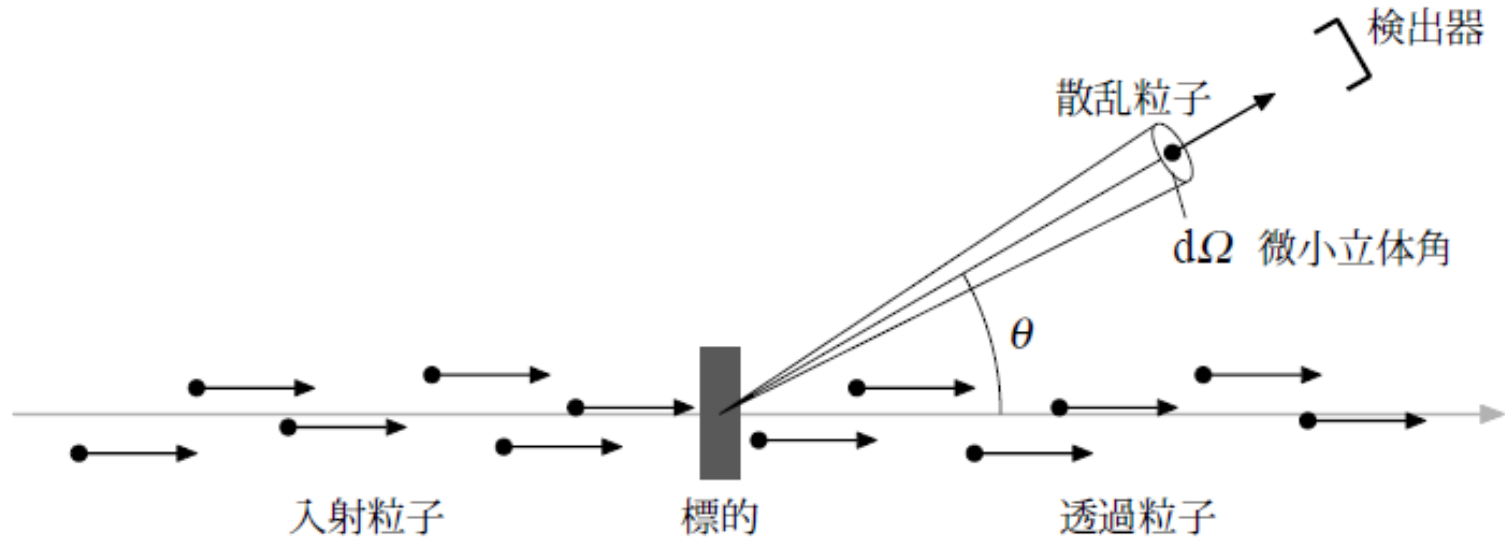
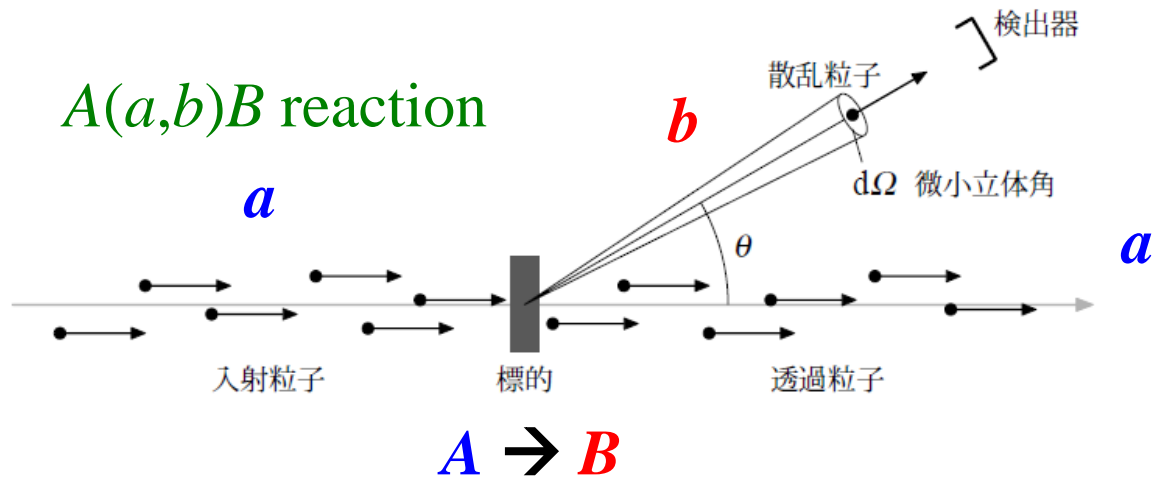


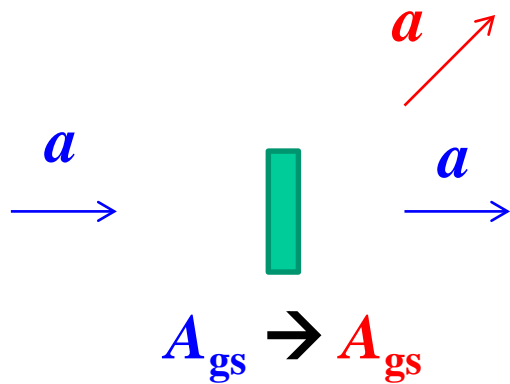
図 21.1: 散乱実験

http://www.th.phys.titech.ac.jp/~muto/lectures/QMII11/QMII11_chap21.pdf

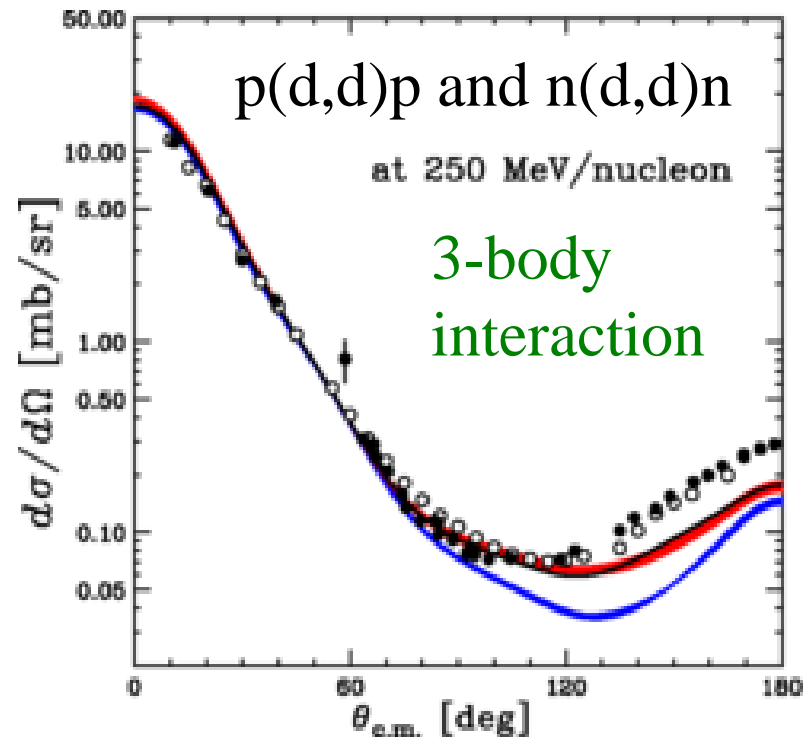
武藤一雄氏(東工大)

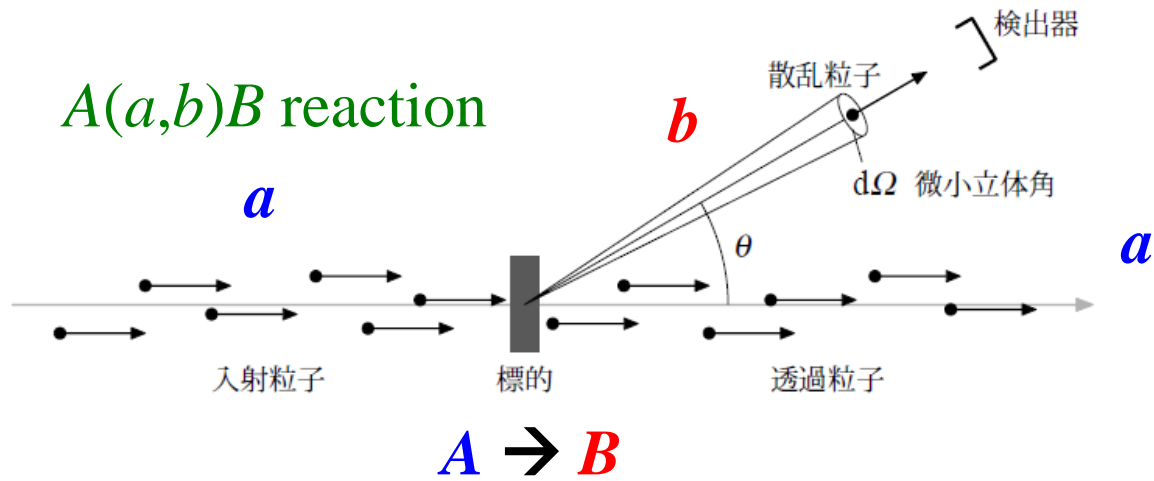


✓ elastic scattering

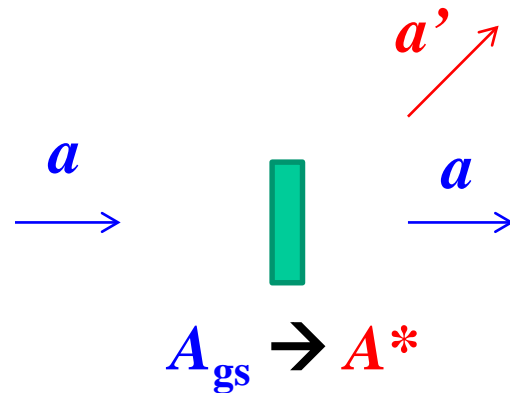


fundamental interaction
between a and A

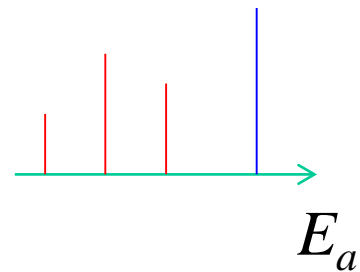


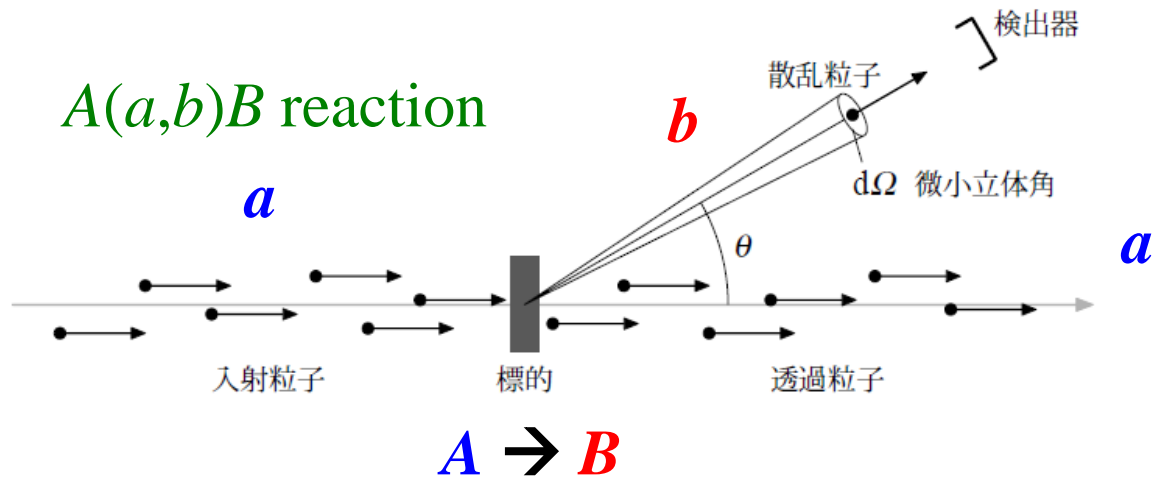


✓ inelastic scattering

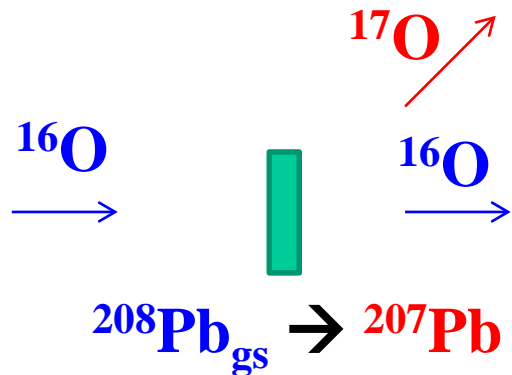


excitation spectrum
of a nucleus A



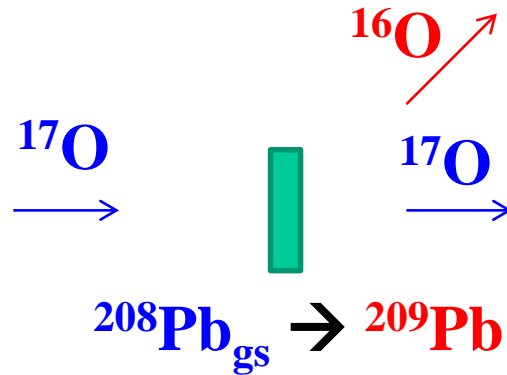


✓ transfer reaction
(pick-up reaction)

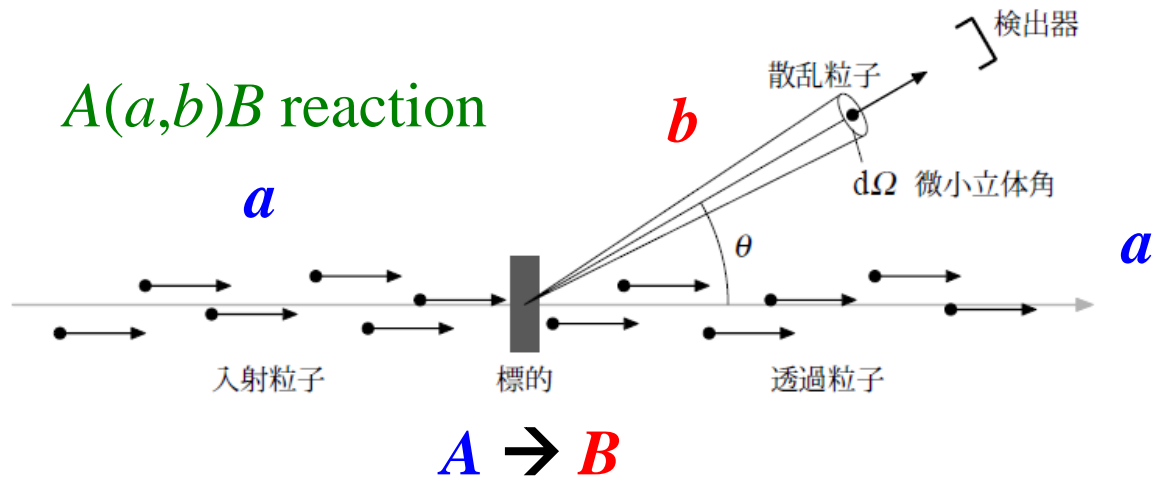


level schem of ^{207}Pb

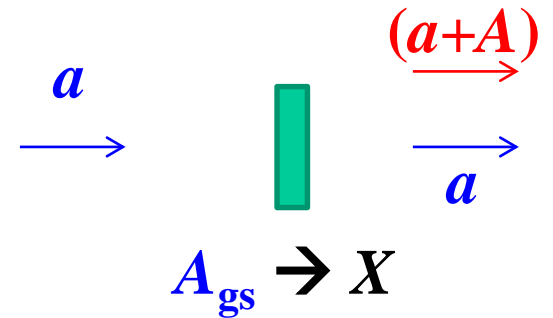
✓ transfer reaction
(stripping reaction)



level schem of ^{209}Pb

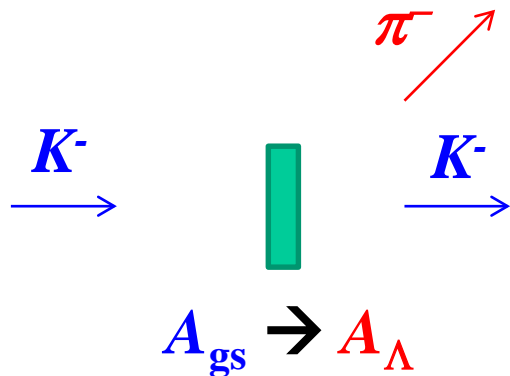


✓ fusion reaction

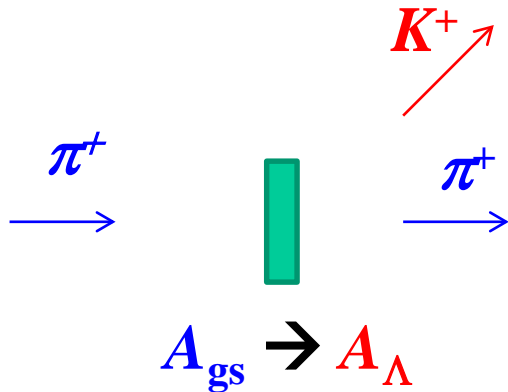


- interaction between a and A
- structure of a and A

✓ (K^-, π^-) reaction

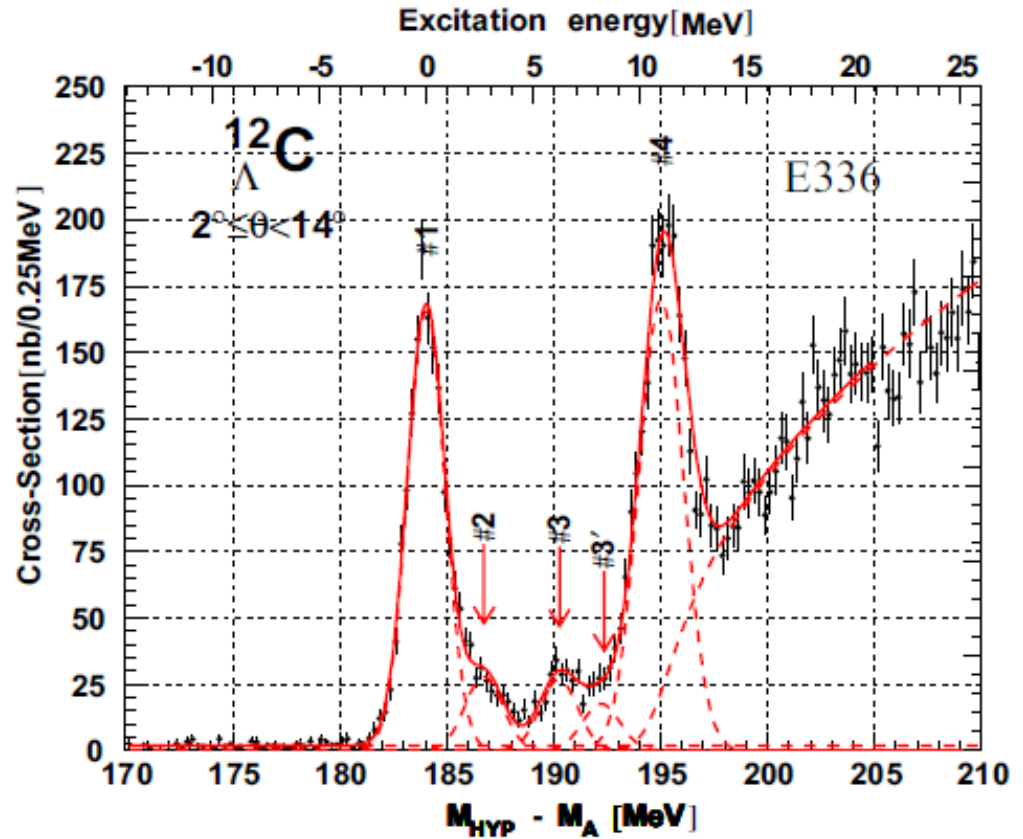


✓ (π^+, K^+) reaction



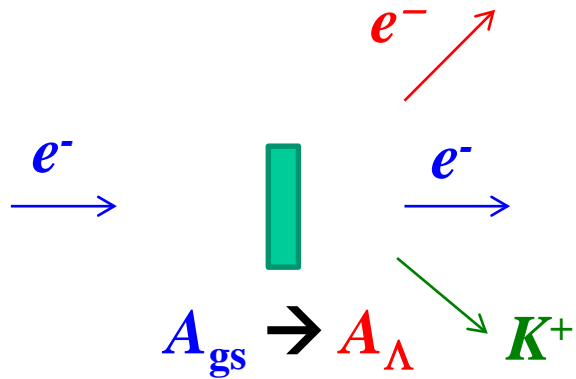
excitation spectrum
of a hypernucleus A_{Λ}

$^{12}\text{C} (\pi^+, K^+) ^{12}_{\Lambda}\text{C}$ reaction



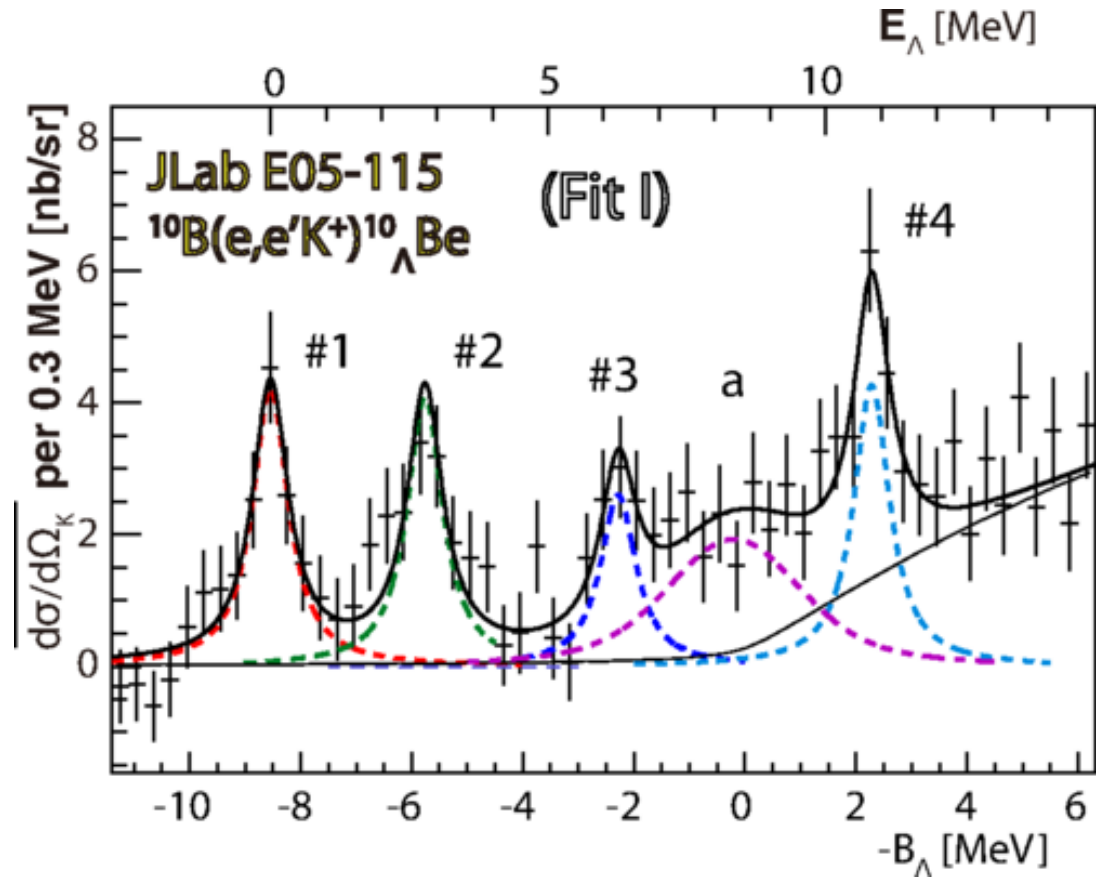
O. Hashimoto and H. Tamura,
Prog. in Part. and Nucl. Phys. 57 ('06)564

✓(e,e'K⁺) reaction



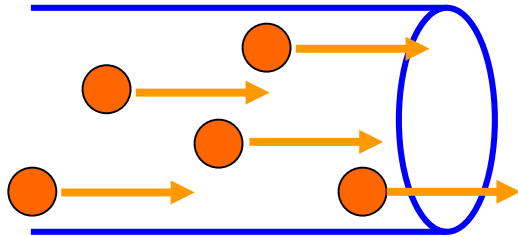
S.N. Nakamura et al.,
PRL110('13)012502

$^{10}\text{B}(e, e' K^+) ^{10}_{\Lambda}\text{Be}$



T. Gogami et al.,
PRC93 ('16) 034314

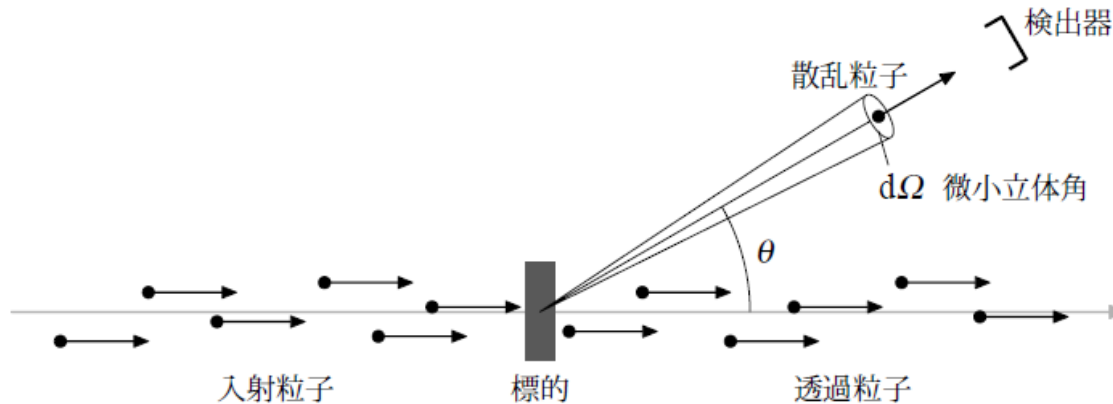
Cross sections



incident beam

flux = the number of particles
crossing unit area
per unit time

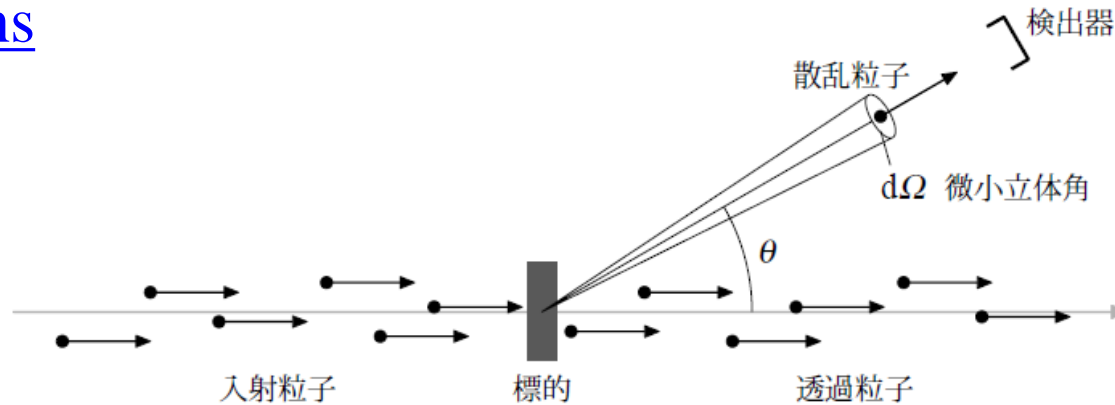
$$j = \rho_P \cdot v$$



event rate (the number of event per unit time per target nucleus)
: proportional to the incident flux

→ $R = N_T \cdot \sigma \cdot j$ ← cross section

Cross sections



event rate (the number of event per unit time per target nucleus)
: proportional to the incident flux

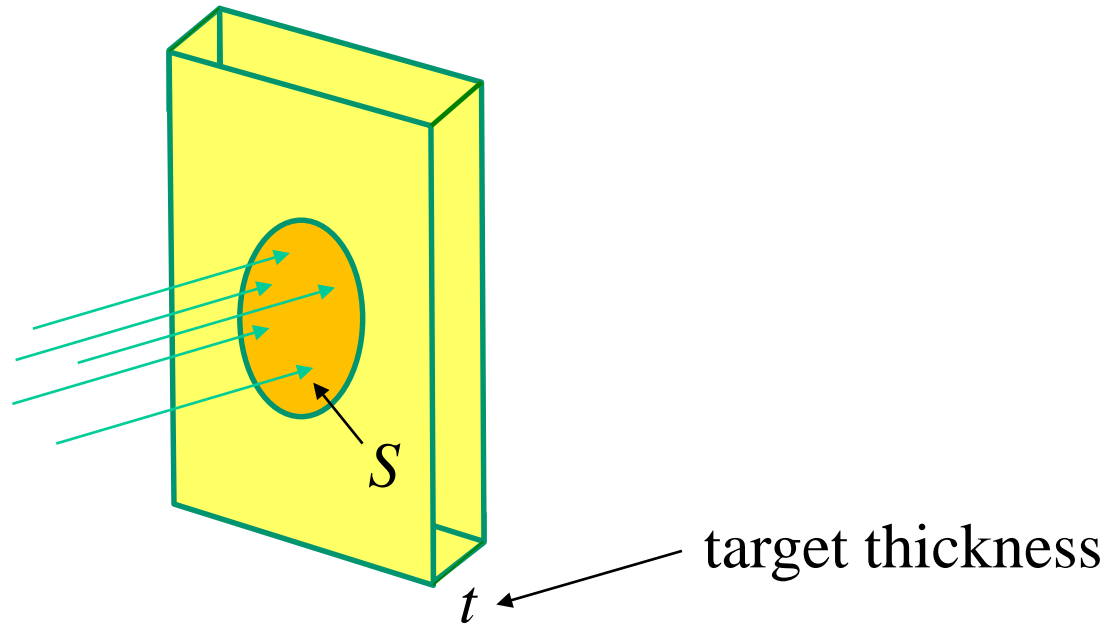
$$\longrightarrow R = N_T \cdot \sigma \cdot j \quad \text{cross section}$$

differential cross sections (angular distribution)

$$dR(\theta, \phi) = N_T \cdot \frac{d\sigma}{d\Omega} \cdot j \cdot d\Omega \quad \sigma = \int d\Omega \frac{d\sigma}{d\Omega}$$

units: 1 barn = $10^{-24} \text{ cm}^2 = 100 \text{ fm}^2$ (1 mb = $10^{-3} \text{ b} = 0.1 \text{ fm}^2$)

Cross sections (experiments)



$$dR(\theta, \phi) = N_{\text{T}} \cdot \frac{d\sigma}{d\Omega} \cdot j \cdot d\Omega$$

beam intensity: $I = j \cdot S$

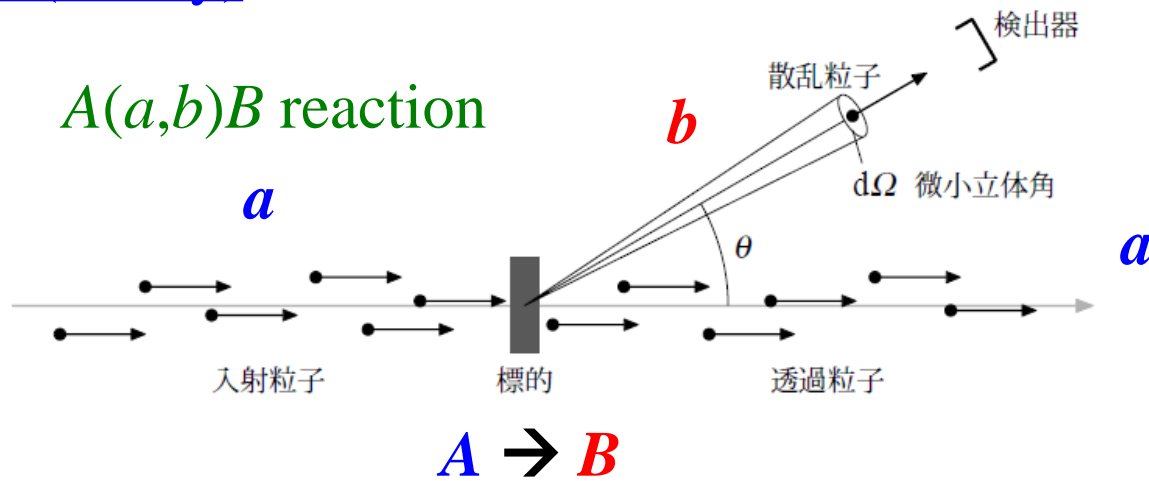
the number of target nucleus: $N_{\text{T}} = S \cdot t \cdot \rho_{\text{T}}$



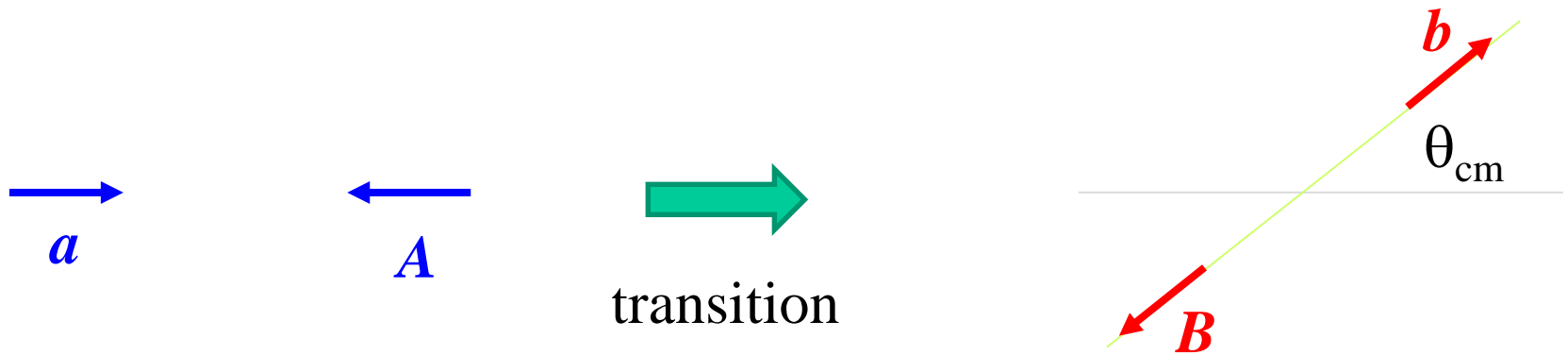
$$dR(\theta, \phi) = I \cdot \frac{d\sigma}{d\Omega} \cdot t \cdot \rho_{\text{T}} \cdot d\Omega \cdot \epsilon$$

← detection efficiency

Cross sections (theory)



center of mass frame



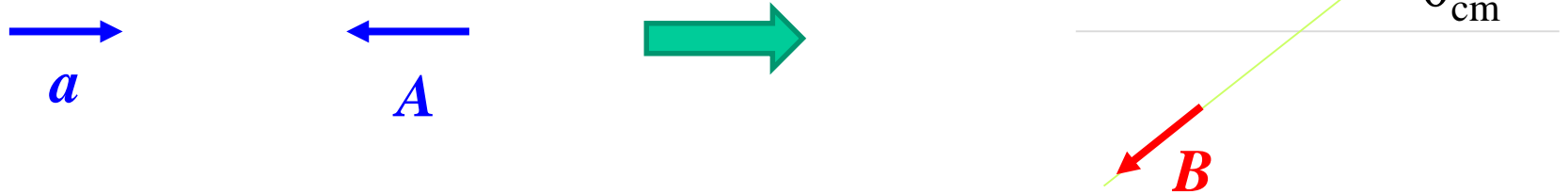
$$\frac{d\sigma}{d\Omega} = \frac{R}{j_{in}}$$

Cross sections

✓ laboratory frame



✓ center of mass frame



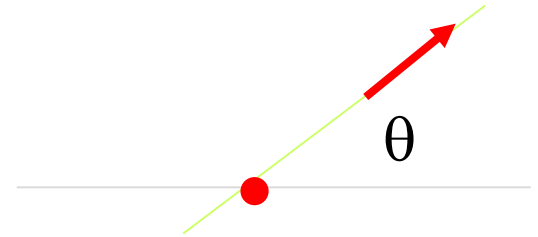
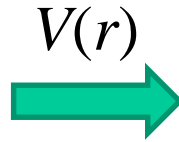
□ transformation ← energy and momentum conservations

$$\tan \theta_{\text{lab}} = \frac{\sin \theta_{\text{cm}}}{\gamma + \cos \theta_{\text{cm}}}, \quad d\Omega_{\text{lab}} = \frac{|1 + \gamma \cos \theta_{\text{cm}}|}{(1 + \gamma^2 + 2\gamma \cos \theta_{\text{cm}})^{3/2}} d\Omega_{\text{cm}}$$
$$E_{\text{cm}} = \frac{M_A}{M_a + M_A} E_{\text{lab}}, \quad \gamma = \sqrt{\frac{M_a M_b}{M_A M_B} \frac{E_{\text{cm}}}{E_{\text{cm}} + Q}}$$

Born approximation

$$\psi_f(\mathbf{r}) = e^{i\mathbf{p}_f \cdot \mathbf{r} / \hbar}$$

$$\psi_i(\mathbf{r}) = e^{i\mathbf{p}_i \cdot \mathbf{r} / \hbar}$$



$$\left(-\frac{\hbar^2}{2\mu} \nabla^2 + \underbrace{V(r)}_{\text{perturbation}} - E \right) \psi(\mathbf{r}) = 0$$

perturbation

transition rate for elastic scattering:

$$W_{fi} = \frac{2\pi}{\hbar} \int \frac{d\mathbf{p}_f}{(2\pi\hbar)^3} |\langle \psi_f | V | \psi_i \rangle|^2 \delta(E_f - E_i)$$

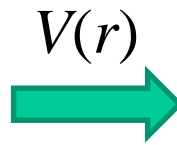
$$= \frac{\mu p_i}{4\pi^2 \hbar^4} \int d\Omega |\tilde{V}(\mathbf{q})|^2$$

$$\tilde{V}(\mathbf{q}) = \int d\mathbf{r} \psi_f^*(\mathbf{r}) V(\mathbf{r}) \psi_i(\mathbf{r})$$

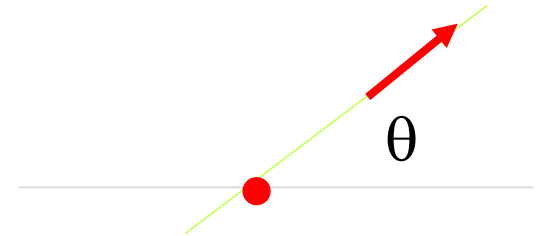
$$= \int d\mathbf{r} e^{i(\mathbf{p}_i - \mathbf{p}_f) \cdot \mathbf{r} / \hbar} V(\mathbf{r}) \equiv \int d\mathbf{r} e^{-i\mathbf{q} \cdot \mathbf{r}} V(\mathbf{r})$$

Born approximation

$$\psi_i(\mathbf{r}) = e^{i\mathbf{p}_i \cdot \mathbf{r} / \hbar}$$



$$\psi_f(\mathbf{r}) = e^{i\mathbf{p}_f \cdot \mathbf{r} / \hbar}$$



$$W_{fi} = \frac{\mu p_i}{4\pi^2 \hbar^4} \int d\Omega |\tilde{V}(\mathbf{q})|^2$$

$$\tilde{V}(\mathbf{q}) = \int d\mathbf{r} e^{i(\mathbf{p}_i - \mathbf{p}_f) \cdot \mathbf{r} / \hbar} V(r) \equiv \int d\mathbf{r} e^{-i\mathbf{q} \cdot \mathbf{r}} V(r)$$

momentum transfer

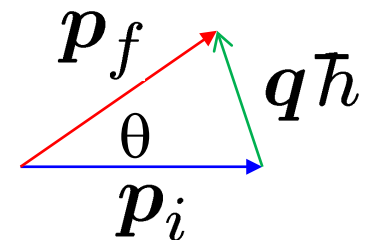


incident flux: $j_{\text{inc}} = \rho_i v = p_i / \mu$



$$\sigma = \frac{W_{fi}}{j_{\text{inc}}} = \int d\Omega \frac{\mu^2}{4\pi^2 \hbar^4} |\tilde{V}(\mathbf{q})|^2$$

$$= \frac{d\sigma}{d\Omega}$$



$$q\hbar = 2p_i \sin \frac{\theta}{2}$$

Electron scattering

$$V(r) = -e^2 \int d\mathbf{r}' \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

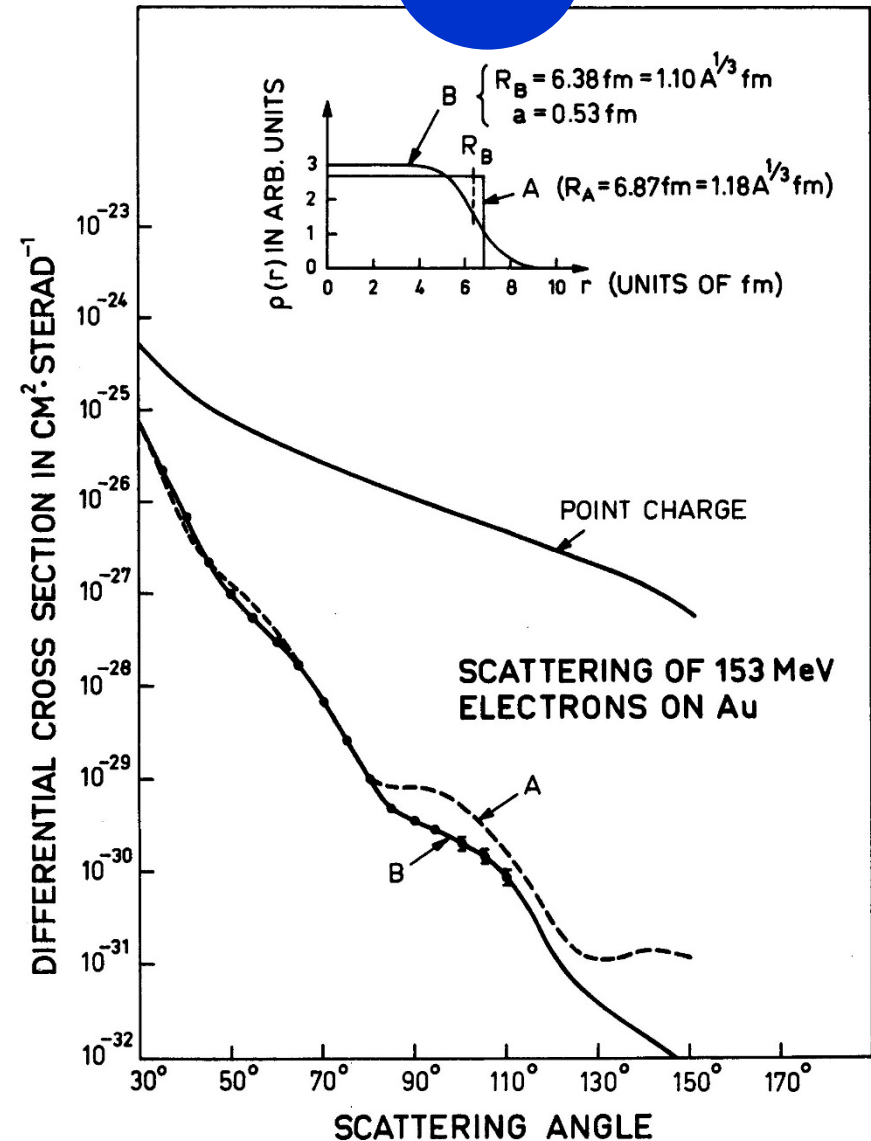
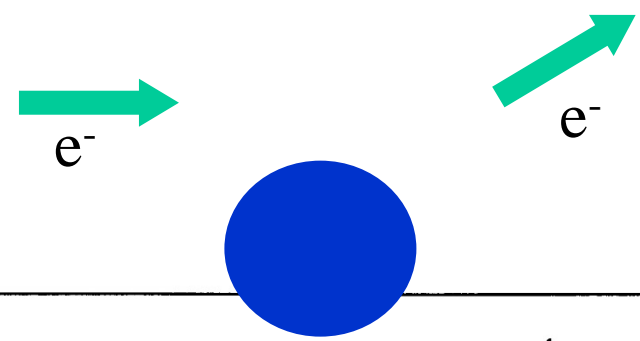
$$\begin{aligned} \frac{d\sigma}{d\Omega} &= \frac{Z_P^2 e^4}{(4E \sin^2 \theta/2)^2} |F(\mathbf{q})|^2 \\ &= \left(\frac{d\sigma_{\text{Ruth}}}{d\Omega} \right) |F(\mathbf{q})|^2 \end{aligned}$$

Form factor

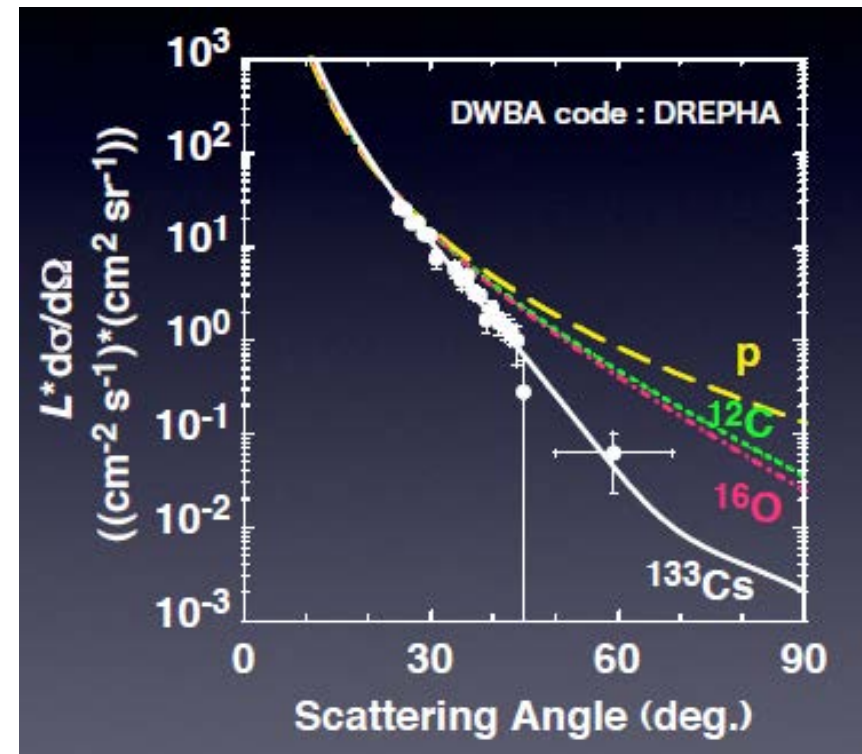
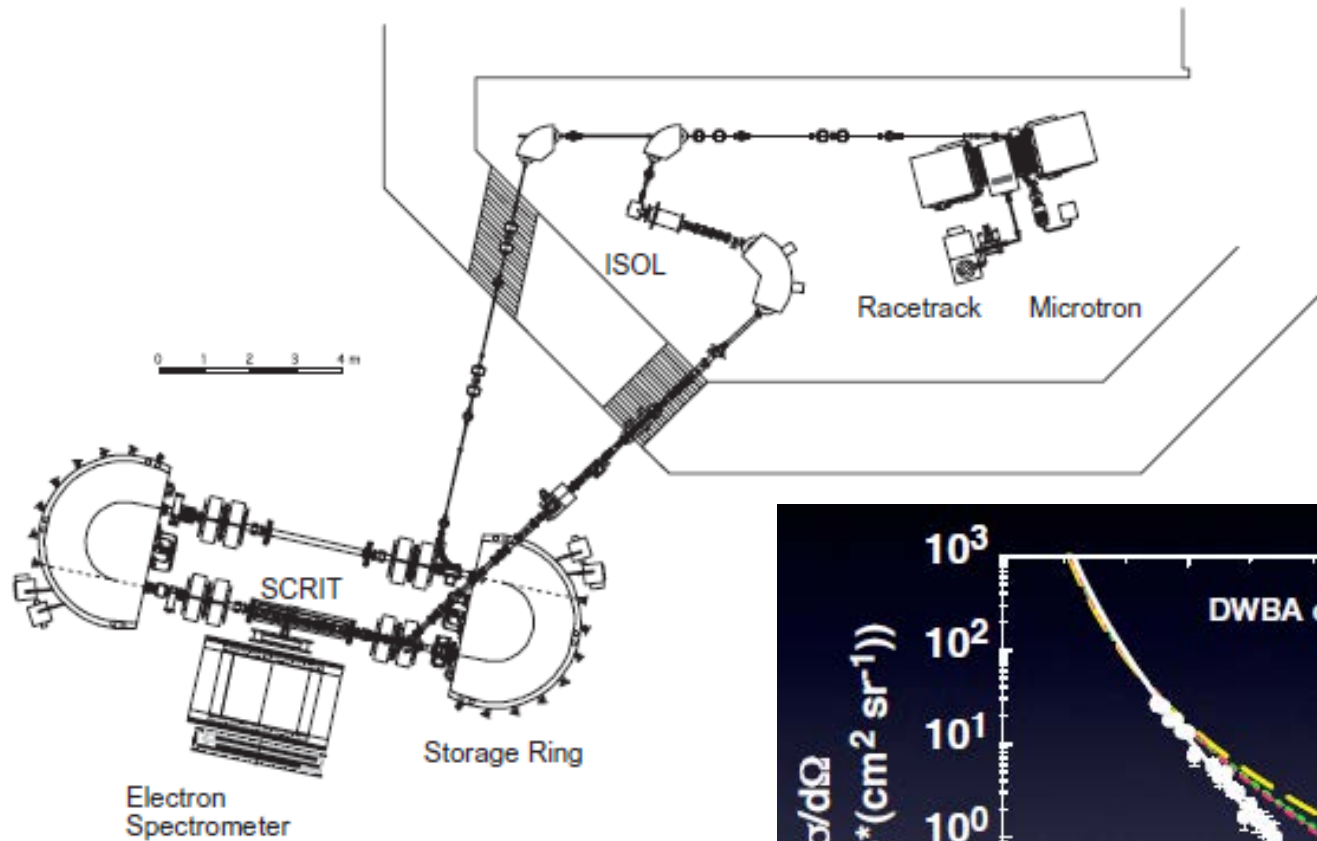
$$F(\mathbf{q}) = \int e^{-i\mathbf{q}\cdot\mathbf{r}} \rho(\mathbf{r}) d\mathbf{r}$$

* relativistic correction:

$$\begin{aligned} \frac{d\sigma_{\text{Ruth}}}{d\Omega} &\rightarrow \frac{d\sigma_{\text{Mott}}}{d\Omega} \\ &= \frac{d\sigma_{\text{Ruth}}}{d\Omega} \cdot \left(1 - \frac{v^2}{c^2} \sin^2 \frac{\theta}{2} \right) \\ &\sim \frac{d\sigma_{\text{Ruth}}}{d\Omega} \cdot \cos^2 \frac{\theta}{2} \quad (v \rightarrow c) \end{aligned}$$



cf. electron scattering off unstable nuclei (SCRIT)

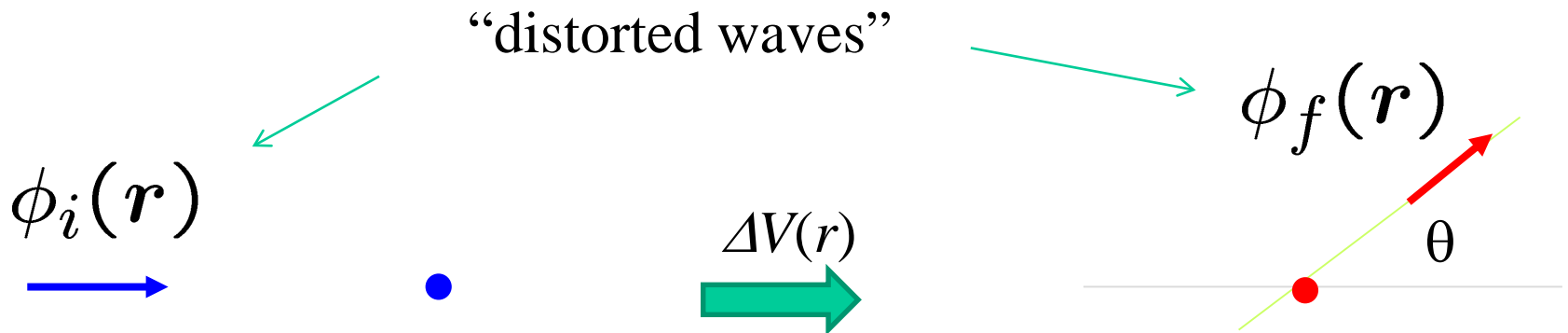


T. Suda et al.,
PTEP 2012, 03C008 (2012)
PRL102, 102501 (2009)

Distorted Wave Born approximation (DWBA)

$$\left(-\frac{\hbar^2}{2\mu} \nabla^2 + \underbrace{V(r)}_{\text{perturbation}} - E \right) \psi(\mathbf{r}) = 0$$

→
$$\left(-\frac{\hbar^2}{2\mu} \nabla^2 + V_0(r) + \underbrace{V(r) - V_0(r)}_{\text{perturbation}} - E \right) \psi(\mathbf{r}) = 0$$

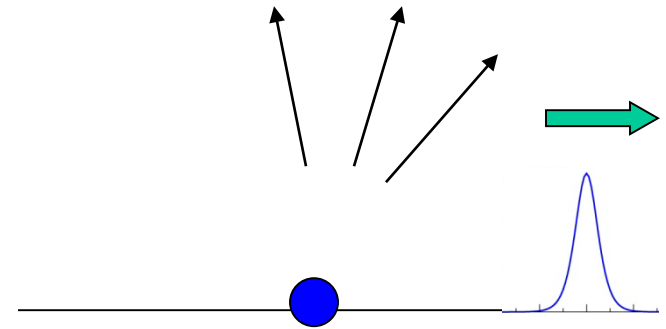
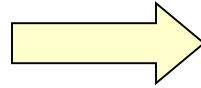


- ✓ inelastic scattering
- ✓ transfer reactions

Optical model

Reaction processes

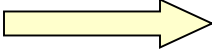
- Elastic scatt.
- Inelastic scatt.
- Transfer reaction
- Compound nucleus formation (fusion)



Loss of incident flux
(absorption)

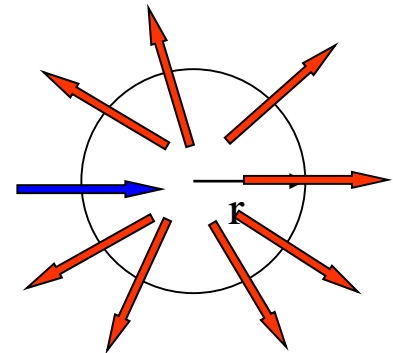
Optical potential

$$V_{\text{opt}}(\mathbf{r}) = V(\mathbf{r}) - iW(\mathbf{r}) \quad (W > 0)$$


$$\nabla \cdot \mathbf{j} = \dots = -\frac{2}{\hbar}W|\psi|^2$$

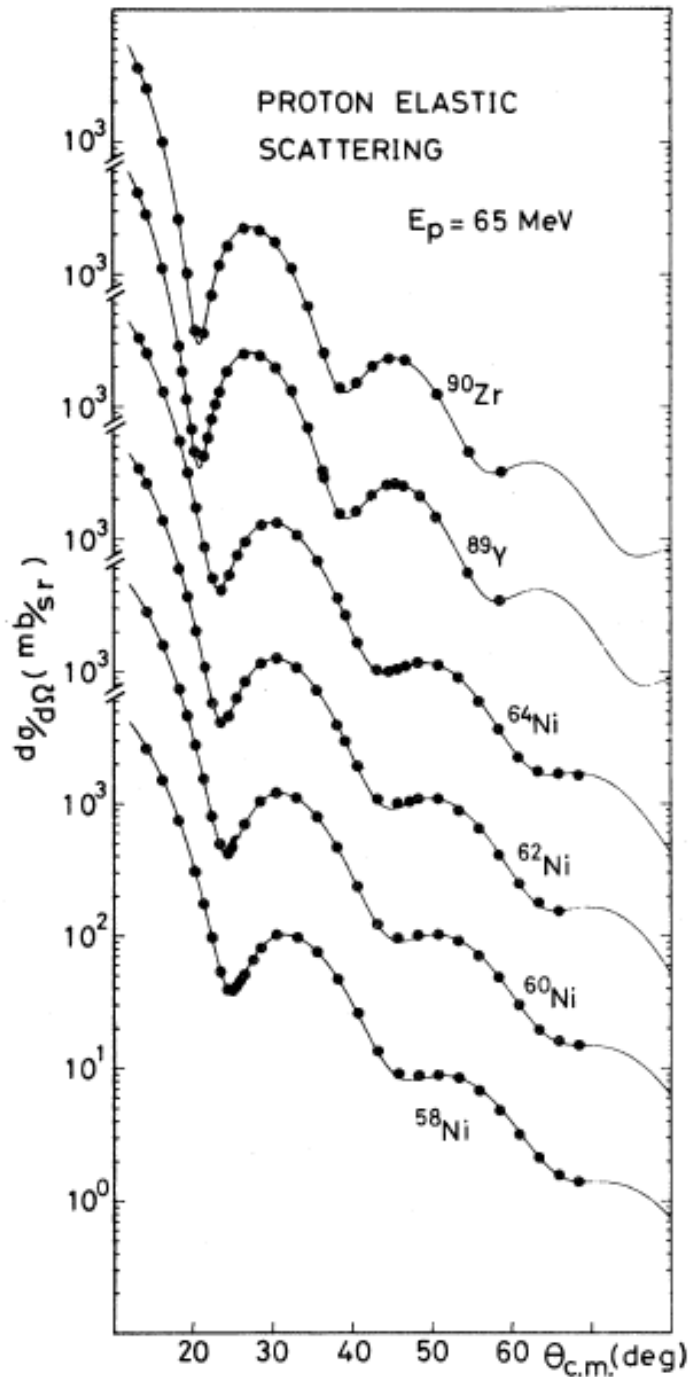
(note) Gauss's theorem

$$\int_S \mathbf{j} \cdot \mathbf{n} dS = \int_V \nabla \cdot \mathbf{j} dV$$



$$\left(-\frac{\hbar^2}{2\mu} \nabla^2 + \frac{Z_P Z_T e^2}{r} + V_{\text{opt}}(r) - E \right) \psi(\mathbf{r}) = 0$$

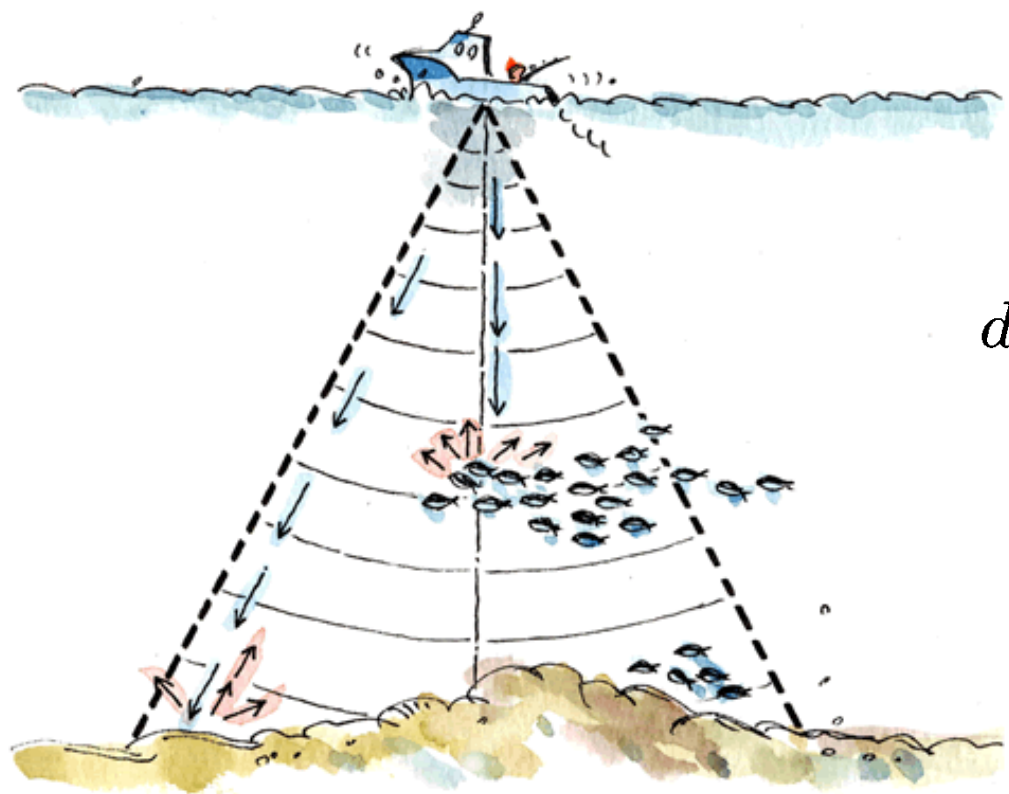
Woods-Saxon + volume & surface
imaginary parts



H. Sakaguchi et al.,
PRC26 (1982) 944

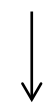
おまけ: 海洋音響学におけるDWBA

魚群探知機



散乱体(魚など)による
(超)音波の(後方)散乱

$$dR(\theta, \phi) = N_T \cdot \frac{d\sigma}{d\Omega} \cdot j \cdot d\Omega$$



$$N_T = \frac{\frac{dR}{d\Omega}}{j \cdot \frac{d\sigma}{d\Omega}}$$

微分散乱断面積を知って
いれば魚の数 N_T がわかる

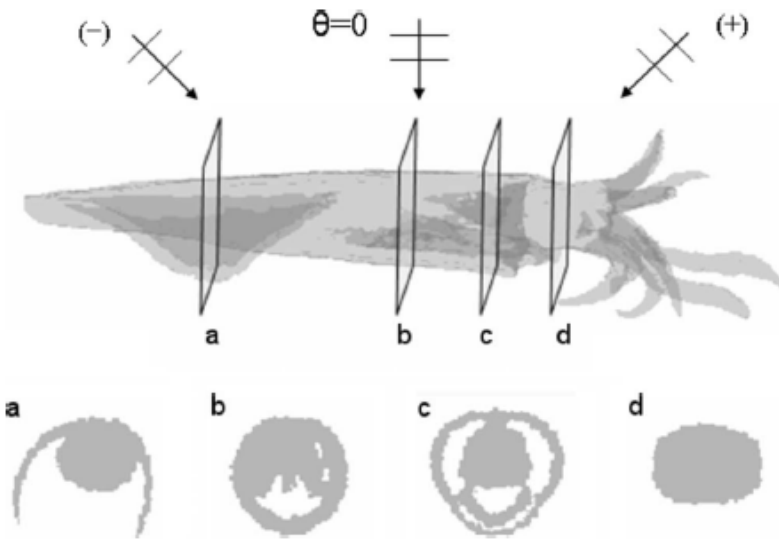
Use of the distorted wave Born approximation to predict scattering by inhomogeneous objects: Application to squid

Benjamin A. Jones,^{a)} Andone C. Lavery, and Timothy K. Stanton

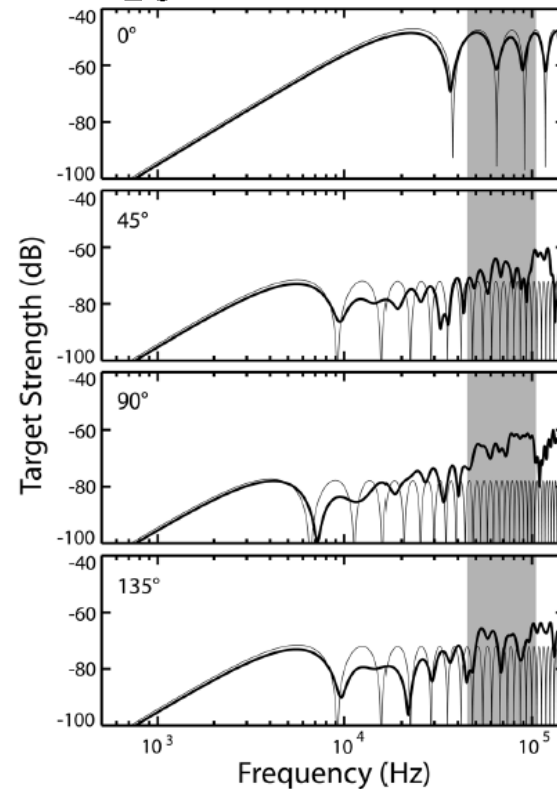
Department of Applied Ocean Physics and Engineering, Woods Hole Oceanographic Institution,
Woods Hole, Massachusetts 02543-1053

J. Acoust. Soc. Am. 125 ('09) 73

$10 \log_{10} \sigma$



イカのモデル化



- Arms-folded numerical model (no fins)
- - - Analytical prolate spheroid model ← !
- Usable band in the experiment

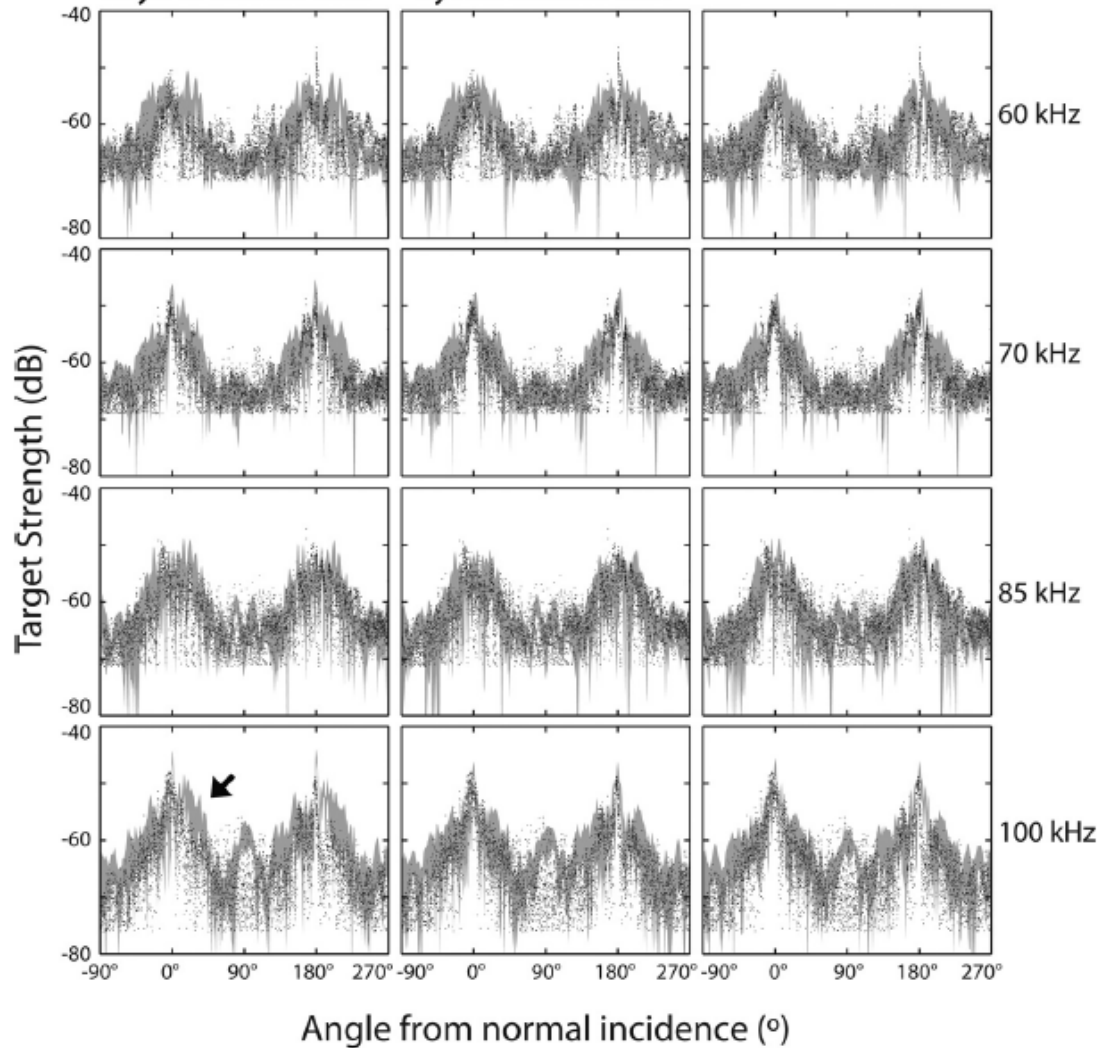
DWBA: イカの内部では局所的な波数を用いる



(A) Original asymmetrical fins

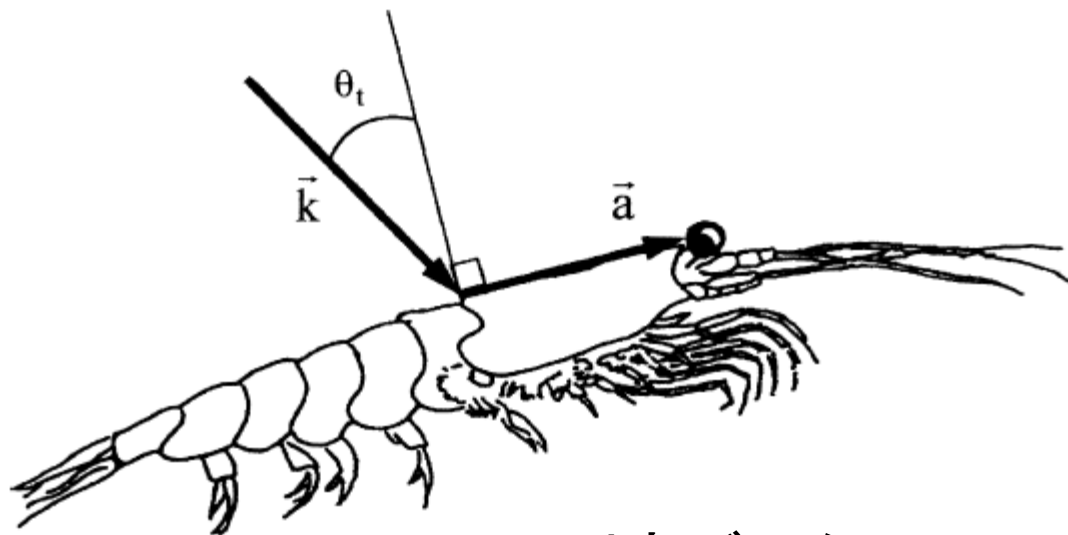
(B) Artificial symmetrical fins

(C) No fins



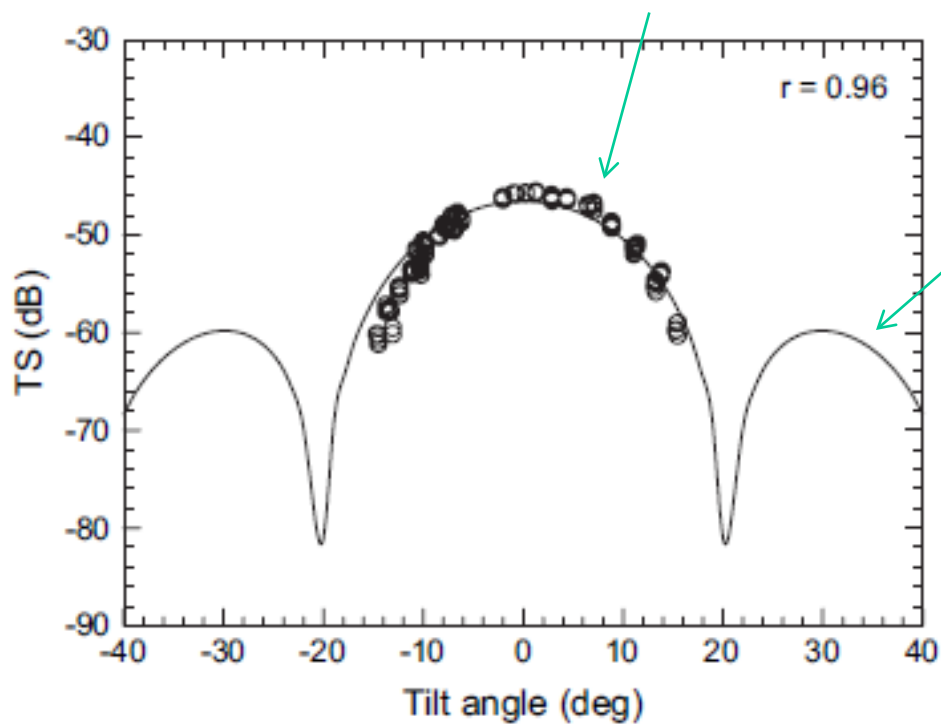
- Experimental data
- Numerical model + noise

W.-J. Lee, A.C. Lavery, T. Stanton,
J. Acoust. Soc. Am. 131 ('12) 4461



オキアミ

測定データ

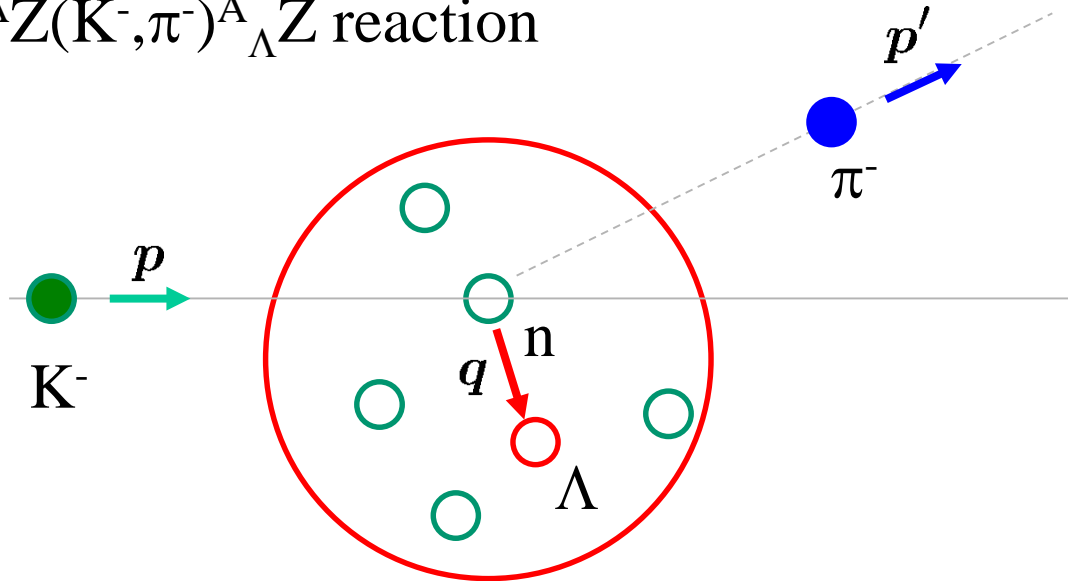


DWBA

K. Akamatsu and M. Furusawa,
 ICES J. of Marine Science 63 ('06) 36

Impulse approximation

example: ${}^A_Z(K^-, \pi^-) {}^A_{\Lambda}Z$ reaction



- ✓ high energy
- ✓ single scattering approximation
- ✓ (other nucleons: spectator)

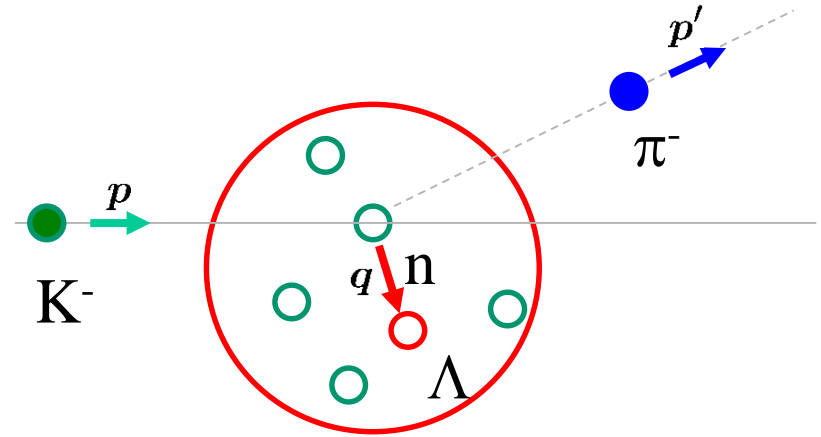
$$T_{fi} \sim \left\langle \psi_{\pi^-} \left| \left\langle \Psi_{\Lambda Z} \left| \sum_j v_{eff}(j) \right| \Psi_{AZ} \right\rangle \right| \psi_{K^-} \right\rangle$$

effective K-n interaction
(including multiple scattering)

Impulse approximation

example: ${}^A Z(K^-, \pi^-) {}^A_{\Lambda} Z$ reaction

- ✓ high energy
- ✓ single scattering approximation



$$T_{fi} \sim \left\langle \psi_{\pi^-} \left| \left\langle \Psi_{\Lambda Z} \left| \sum_j v_{eff}(j) \right| \Psi_{AZ} \right\rangle \right| \psi_{K^-} \right\rangle$$

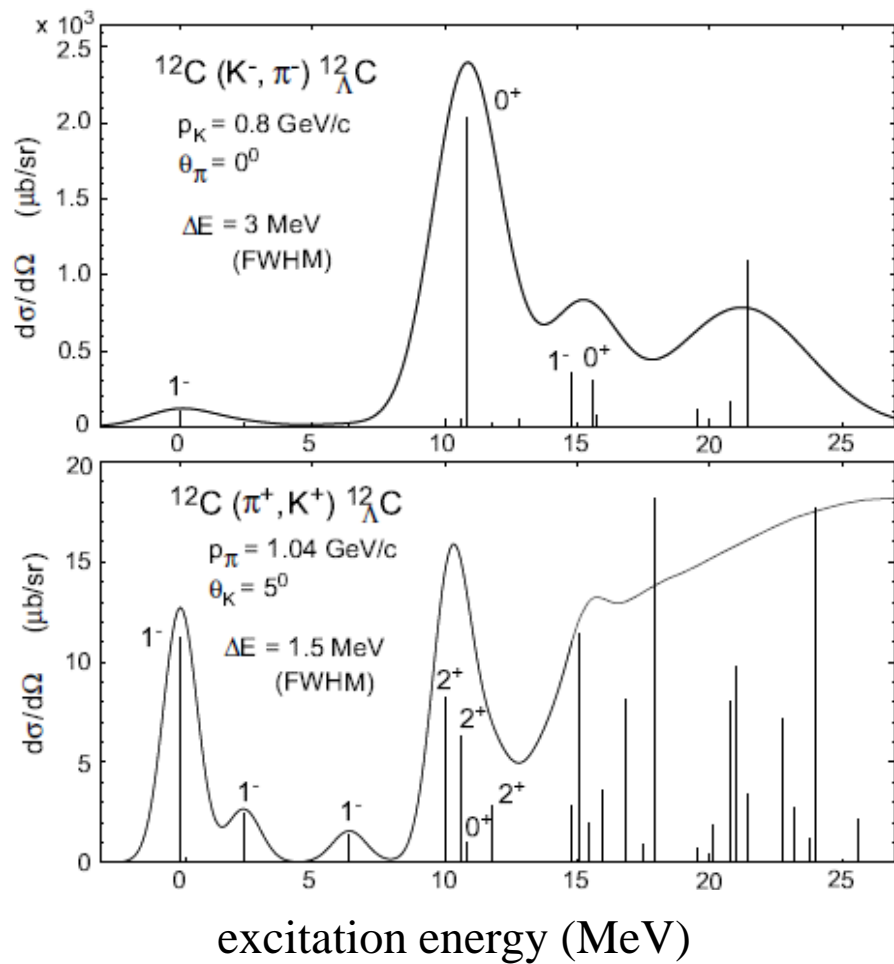
$$\frac{d\sigma}{d\Omega} \sim \underbrace{\alpha_{\text{kin}}}_{\text{kinematical factor}} \underbrace{\left(\frac{d\sigma}{d\Omega} \right)_{K^- n \rightarrow \pi^- \Lambda}}_{\text{elementary process}} N_{\text{eff}}(\theta; i \rightarrow f)$$

kinematical
factor

elementary process

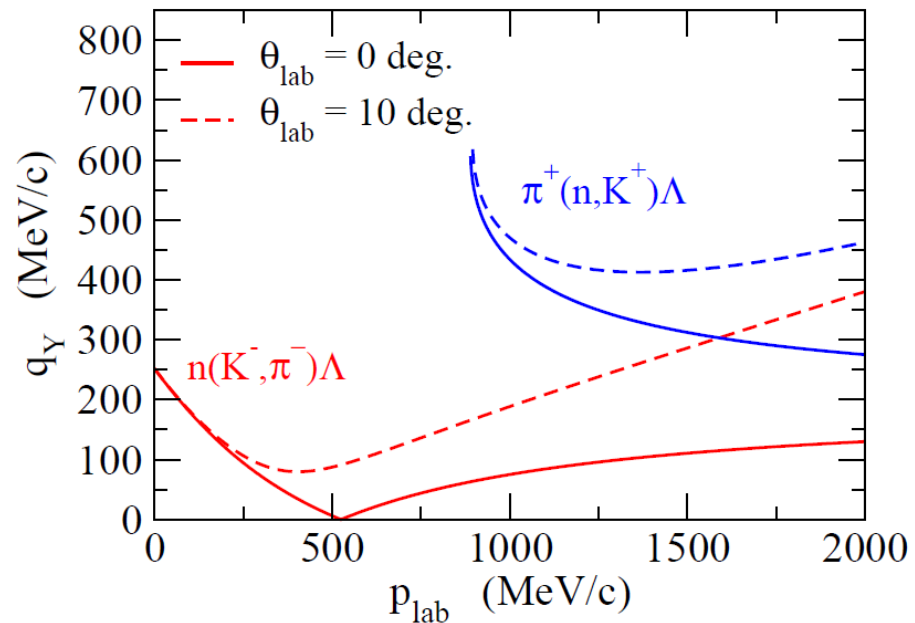
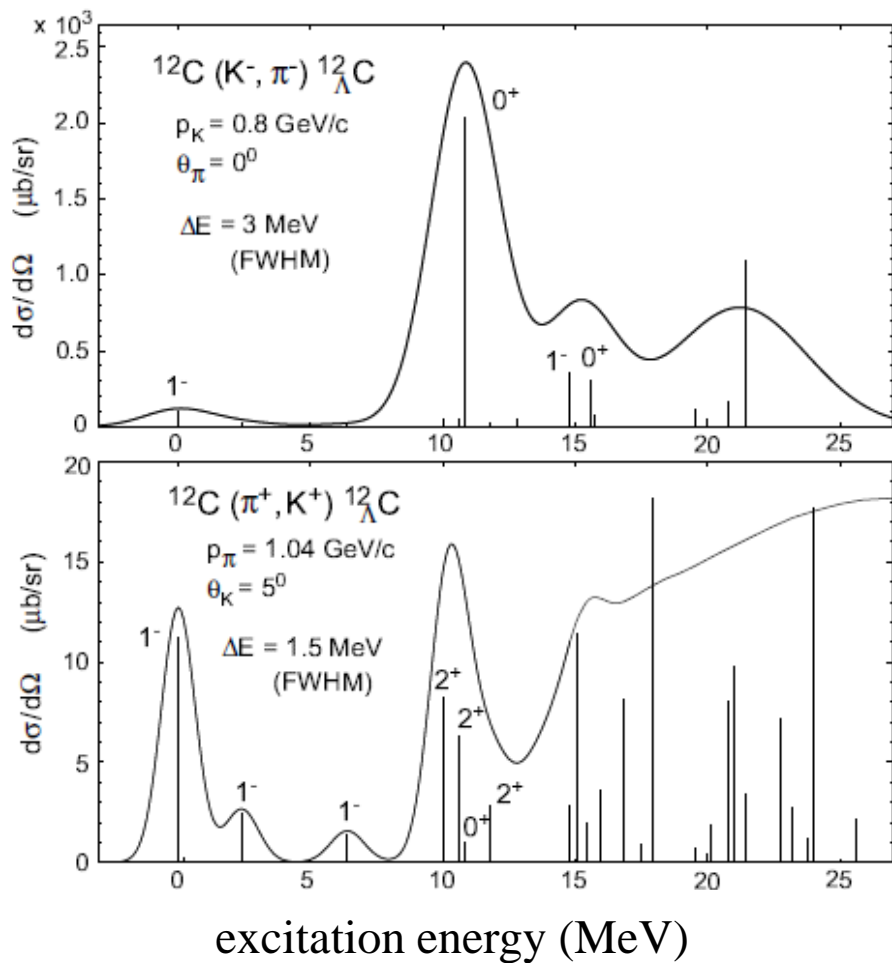
$$N_{\text{eff}}(\theta; i \rightarrow f) \sim \left| \int d\mathbf{r} \psi_{\pi^-}^*(\mathbf{r}) \underbrace{\varphi_{j\Lambda l_{\Lambda} m_{\Lambda}}^{(\Lambda)*}(\mathbf{r}) \varphi_{j n l_n m_n}^{(n)}(\mathbf{r})}_{\text{elementary process}} \psi_{K^-}(\mathbf{r}) \right|^2$$

- Plane wave impulse approximation (PWIA)
- Distorted wave impulse approximation (DWIA)



O. Hashimoto and H. Tamura,
 Prog. in Part. and Nucl. Phys. 57 ('06)564

T. Motoba et al., PRC38('88)1322



$$m_n + m_{\text{K}} = 1432 \text{ MeV} \quad \left. \begin{array}{l} \leftarrow \\ \leftarrow \end{array} \right\} Q > 0$$

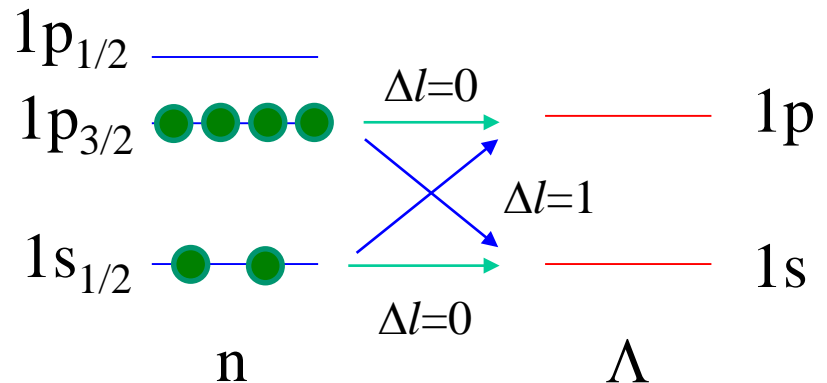
$$m_{\pi} + m_{\Lambda} = 1255.3 \text{ MeV} \quad \leftarrow$$

$$m_{\pi} + m_n = 1079.2 \text{ MeV} \quad \left. \begin{array}{l} \leftarrow \\ \leftarrow \end{array} \right\} Q < 0$$

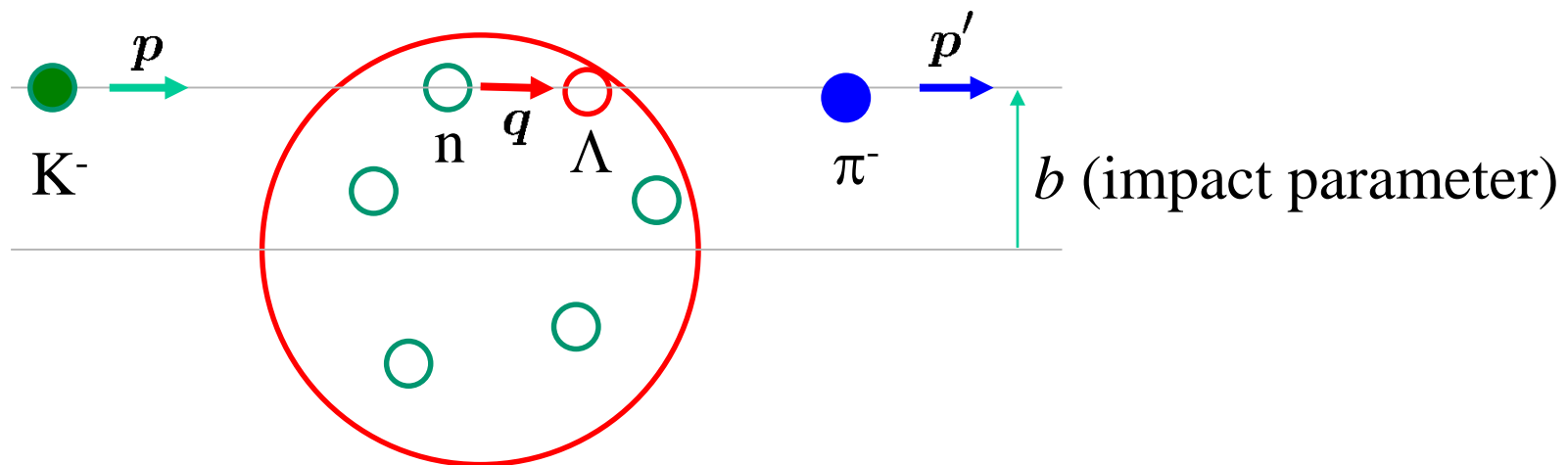
$$m_{\text{K}} + m_{\Lambda} = 1609.4 \text{ MeV} \quad \leftarrow$$

O. Hashimoto and H. Tamura,
 Prog. in Part. and Nucl. Phys. 57 ('06)564

T. Motoba et al., PRC38('88)1322

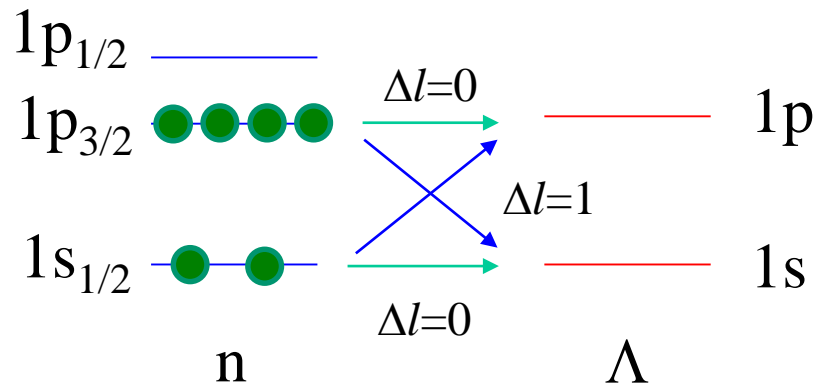
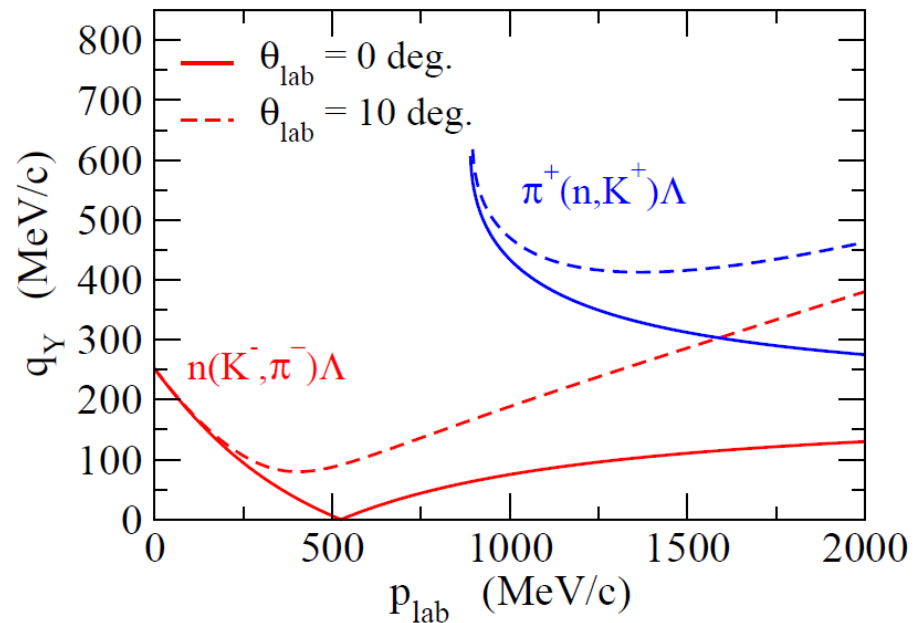
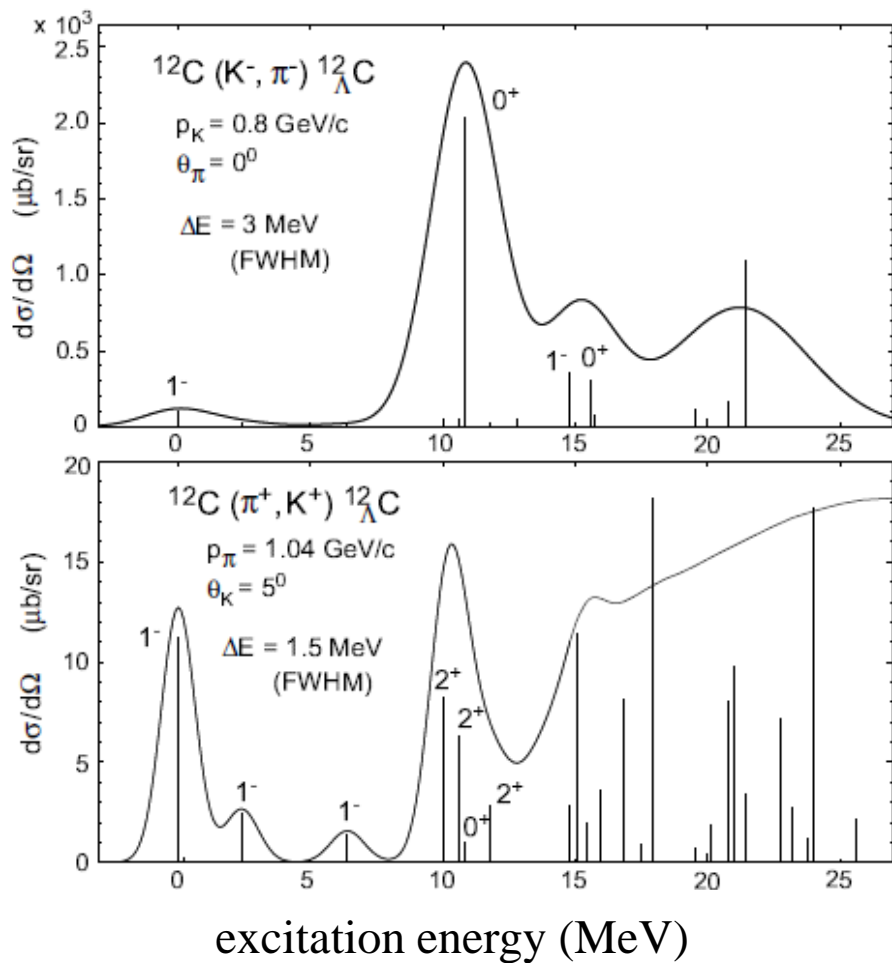


relation between q and Δl



$$l \sim kb \text{ (classically)}$$

➡ $\Delta l \sim b(p' - p) = bq$



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$$\Delta l \sim b(p' - p) = bq$$