

Advanced Nuclear Physics

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Tohoku University
Kouichi Hagino

原子核理論特論

東北大学
原子核理論研究室
萩野浩一

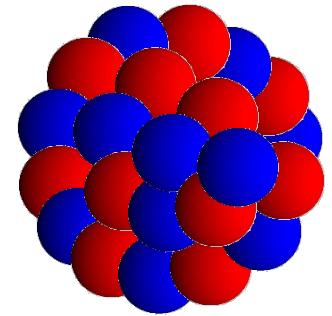
Contents

Nuclei: many-body systems of nucleons (protons and neutrons)



*Nuclear Many-Body Problems
with strong interaction*

(Low-energy) Nuclear Physics



to understand rich nature of atomic nuclei starting from nucleon-nucleon interactions

- size, mass, density, shape
- excitations
- decays
- nuclear reactions

two kinds of particle: protons and neutrons

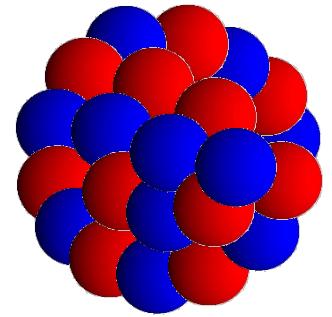
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Nuclei: many-body systems of nucleons (protons and neutrons)

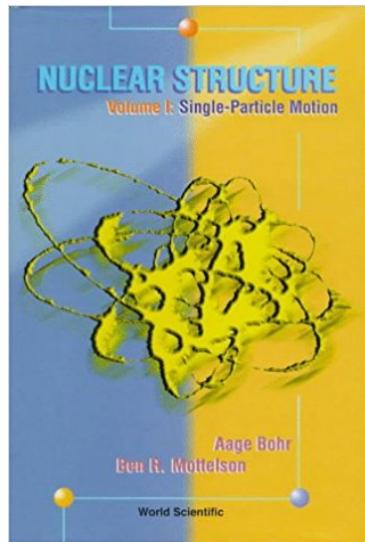
→ *Nuclear Many-Body Problems*

this lecture: microscopic descriptions of atomic nuclei

- Liquid drop model
- Single-particle motion and shell structure
- **Hartree-Fock approximation**
- Bruckner theory
- Pairing correlations and superfluid Nuclei
- nucleosynthesis and physics of neutron-rich nuclei
- **Random phase approximation (RPA)**
- **Nuclear reactions and superheavy elements
(physics of Nihonium)**



References



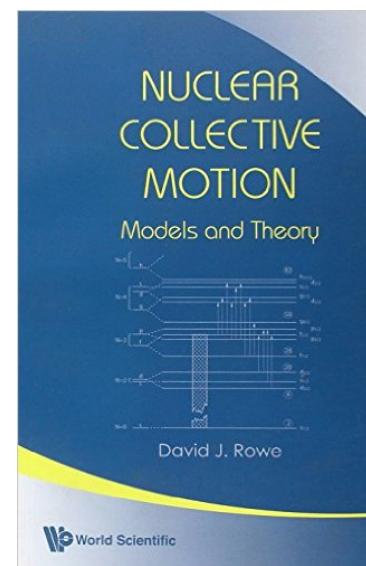
Bohr-Mottelson



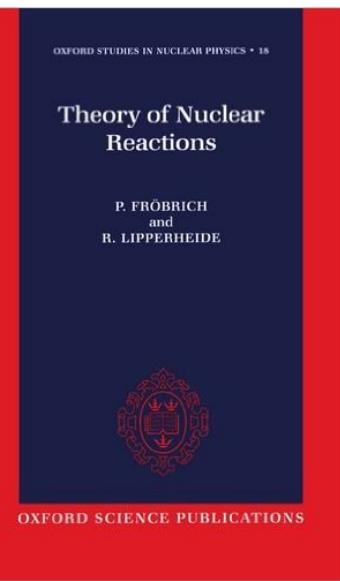
Ring-Schuck



Rowe



Frobrich
-Lipperheide



Lecture notes:

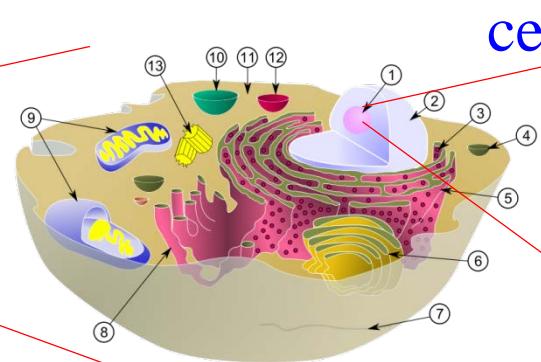
<http://www.nucl.phys.tohoku.ac.jp/~hagino/lecture.html>

(Tohoku University → Physics → Nuclear Theory
→ Kouichi Hagino → Lectures)

Introduction: atoms and atomic nuclei

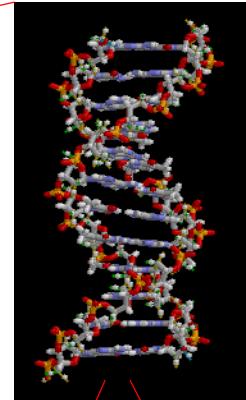


↔
 $\sim 50 \text{ cm}$



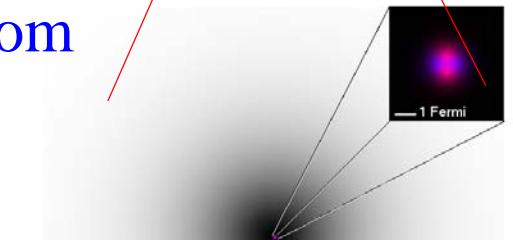
↔
 $\sim \mu\text{m} = 10^{-6} \text{ m}$

cells



$\sim 10^{-8} \text{ m}$

atom



↔
1 Ångstrom ($= 100,000 \text{ Fermi}$)



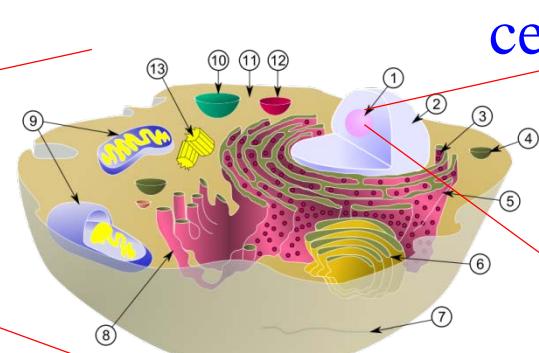
$\sim 10^{-10} \text{ m}$

everything is made of atoms.

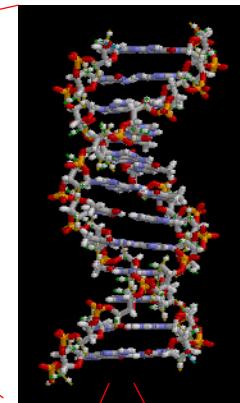
Introduction: atoms and atomic nuclei



↔
 $\sim 50 \text{ cm}$

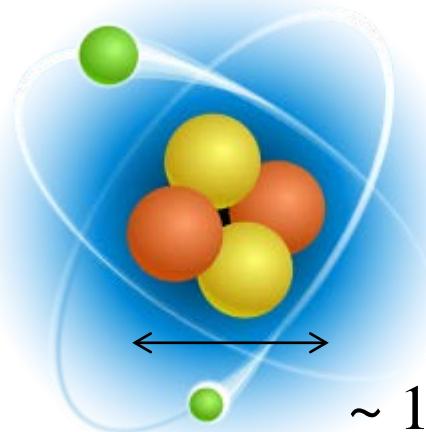


cells



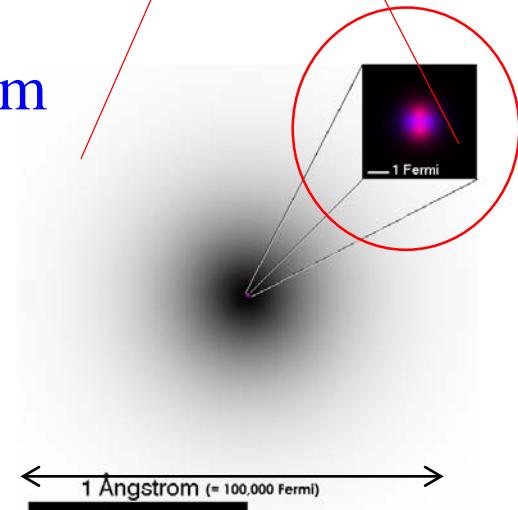
$\sim 10^{-8} \text{ m}$

atomic nucleus

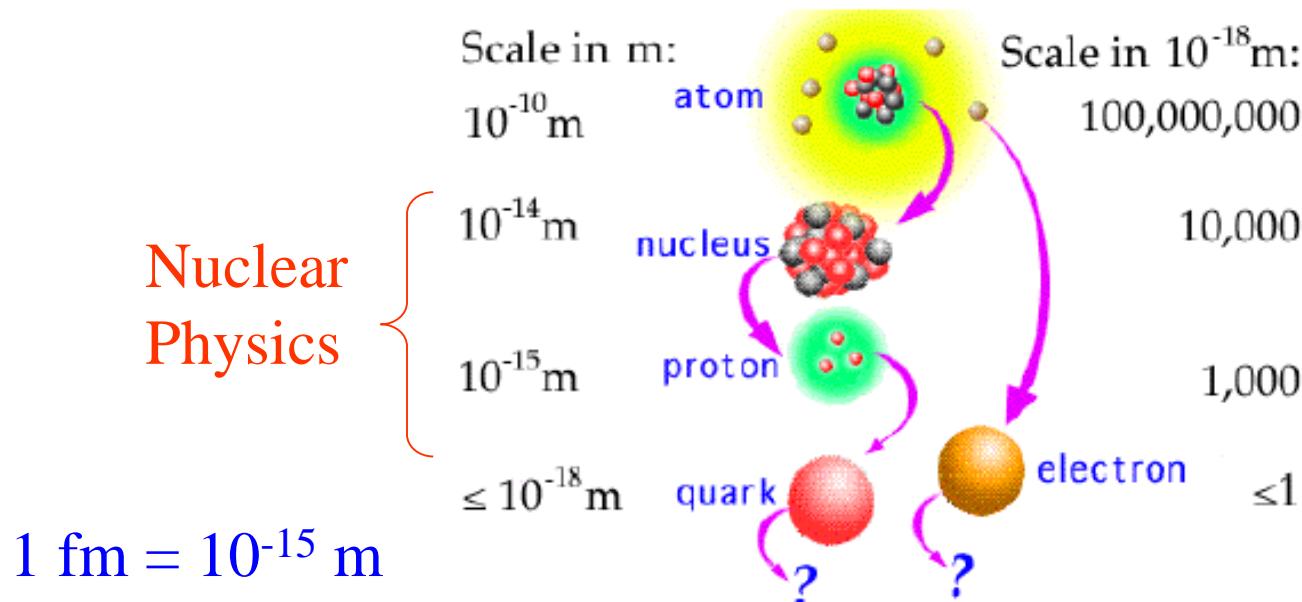


↔
 $\sim 10^{-15} \text{ m}$

atom



$\sim 10^{-10} \text{ m}$

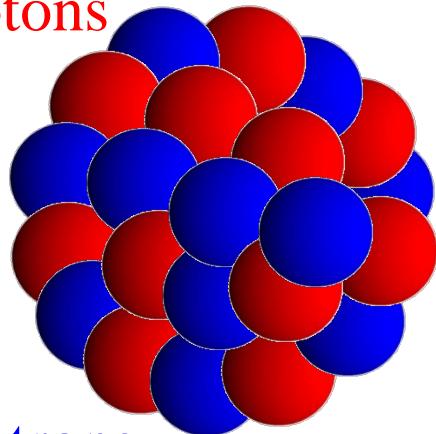


Nucleus as a *quantum many body system*

Basic ingredients:	charge	mass (MeV)	spin, parity
Proton	+e	938.256	$1/2^+$
Neutron	0	939.550	$1/2^+$

(note) $\text{n} \rightarrow \text{p} + \text{e}^- + \bar{\nu}_e$ (10.4 min)

protons

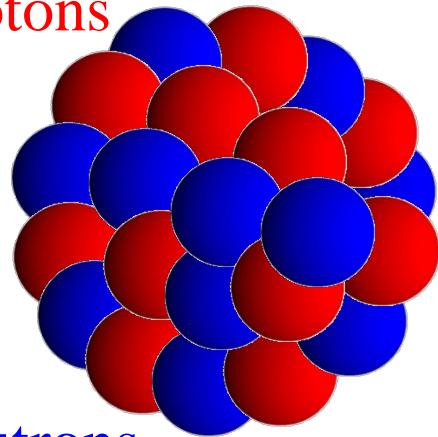


neutrons

- Nucleons are not stopping inside a nucleus.
(they move relatively freely)
- Yet, they are not completely independent.
A nucleus keeps its shape while nucleons influence among themselves so that a nucleon does not escape.

a self-bound system

protons



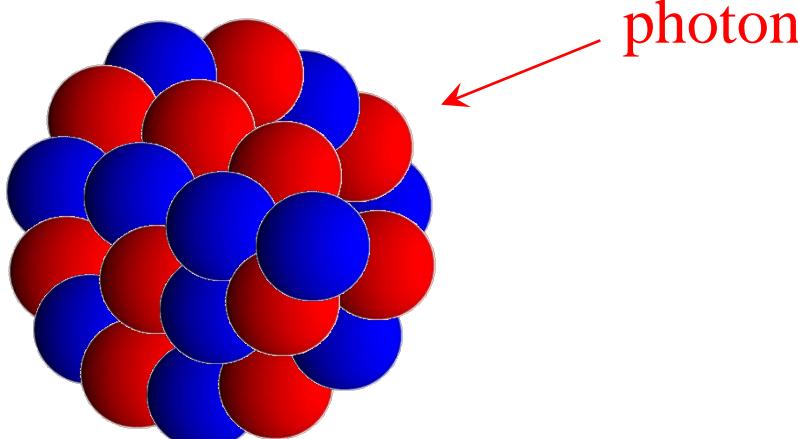
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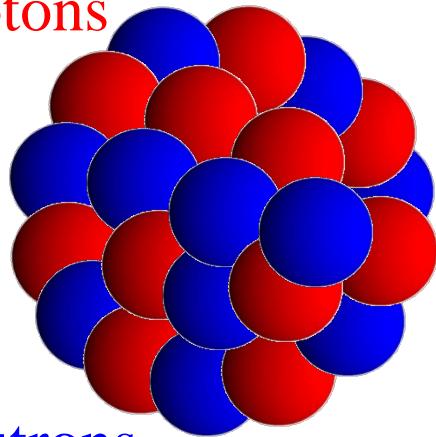
a self-bound system

What happens if a photon is absorbed into a nucleus?

- one nucleon simply starts moving faster?



protons



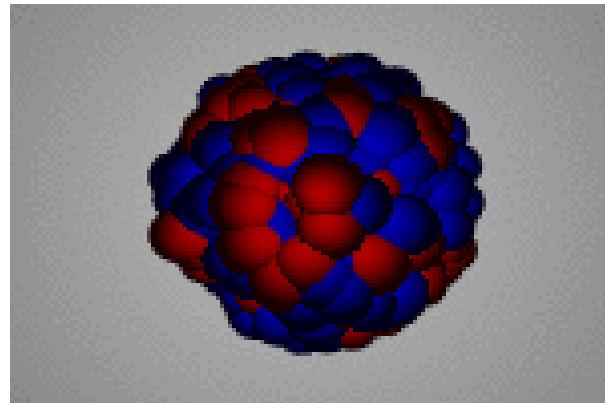
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a self-bound system

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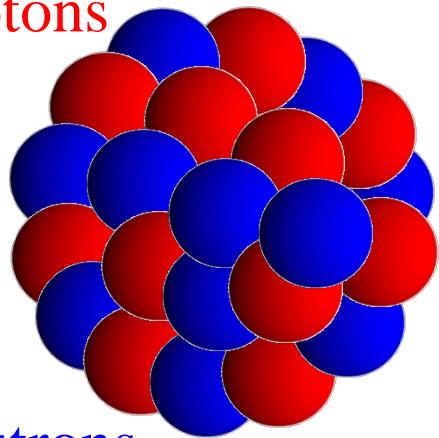
- one nucleon simply starts moving faster?



Very coherent
motion can happen

Collective motions

protons

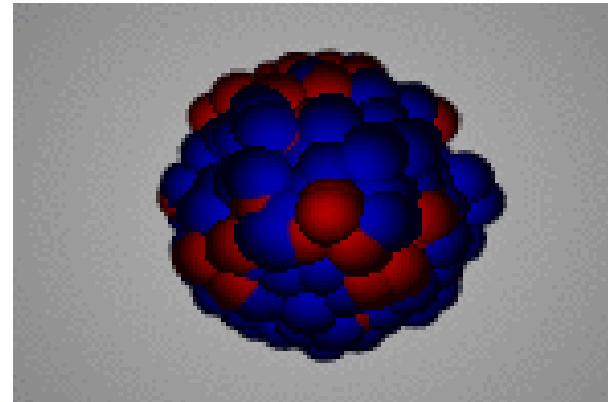
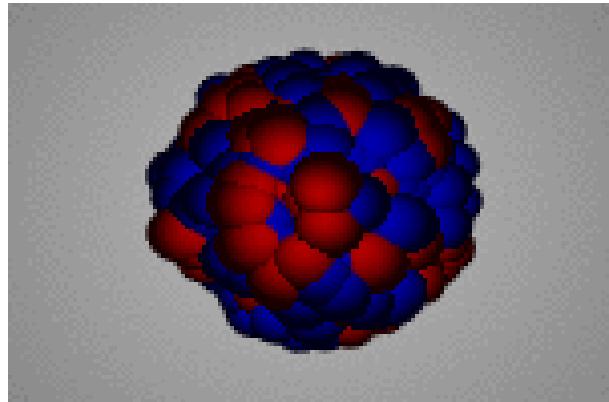


neutrons

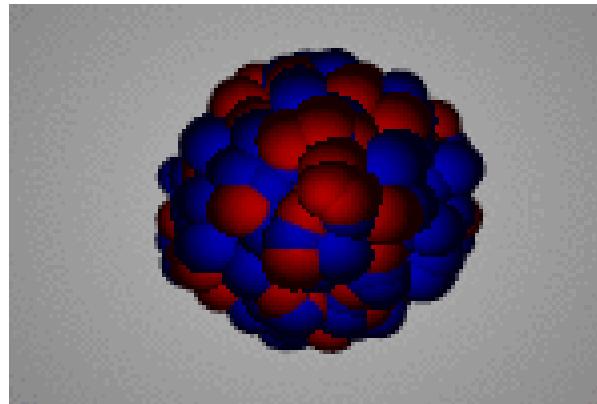
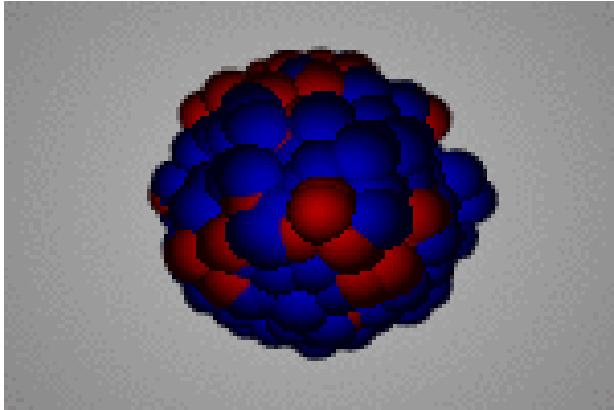
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a self-bound system

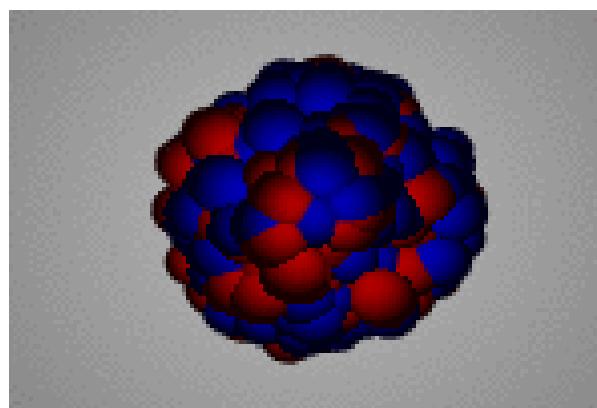
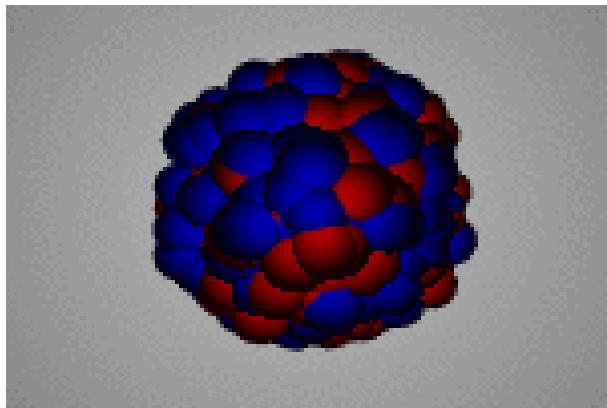
What happens if a photon is absorbed into a nucleus?
- one nucleon simply starts moving faster?



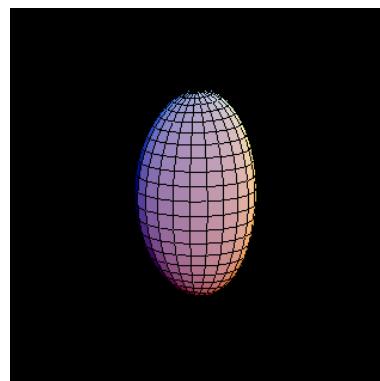
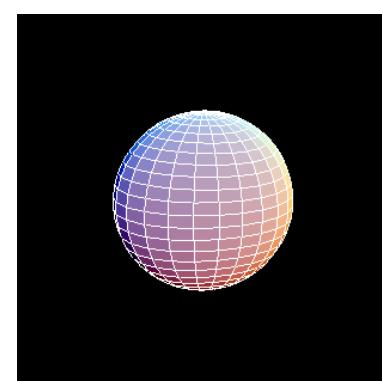
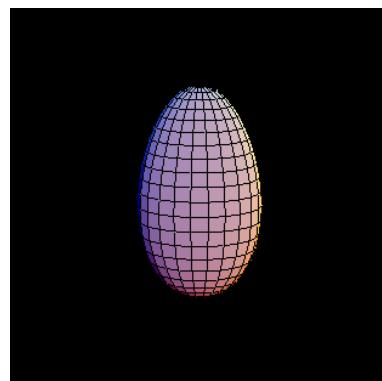
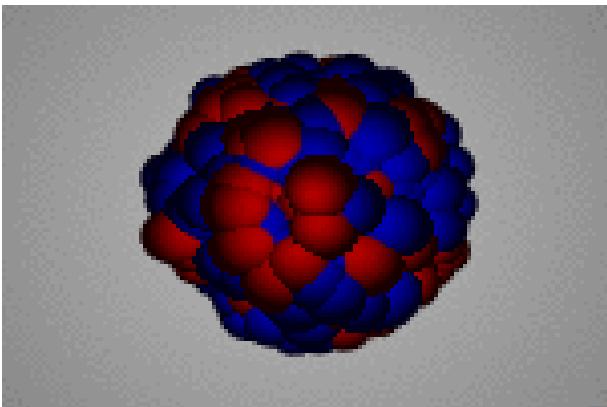
Very coherent
motion can happen
Collective motions



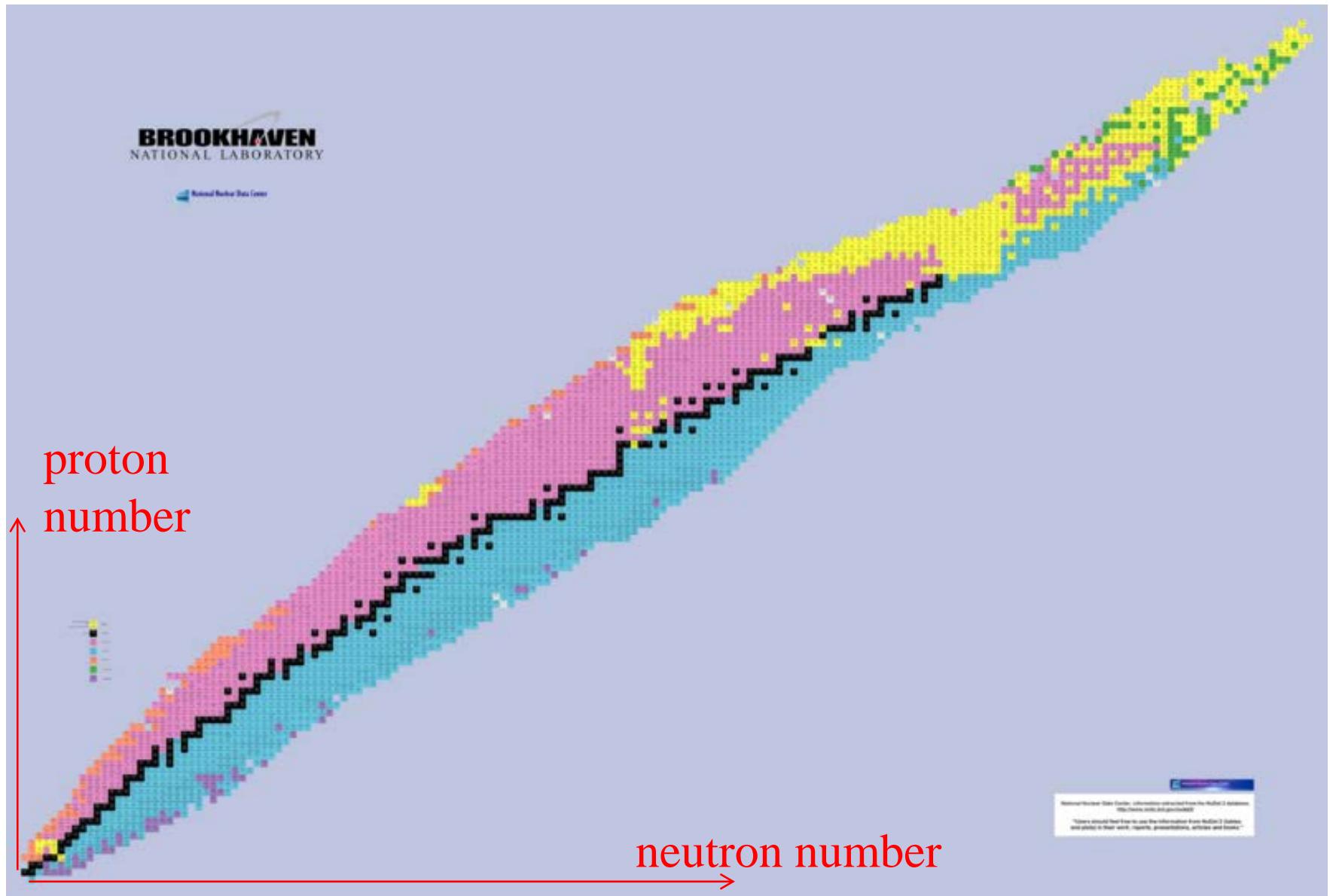
Very coherent
motion can happen
Collective motions



a variety of
motions
→ very rich!



Nuclear Chart: 2D map of atomic nuclei

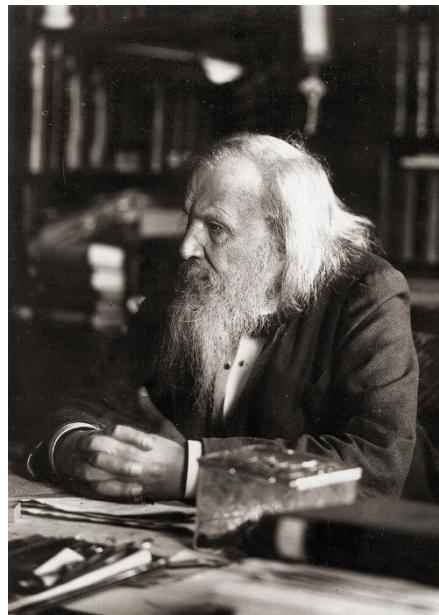


Periodic Table of elements (1869)

Group	→ 1	Period
↓		
1	1 H	2
2	3 Li 4 Be	
3	11 Na 12 Mg	
4	19 K 20 Ca 21 Sc 22 Ti 23 V 24 Cr 25 Mn 26 Fe 27 Co 28 Ni 29 Cu 30 Zn 31 Ga 32 Ge 33 As 34 Se 35 Br 36 Kr	
5	37 Rb 38 Sr 39 Y 40 Zr 41 Nb 42 Mo 43 Tc 44 Ru 45 Rh 46 Pd 47 Ag 48 Cd 49 In 50 Sn 51 Sb 52 Te 53 I 54 Xe	
6	55 Cs 56 Ba 57 La * 72 Hf 73 Ta 74 W 75 Re 76 Os 77 Ir 78 Pt 79 Au 80 Hg 81 Tl 82 Pb 83 Bi 84 Po 85 At 86 Rn	
7	87 Fr 88 Ra 89 Ac * 104 Rf 105 Db 106 Sg 107 Bh 108 Hs 109 Mt 110 Ds 111 Rg 112 Cn 113 Nh 114 Fl 115 Mc 116 Lv 117 Ts 118 Og	



International Year
of the Periodic Table
of Chemical Elements



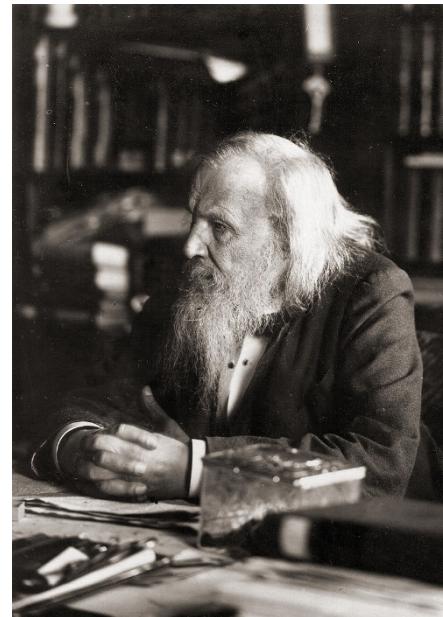
Mendeleev
(1834-1907)

13	14	15	16	17	18
5 B	6 C	7 N	8 O	9 F	10 Ne
13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn
115 Mc	116 Lv	117 Ts	118 Og		
* 58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu
* 90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am
* 96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md
					102 No
					103 Lr

Periodic Table of elements (1869)



5 6 7



Mendeleev
(1834-1907)

13 14 15 16 17 18

2 He

10 Ne

6 Ar

4 e

6 n

8 g

5 B

13 Al

6 C

14 Si

7 N

15 P

8 O

16 S

9 F

17 Cl

18 Ar

Taka and Fuji (a grand
-daughter of Mendeleev)

55 Cs	56 Ba	57 La	*	72 Hf	73 Ta	74 W
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87 Fr	88 Ra	89 Ac	*	104 Rf	105 Db	106 Sg
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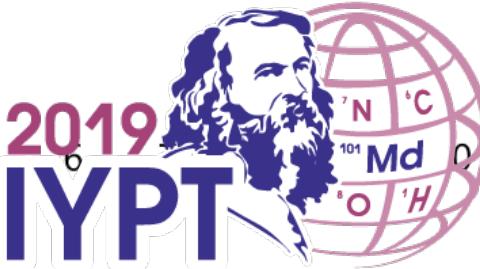
*	58 Ce	59 Pr	60 Nd
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*	90 Th	91 Pa	92 U
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Periodic Table of elements (1869)

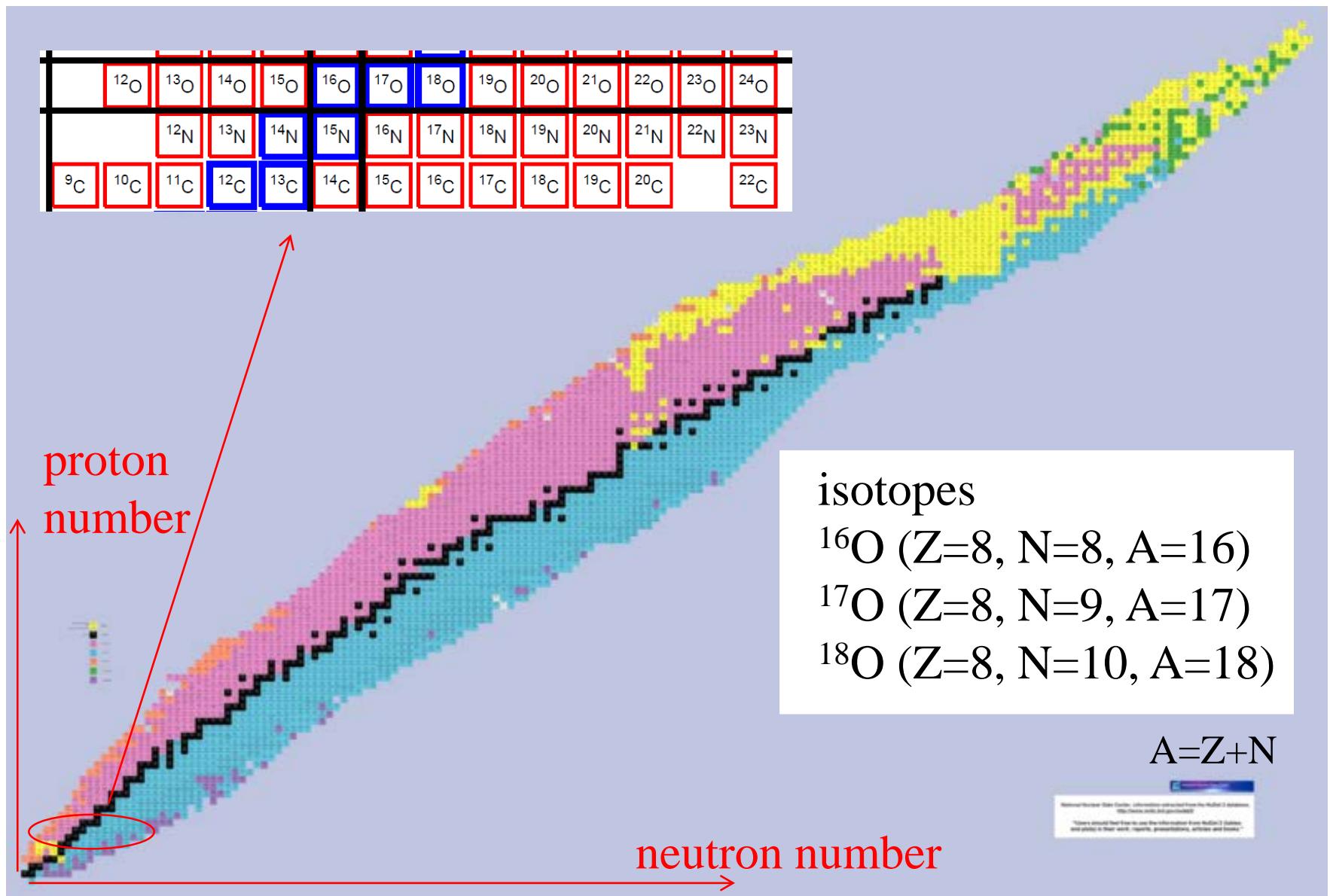
protons only, no information on neutrons



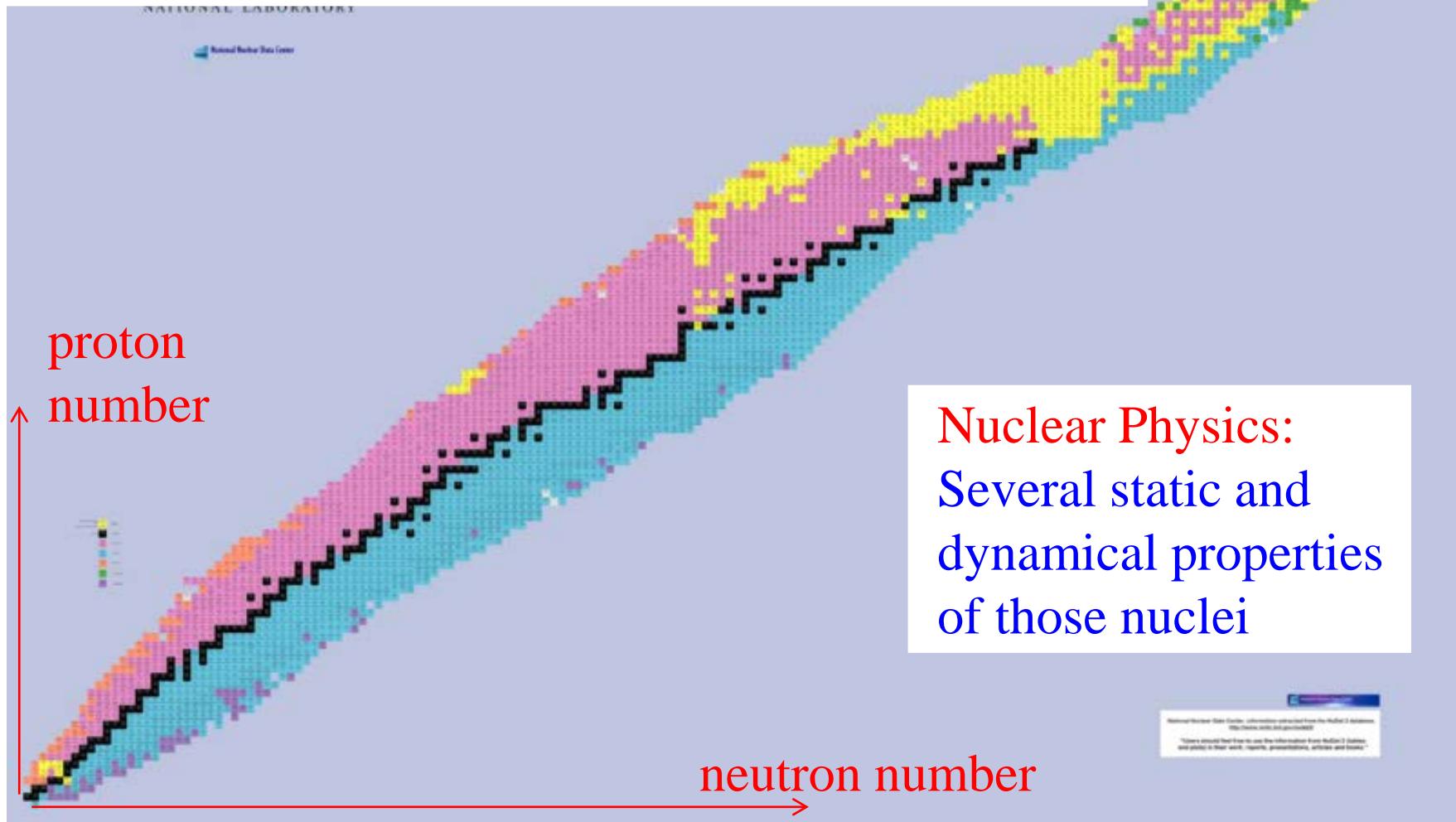
International Year
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of Chemical Elements

Group	→ 1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	
↓ Period																			
1	H																	He	
2	Li	Be																Ne	
3	Na	Mg																Ar	
4	K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr	
5	Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe	
6	Cs	Ba	La	*	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	At	Rn	
7	Fr	Ra	Ac	*	104	105	106	107	108	109	110	111	112	113	114	115	116	117	Og
				*	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb	Lu	
				*	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No	Lr	

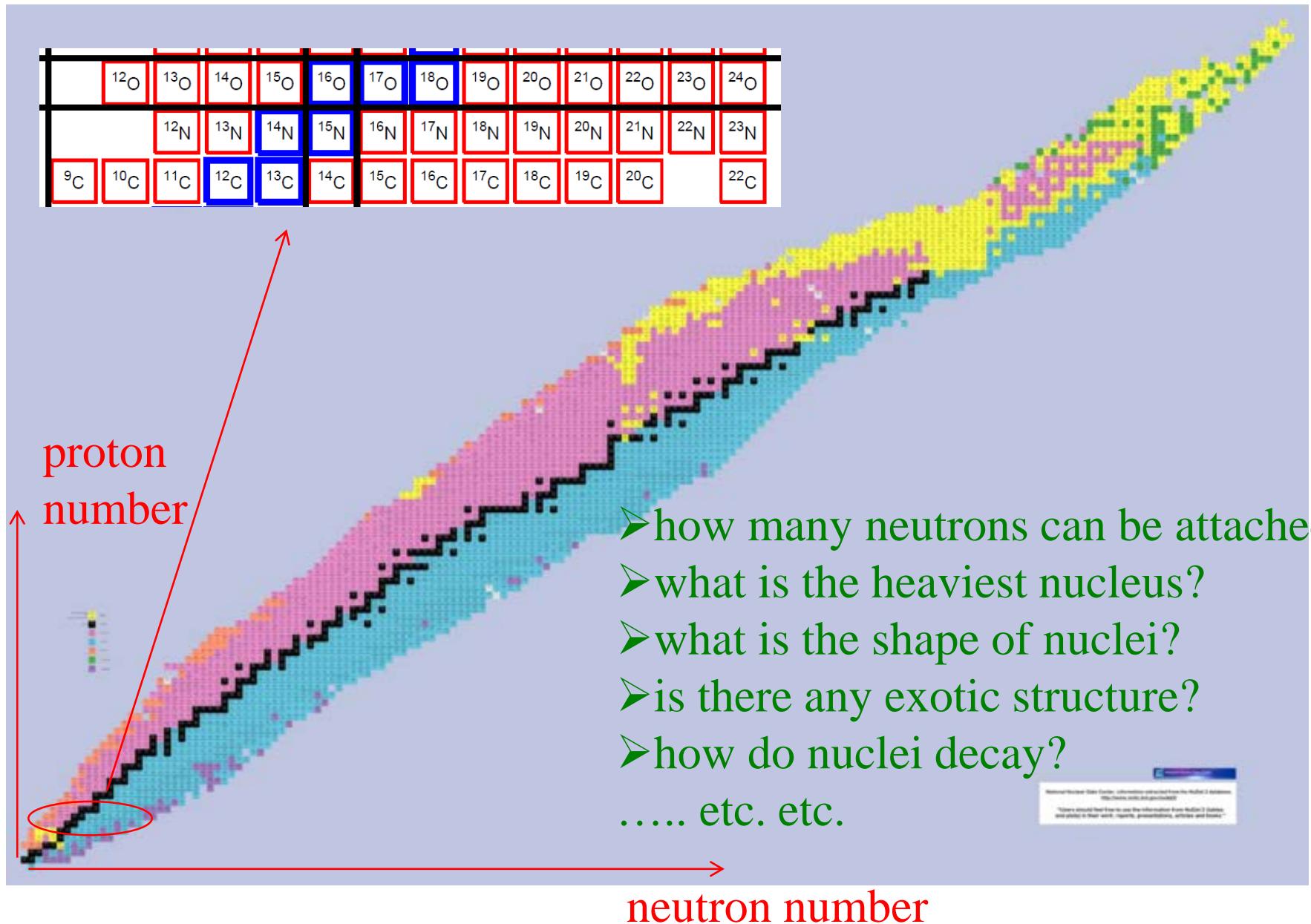
Nuclear Chart: an extended version of periodic table



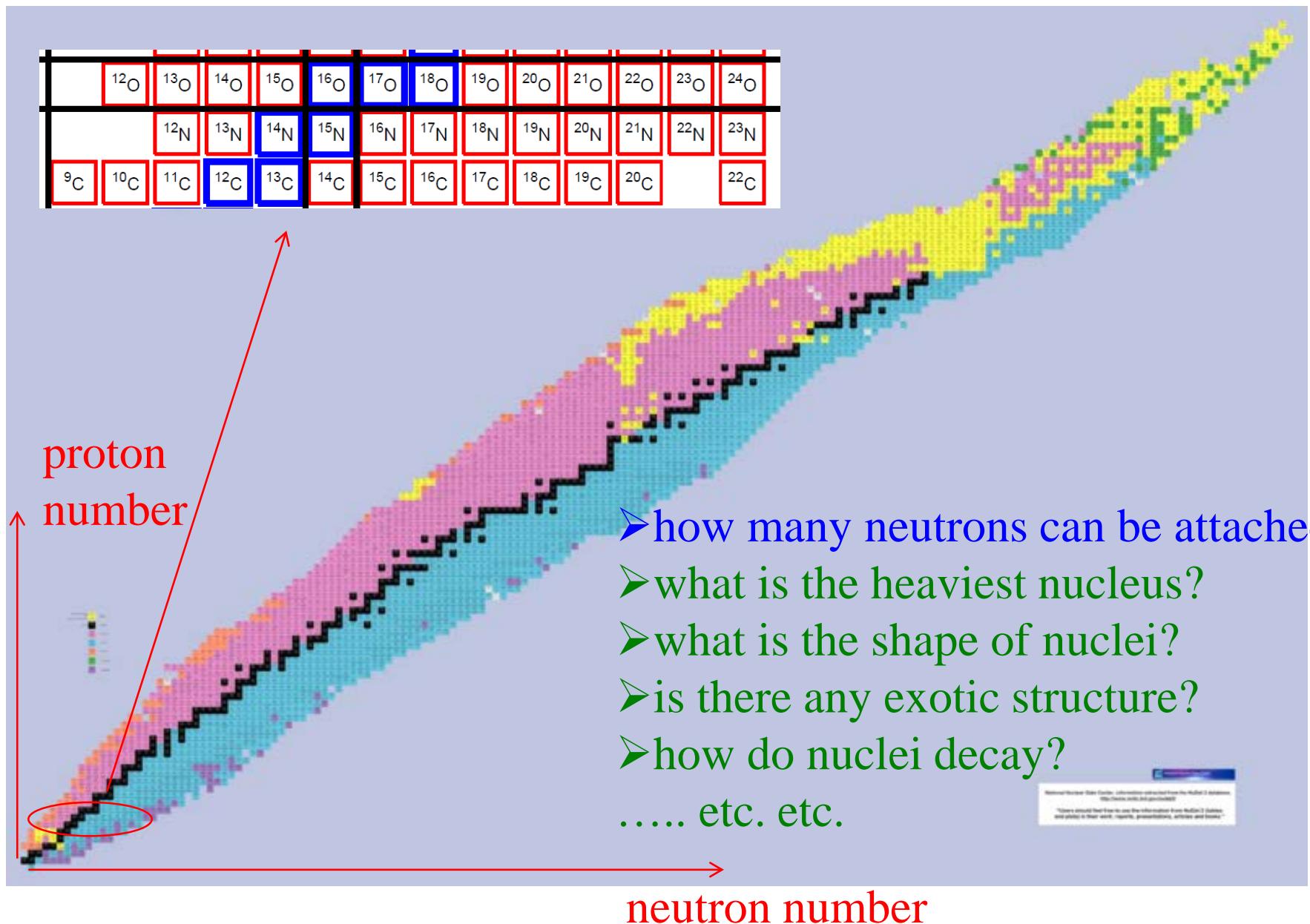
- Stable nuclei in nature : 287
- Nuclei artificially synthesized : about 3,000
- Nuclei predicted : about 7,000 ~ 10,000



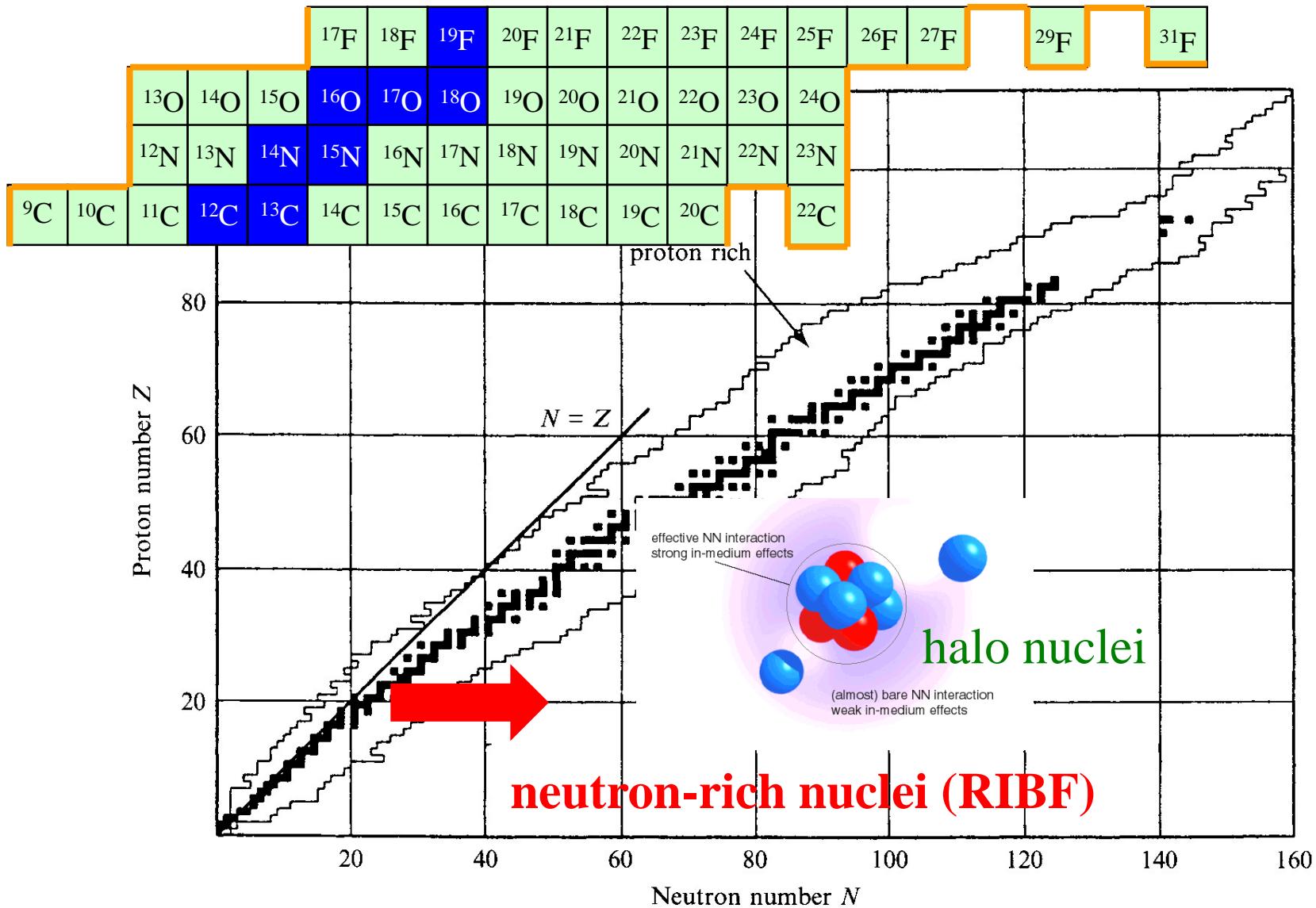
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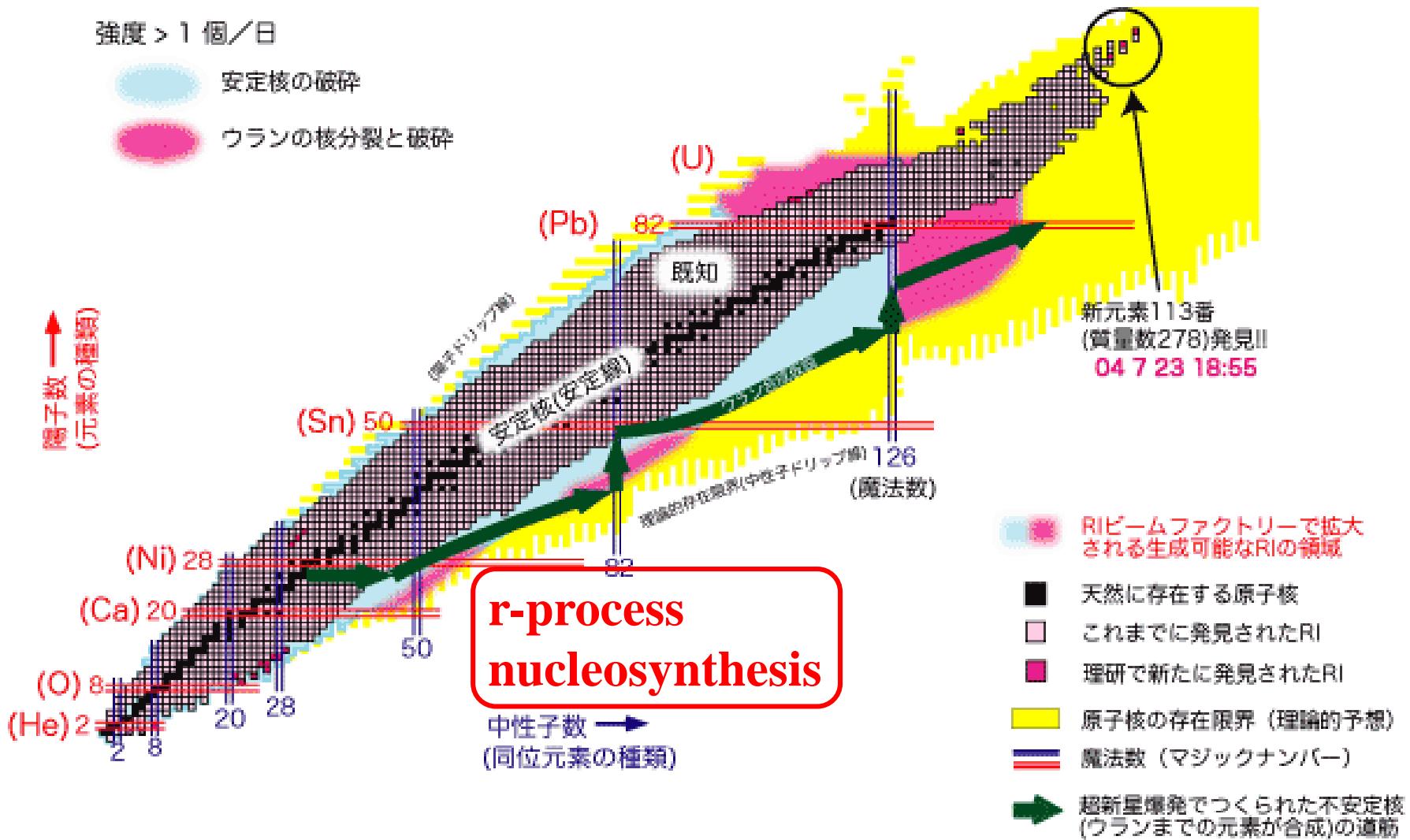
Nuclear Chart: an extended version of periodic table



Extension of nuclear chart: frontier of nuclear physics



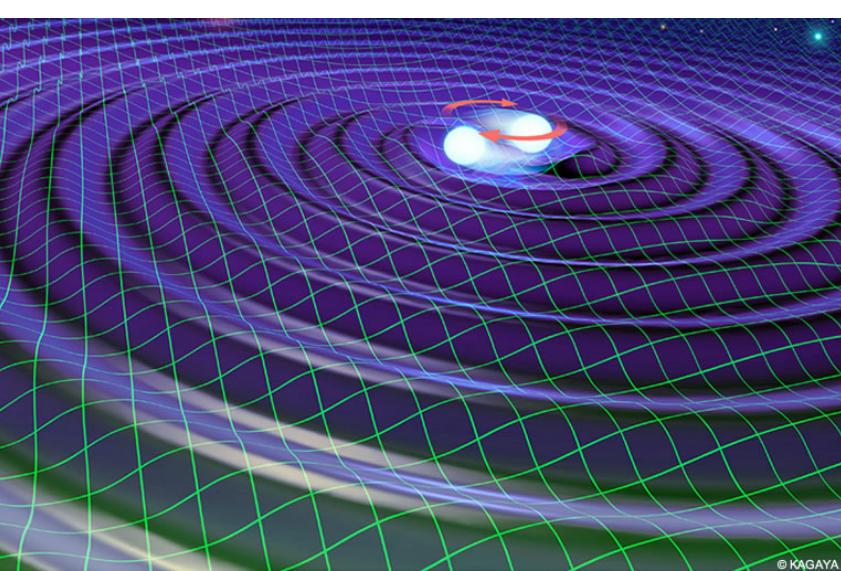
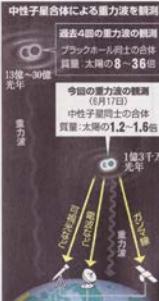
Neutron-rich nuclei (RIBF at RIKEN)



October 17, 2017

星の合体 重力波で観測

星の合体による重力波を観測
中性星子が大きくなり、星が合った
重力をもつて、星が大きくなる
星の合体による重力波を観測
星の合体による重力波を観測



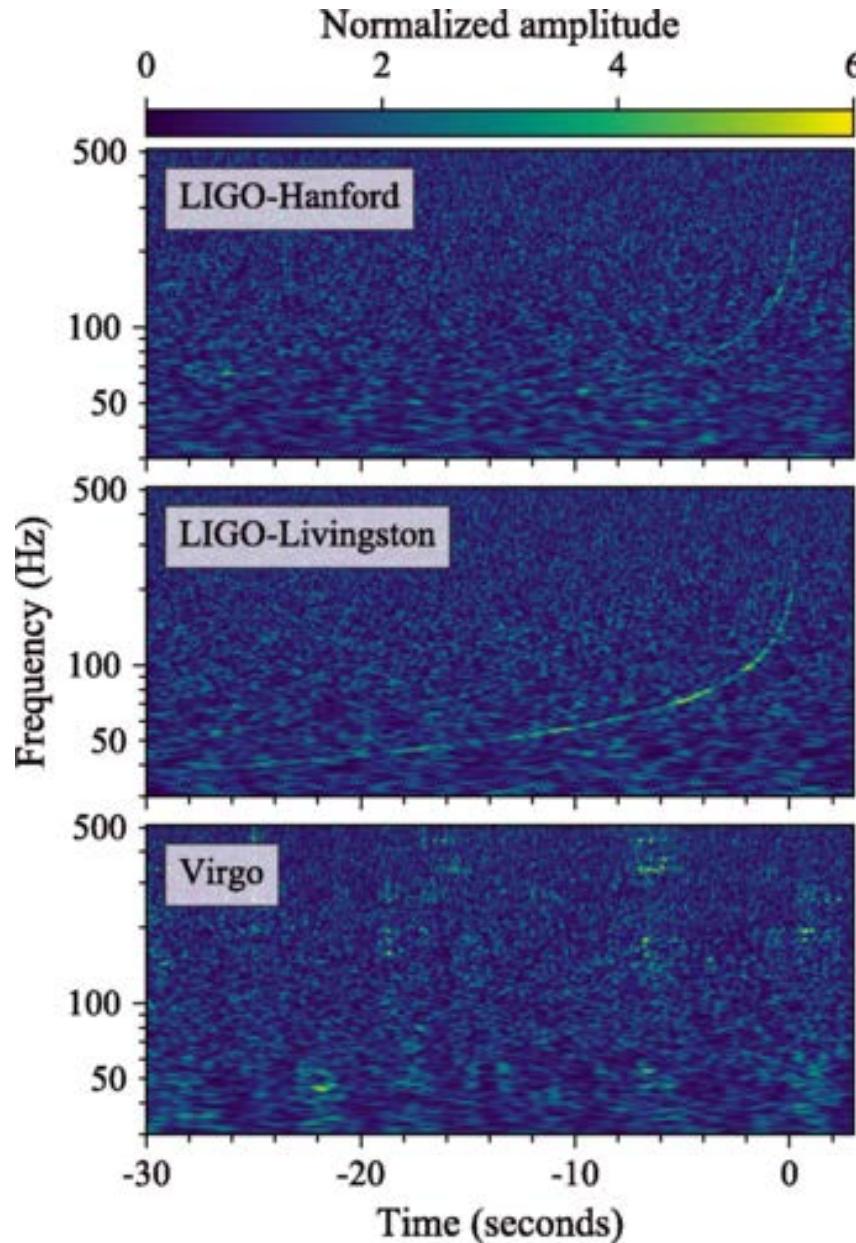
発生源からの光も確認

星の合体による重力波を観測
星の合体による重力波を観測
星の合体による重力波を観測
星の合体による重力波を観測



NAOJ

nucleosynthesis at this time?



B.P. Abbott et al., PRL119 ('17) 161101

October 17, 2017

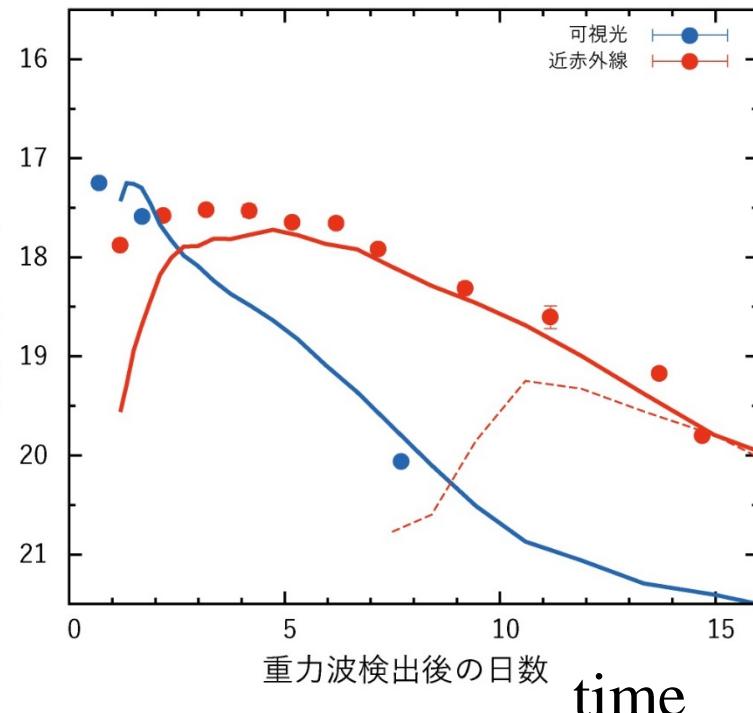
Ele.-mag. wave from the source of GW

星の合体 重力波で観測

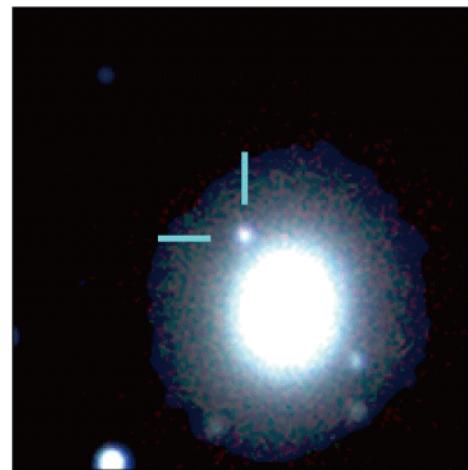


元素の起源
解明に期待

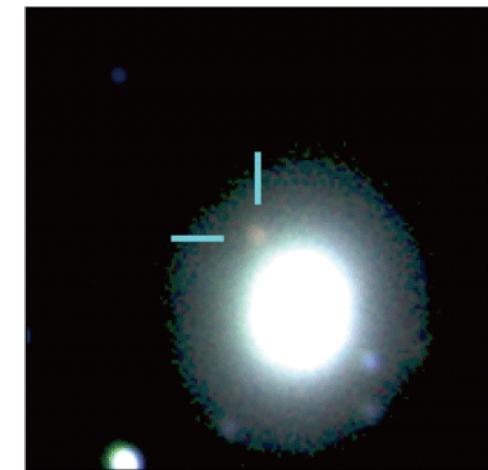
opacity
等級 (明るさ)



2017.08.18-19



2017.08.24-25

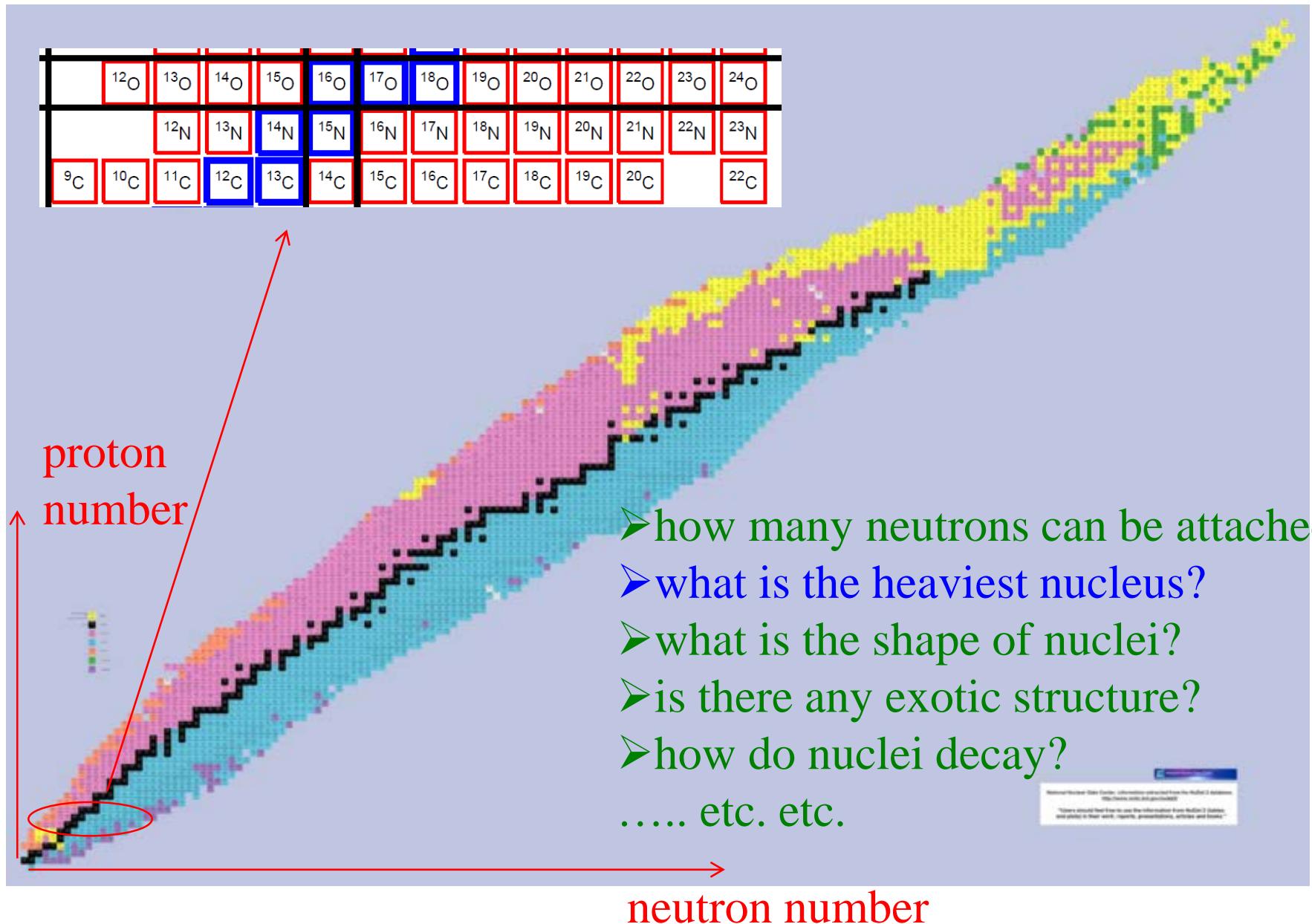


solid lines : with r-process nucleosynthesis
dashed line : no r-process nucleosynthesis

M. Tanaka et al.,
Astron. Soc. Jpn. 69 ('17) 102

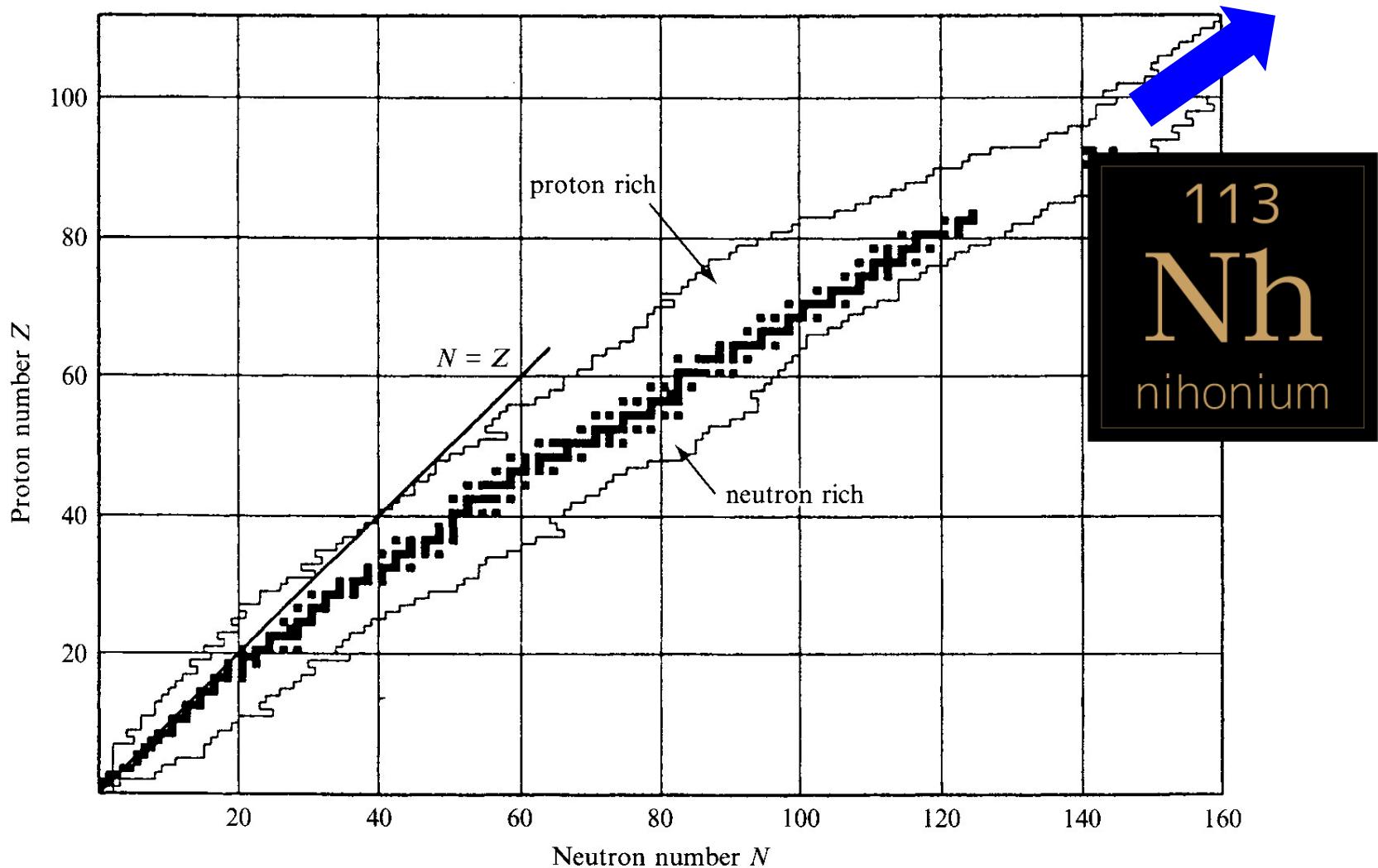
<http://www.cfca.nao.ac.jp/pr/20171016>

Nuclear Chart: an extended version of periodic table

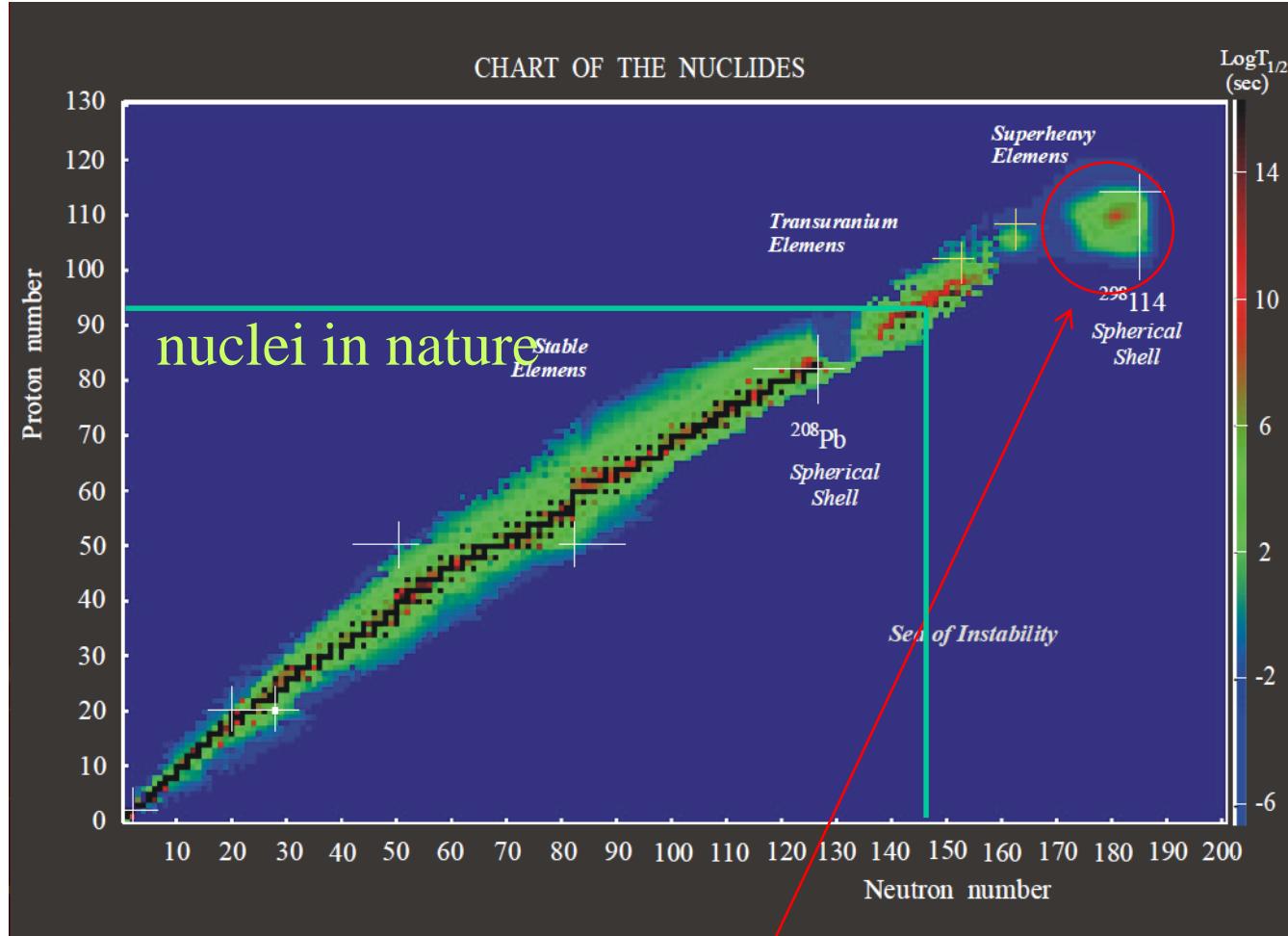


Extension of nuclear chart: frontier of nuclear physics

superheavy
elements



Prediction of island of stability: an important motivation of SHE study

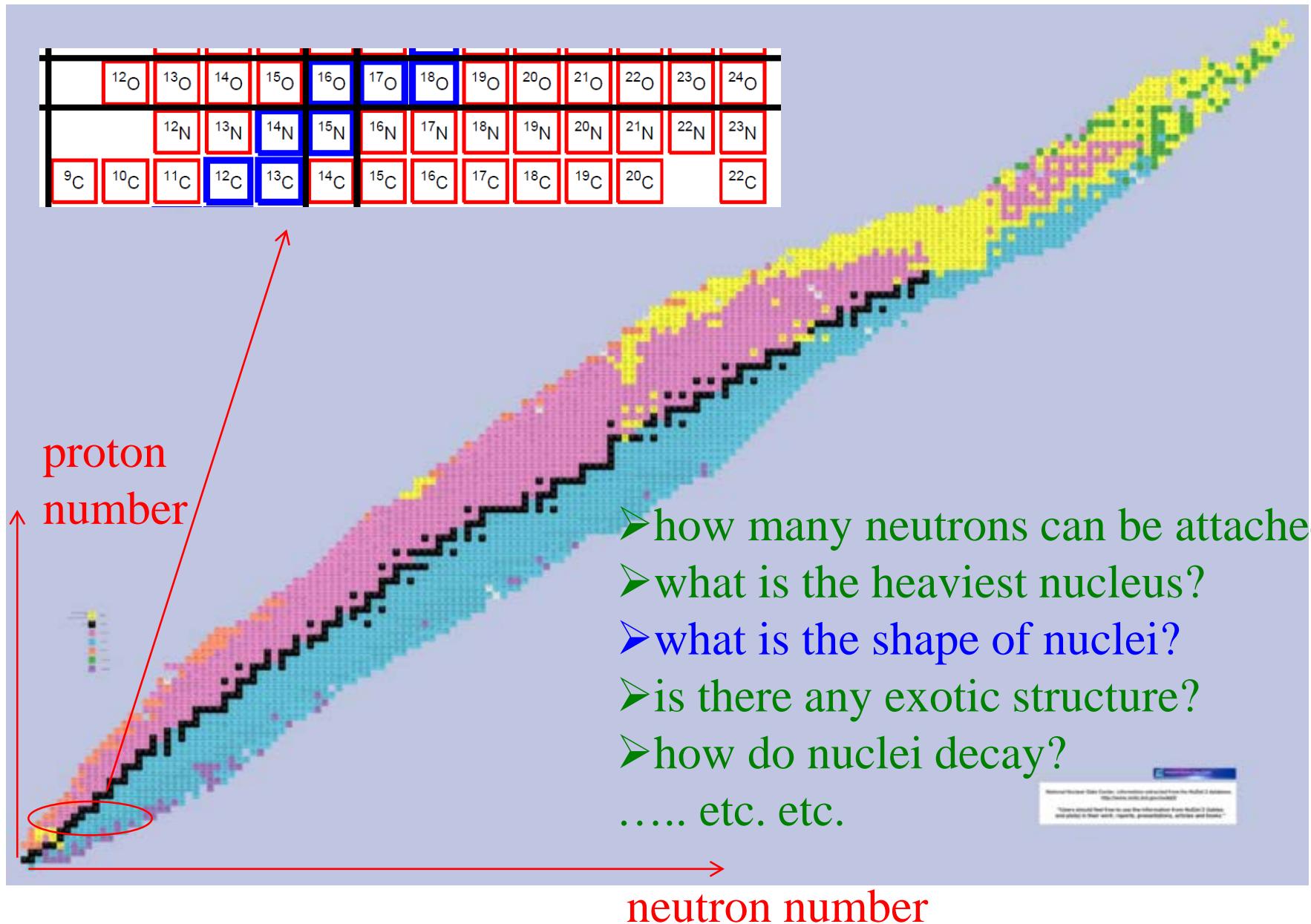


Yuri Oganessian

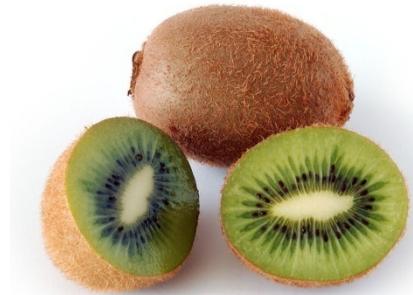
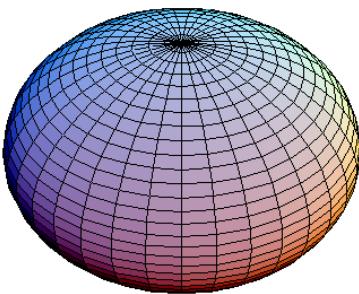
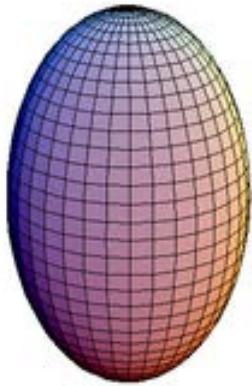
island of stability around $Z=114, N=184$

W.D. Myers and W.J. Swiatecki (1966), A. Sobiczewski et al. (1966)

Nuclear Chart: an extended version of periodic table



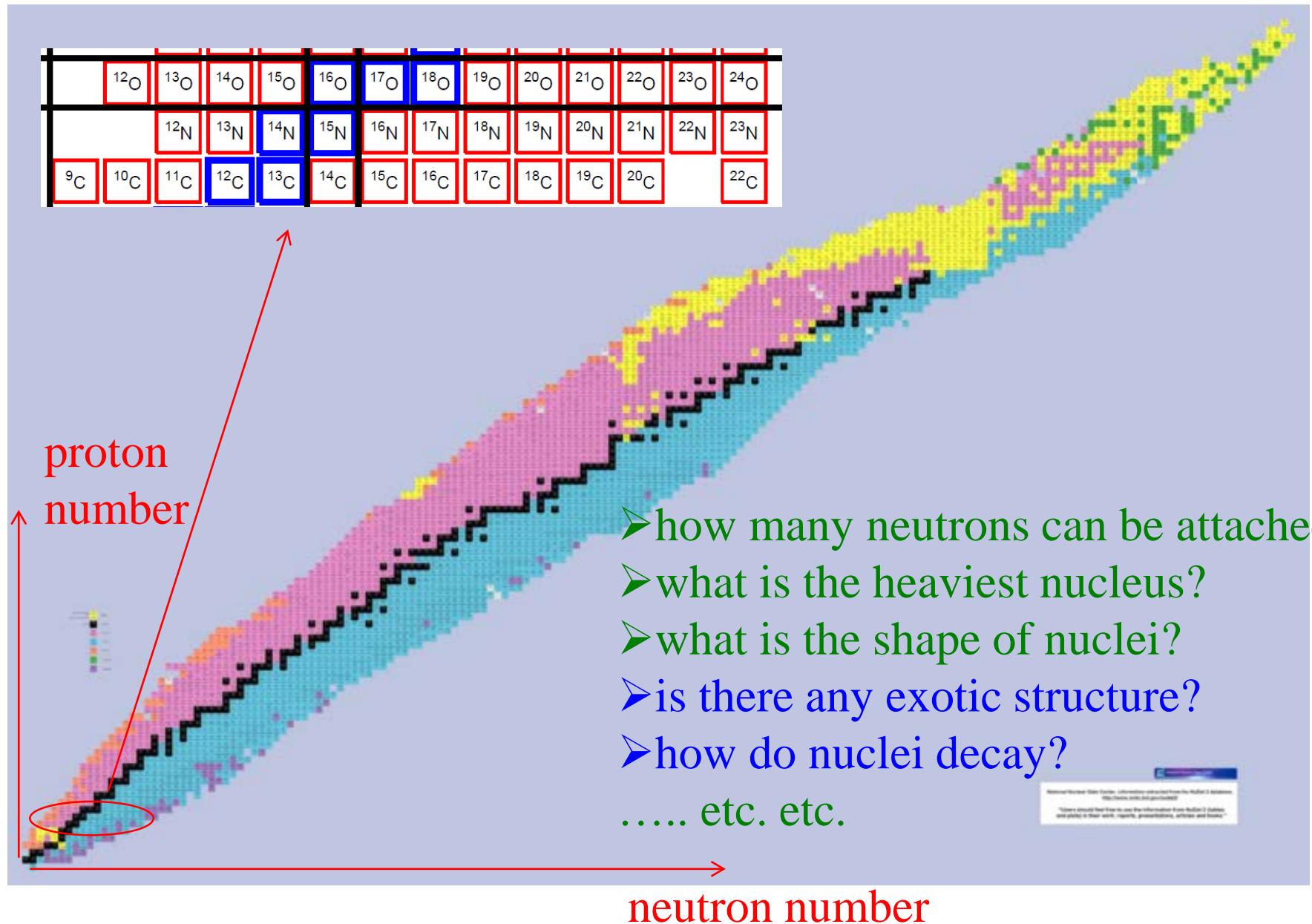
a nucleus is not always spherical



Quantum shape
dynamics

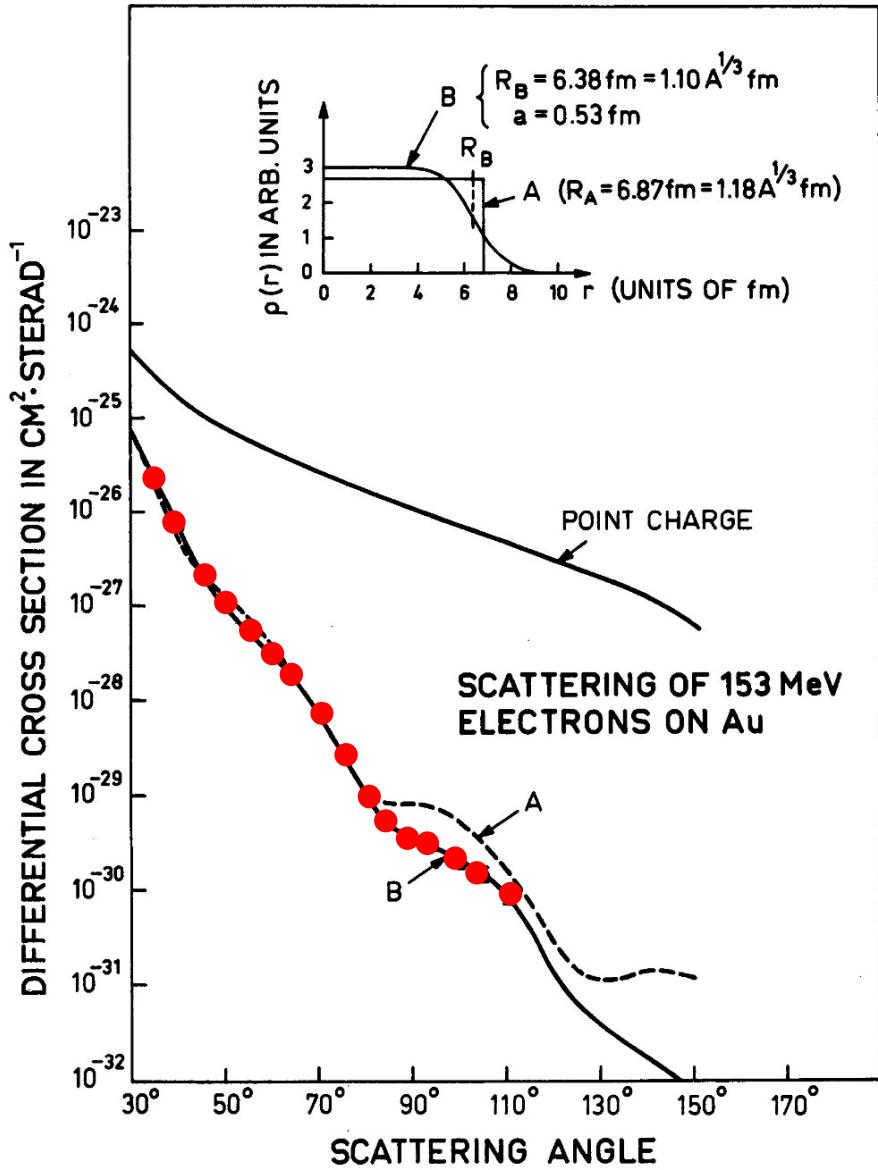
「形の量子論」

Nuclear Chart: an extended version of periodic table



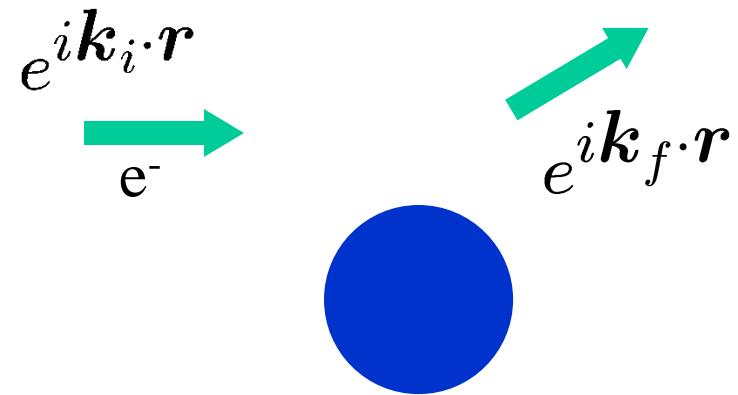
Density Distribution

Density Distribution



High energy electron scattering

Born approximation:



$$\frac{d\sigma}{d\Omega} = \frac{Z_P^2 e^4}{(4E \sin^2 \theta/2)^2} |F(q)|^2$$

Form factor

$$F(q) = \int e^{-iq \cdot r} \rho(r) dr$$

(Fourier transform of the density)

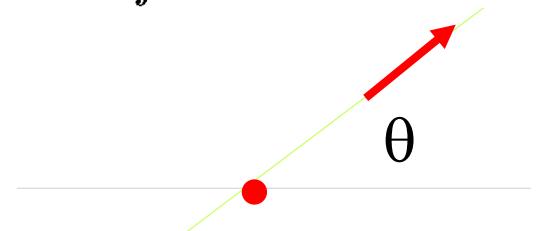
Born approximation

$$\psi_f(r) = e^{i\mathbf{p}_f \cdot \mathbf{r} / \hbar}$$

$$\psi_i(r) = e^{i\mathbf{p}_i \cdot \mathbf{r} / \hbar}$$



$$V(r)$$

$$W_{fi} = \frac{2\pi}{\hbar} |\langle \psi_f | V | \psi_i \rangle|^2 \rho(E_f)$$

incident flux: $j_{\text{inc}} = \rho_i v = p_i / \mu$



$$\sigma = \frac{W_{fi}}{j_{\text{inc}}} = \int d\Omega \boxed{\frac{d\sigma}{d\Omega}}$$

Electron scattering

$$V(r) = -e^2 \int d\mathbf{r}' \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

$$\begin{aligned} \frac{d\sigma}{d\Omega} &= \frac{Z_P^2 e^4}{(4E \sin^2 \theta/2)^2} |F(\mathbf{q})|^2 \\ &= \left(\frac{d\sigma_{\text{Ruth}}}{d\Omega} \right) |F(\mathbf{q})|^2 \end{aligned}$$

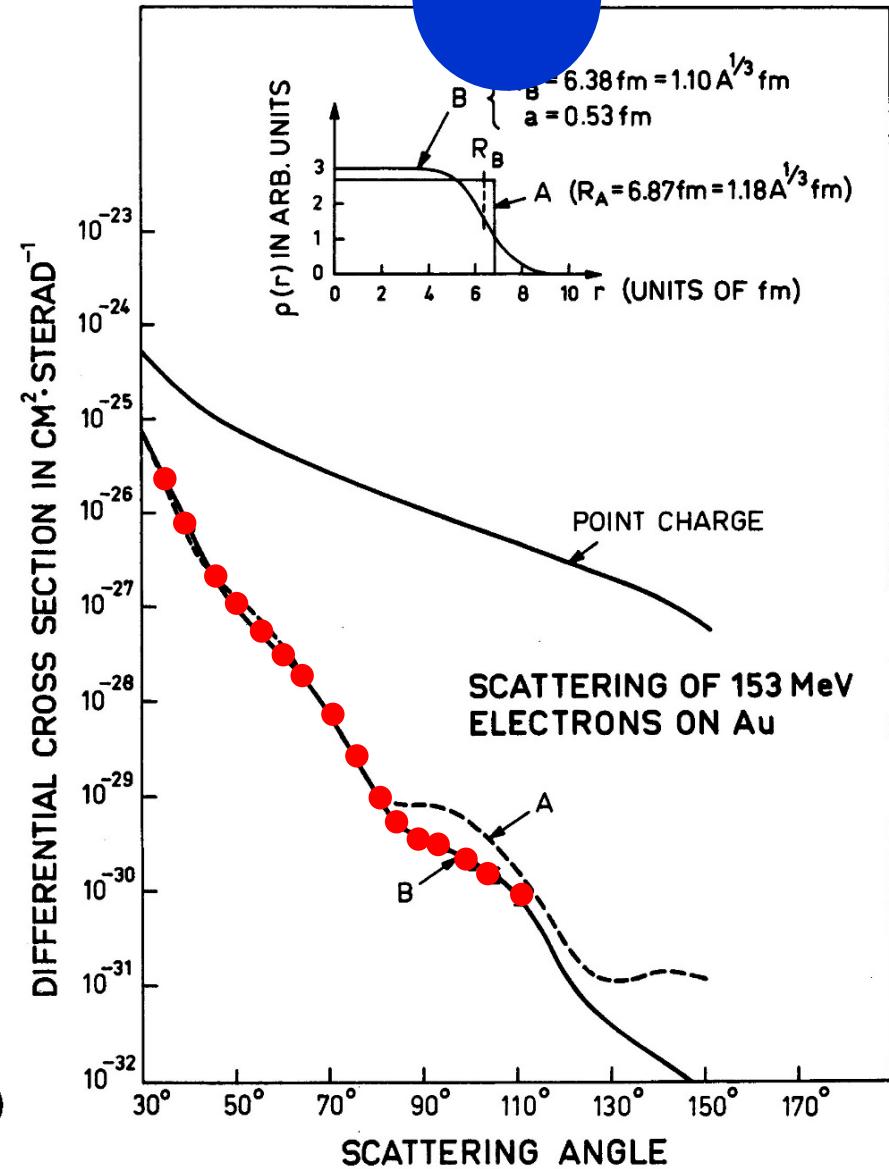


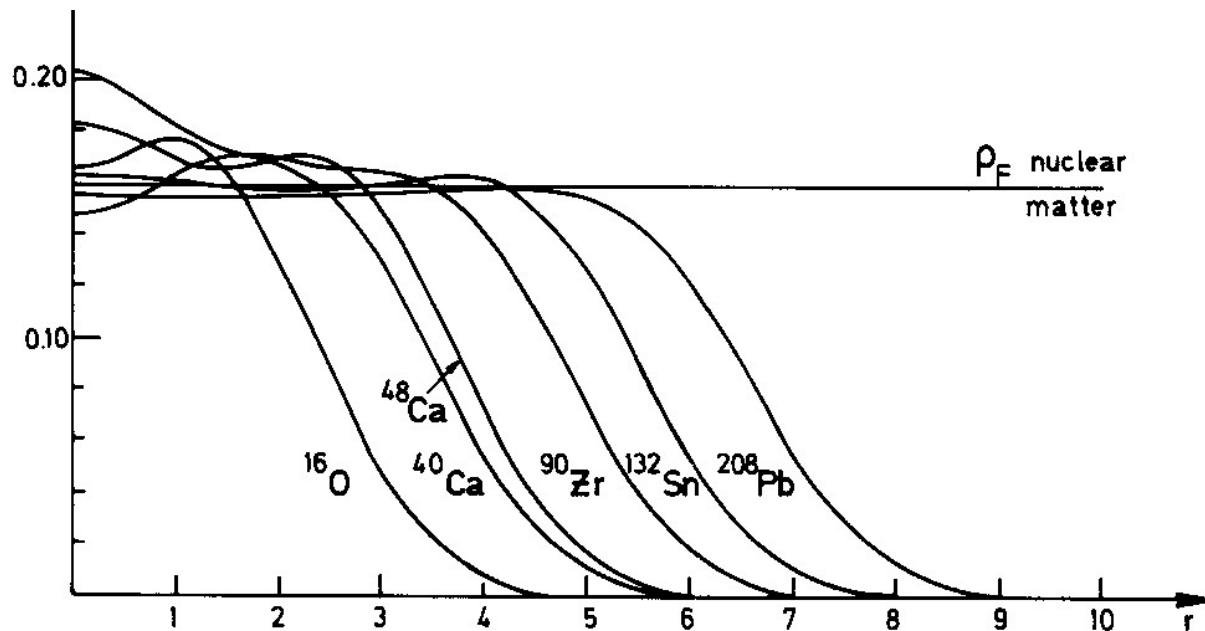
Form factor

$$F(\mathbf{q}) = \int e^{-i\mathbf{q} \cdot \mathbf{r}} \rho(\mathbf{r}) d\mathbf{r}$$

* relativistic correction:

$$\begin{aligned} \frac{d\sigma_{\text{Ruth}}}{d\Omega} &\rightarrow \frac{d\sigma_{\text{Mott}}}{d\Omega} \\ &= \frac{d\sigma_{\text{Ruth}}}{d\Omega} \cdot \left(1 - \frac{v^2}{c^2} \sin^2 \frac{\theta}{2} \right) \\ &\sim \frac{d\sigma_{\text{Ruth}}}{d\Omega} \cdot \cos^2 \frac{\theta}{2} \quad (v \rightarrow c) \end{aligned}$$





Fermi distribution

$$\rho(r) = \frac{\rho_0}{1 + \exp[(r - R_0)/a]}$$

$$\rho_0 \sim 0.17 \text{ (fm}^{-3}\text{)}$$

$$R_0 \sim 1.1 \times A^{1/3} \text{ (fm)}$$

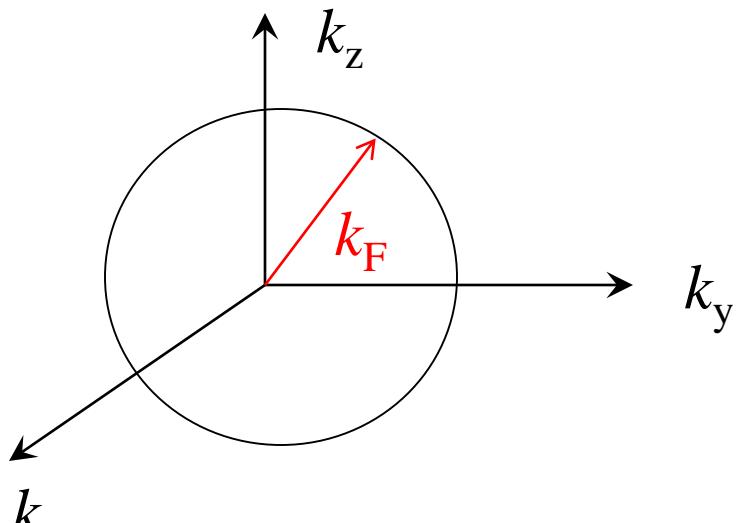
$$a \sim 0.57 \text{ (fm)}$$



Saturation
property

Momentum Distribution

Fermi gas approximation



$$\begin{aligned}\rho &= 2 \times 2 \times 4\pi \int_0^{k_F} \frac{k^2 dk}{(2\pi)^3} \\ &= \frac{2}{3\pi^2} k_F^3\end{aligned}$$

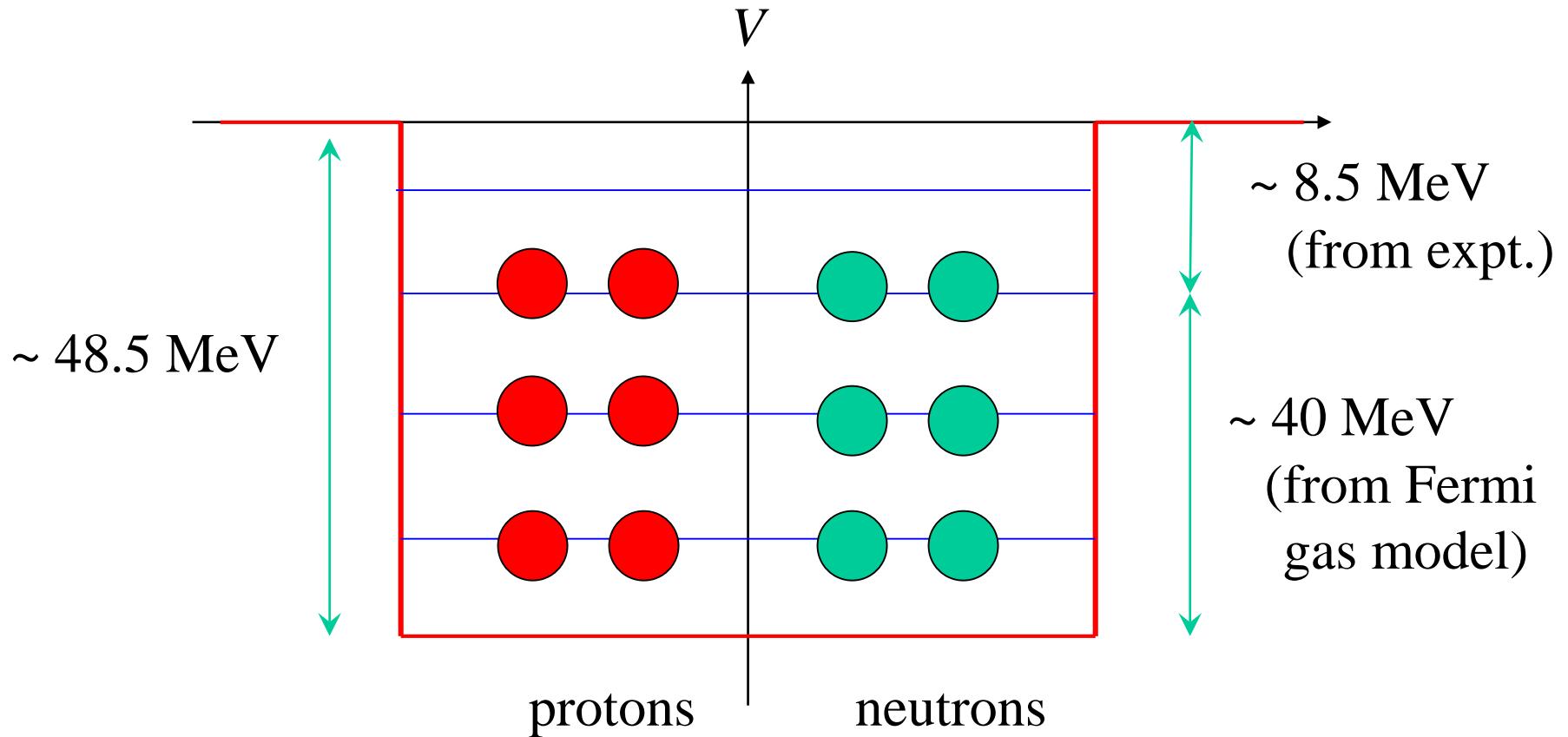
(note: spin-isospin degeneracy)

$$\rho = 0.17 \text{ fm}^{-3} \longrightarrow k_F \sim 1.36 \text{ fm}^{-1}$$

$$\longleftrightarrow \frac{v_F}{c} = \frac{k_F \cdot \hbar c}{mc^2} = 0.285$$

Fermi energy: $\epsilon_F = \frac{k_F^2 \hbar^2}{2m} \sim 37 \text{ (MeV)}$

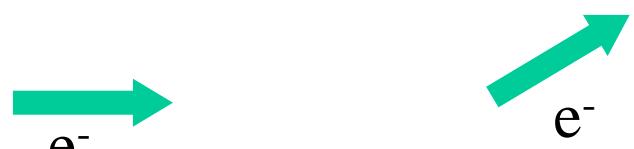
A potential for nucleons inside a nucleus



Discussion: Electron scattering

$$V(r) = -e^2 \int d\mathbf{r}' \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

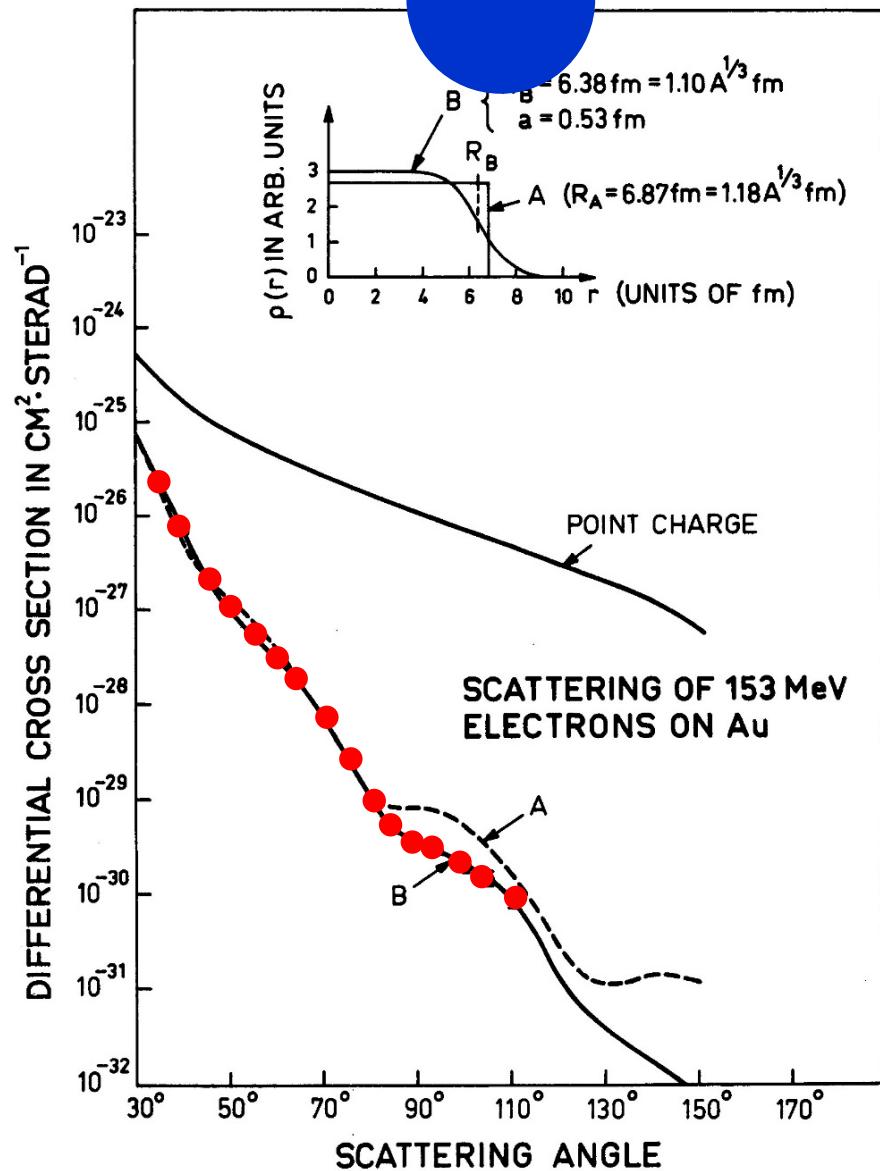
$$\begin{aligned}\frac{d\sigma}{d\Omega} &= \frac{Z_P^2 e^4}{(4E \sin^2 \theta/2)^2} |F(\mathbf{q})|^2 \\ &= \left(\frac{d\sigma_{\text{Ruth}}}{d\Omega} \right) |F(\mathbf{q})|^2\end{aligned}$$



Form factor

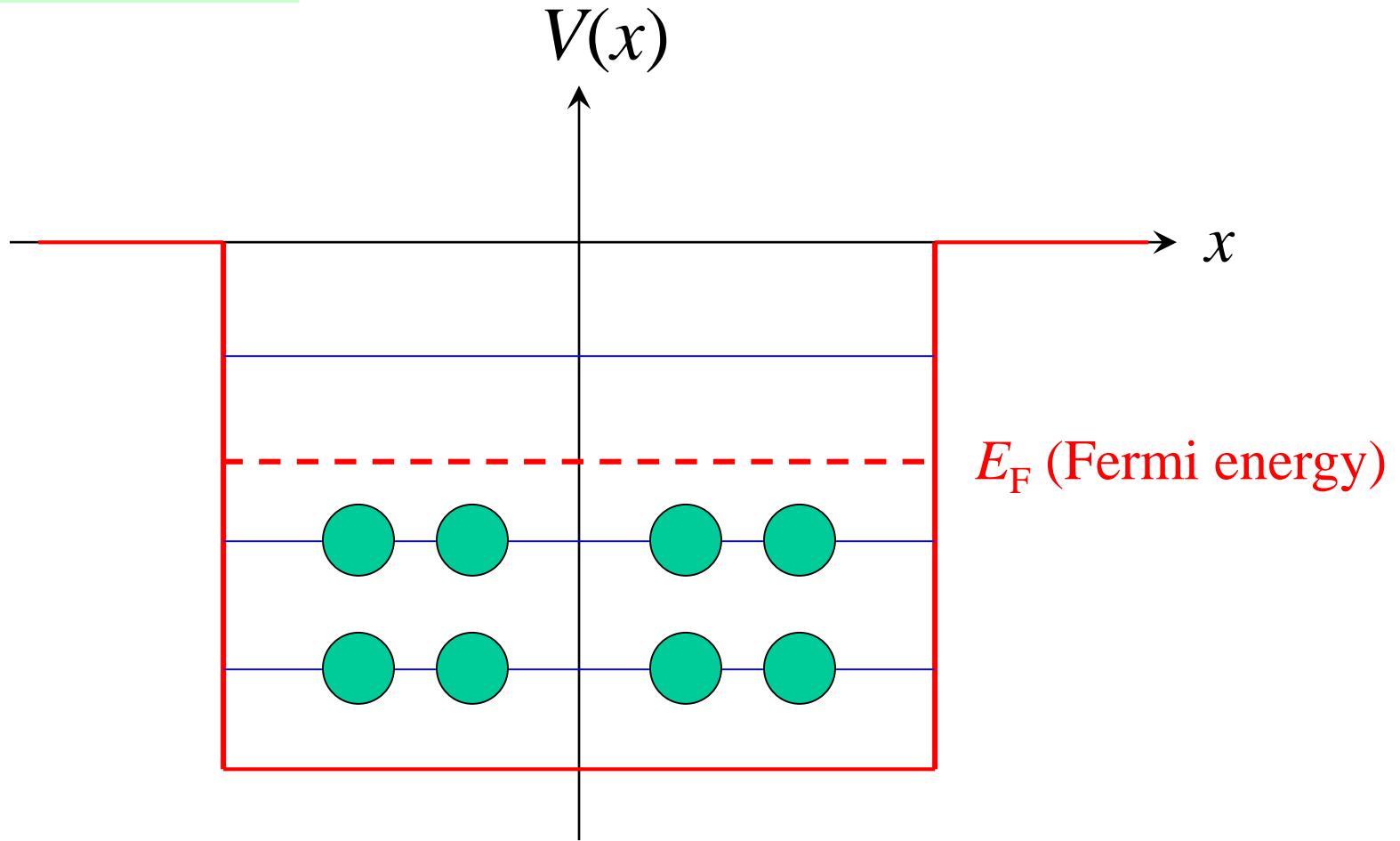
$$F(\mathbf{q}) = \int e^{-i\mathbf{q} \cdot \mathbf{r}} \rho(\mathbf{r}) d\mathbf{r}$$

Why do cross sections decrease for an extended density distribution?



Appendix

Fermi gas model

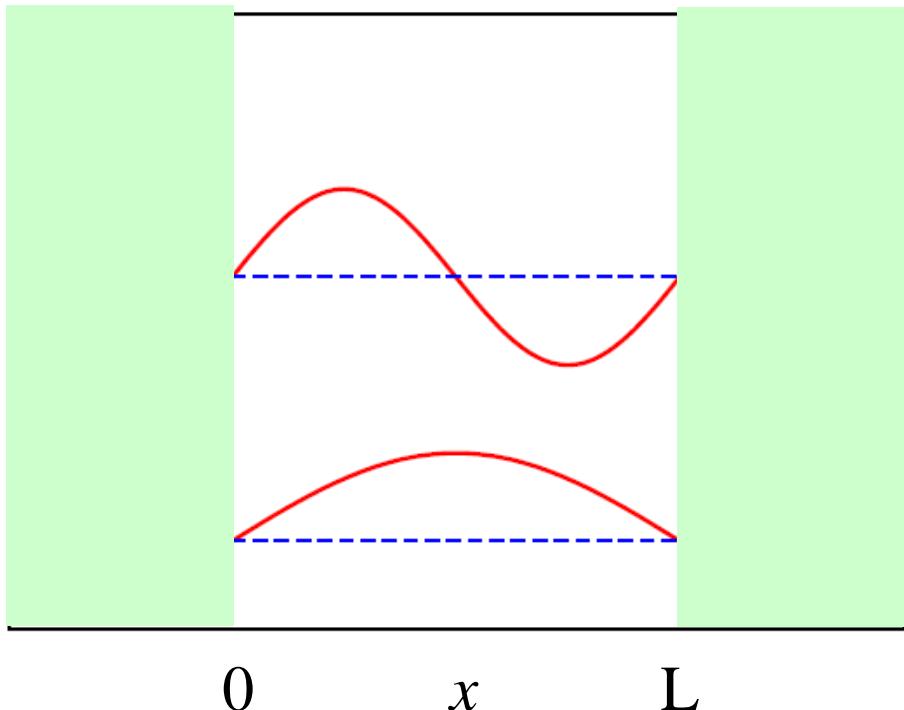


What is the relation between E_F and the particle number?

→ Fermi gas model

Fermi gas model

non-interaction many Fermion system (with no external potential)



put infinite walls at $x = 0$ and $x = L$:

$$\rightarrow \psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi}{L}x\right)$$

$$E_n = \frac{\hbar^2}{2m} \left(\frac{n\pi}{L}\right)^2 \quad (n = 1, 2, \dots)$$

three-dimensional case:

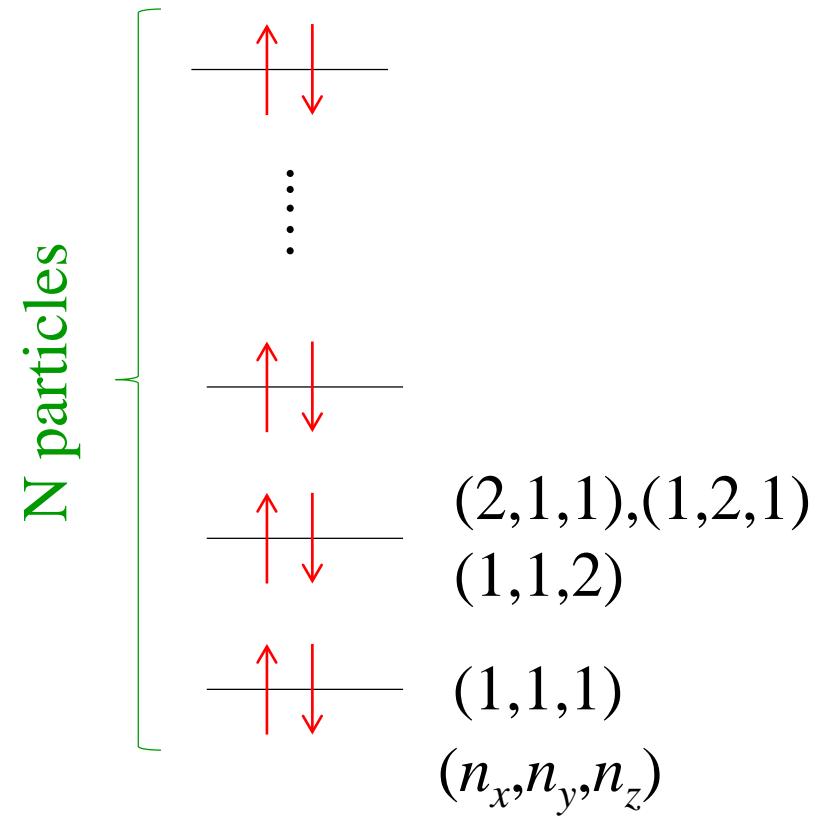
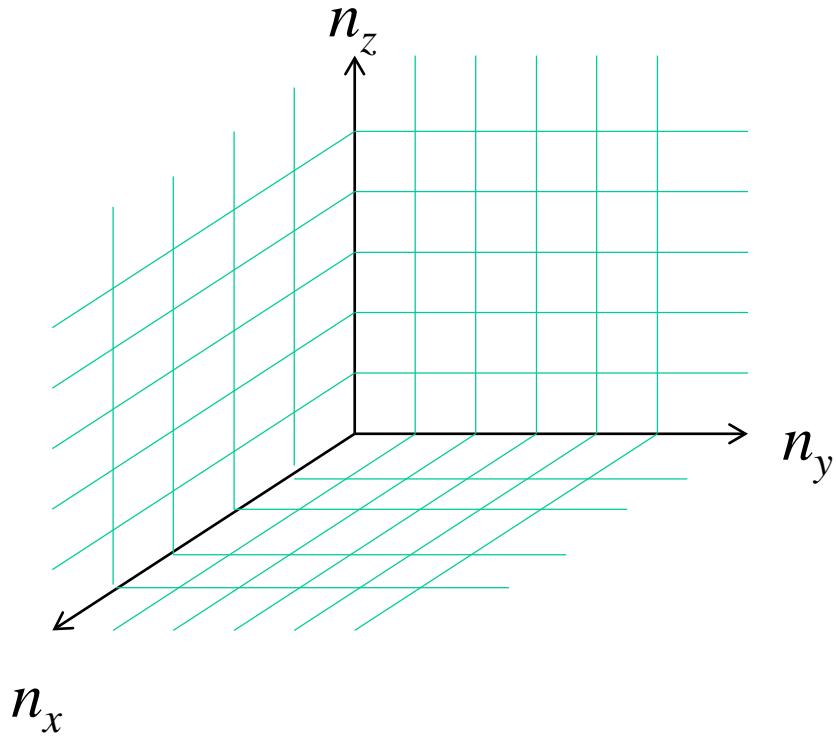
$$\psi_{n_x n_y n_z}(x) = \left(\frac{2}{L}\right)^{3/2} \sin\left(\frac{n_x \pi}{L}x\right) \sin\left(\frac{n_y \pi}{L}y\right) \sin\left(\frac{n_z \pi}{L}z\right)$$

$$E_{n_x n_y n_z} = \frac{\hbar^2}{2m} \left(\frac{\pi}{L}\right)^2 (n_x^2 + n_y^2 + n_z^2)$$

Fermi gas model

$$\psi_{n_x n_y n_z}(x) = \left(\frac{2}{L}\right)^{3/2} \sin\left(\frac{n_x \pi}{L} x\right) \sin\left(\frac{n_y \pi}{L} y\right) \sin\left(\frac{n_z \pi}{L} z\right)$$

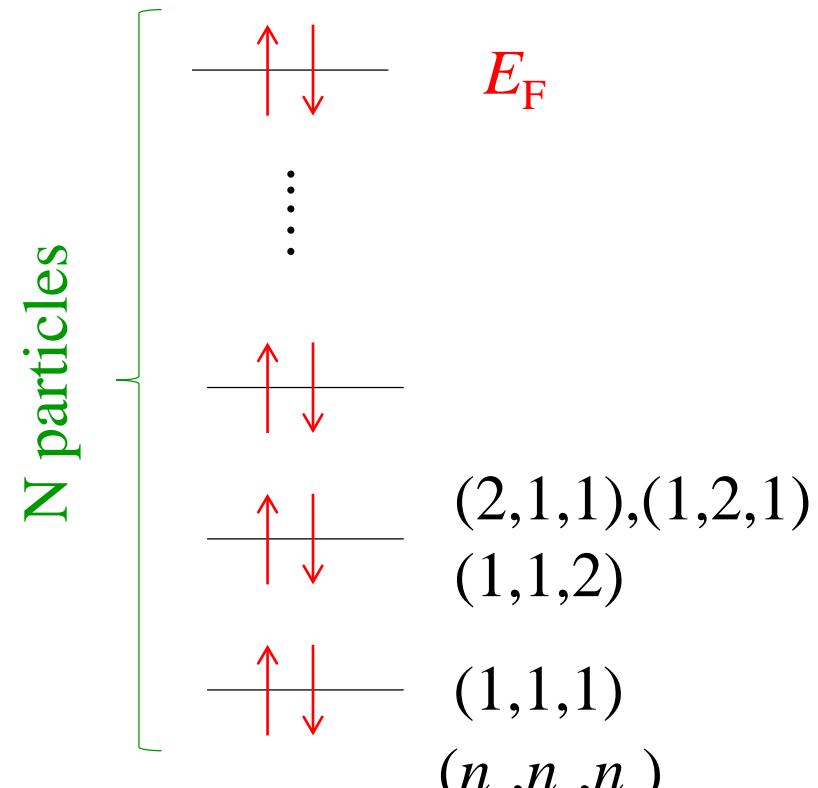
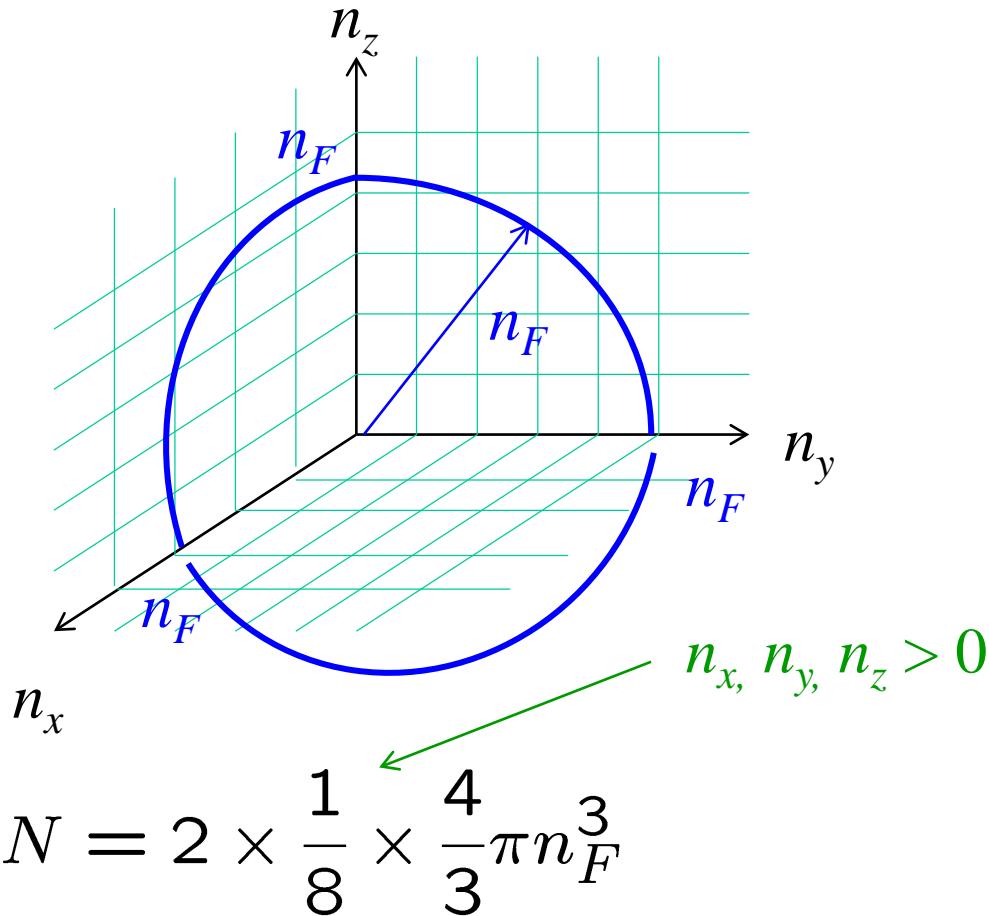
$$E_{n_x n_y n_z} = \frac{\hbar^2}{2m} \left(\frac{\pi}{L}\right)^2 (n_x^2 + n_y^2 + n_z^2)$$



Fermi gas model

the highest energy: $E_F = \frac{\hbar^2}{2m} \left(\frac{\pi}{L}\right)^2 (n_x^2 + n_y^2 + n_z^2) \equiv \frac{\hbar^2}{2m} \left(\frac{\pi}{L}\right)^2 n_F^2$

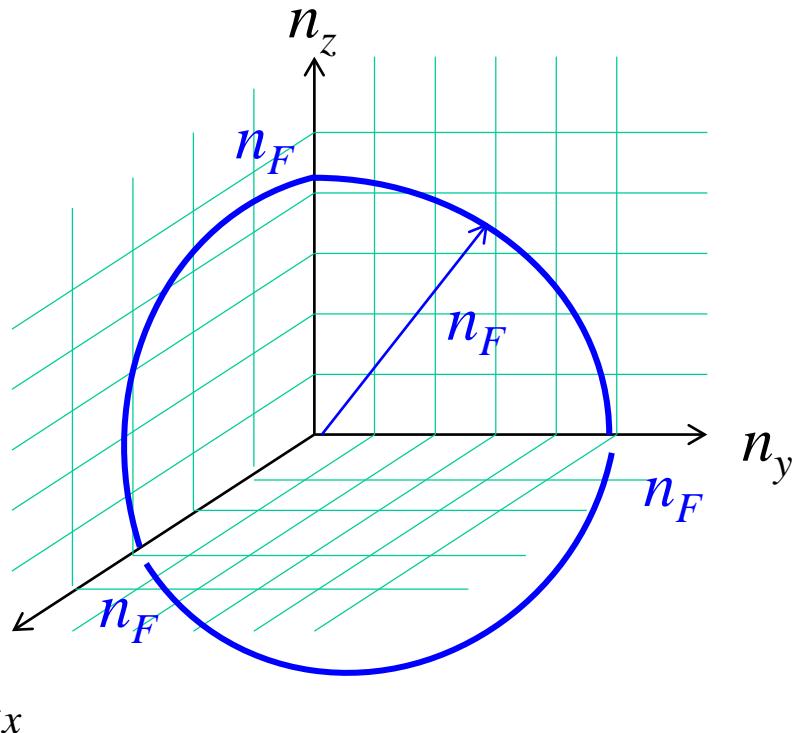
$$\rightarrow n_F^2 = \frac{2mE_F}{\hbar^2\pi^2} L^2$$



Fermi gas model

the highest energy: $E_F = \frac{\hbar^2}{2m} \left(\frac{\pi}{L}\right)^2 (n_x^2 + n_y^2 + n_z^2) \equiv \frac{\hbar^2}{2m} \left(\frac{\pi}{L}\right)^2 n_F^2$

$$\rightarrow n_F^2 = \frac{2mE_F}{\hbar^2\pi^2} L^2$$



$$N = 2 \times \frac{1}{8} \times \frac{4}{3} \pi n_F^3 = \frac{\pi}{3} \left(\frac{2mE_F}{\hbar^2\pi^2} L^2 \right)^{3/2} = \frac{\pi}{3} \left(\frac{2mE_F}{\hbar^2\pi^2} \right)^{3/2} L^3$$

Fermi gas model

$$N = 2 \times \frac{1}{8} \times \frac{4}{3} \pi n_F^3 = \frac{\pi}{3} \left(\frac{2mE_F}{\hbar^2 \pi^2} L^2 \right)^{3/2} = \frac{\pi}{3} \left(\frac{2mE_F}{\hbar^2 \pi^2} \right)^{3/2} L^3$$

→

$$\rho = \frac{N}{V} = \frac{N}{L^3} = \frac{\pi}{3} \left(\frac{2mE_F}{\hbar^2 \pi^2} \right)^{3/2}$$

or

$$E_F = \frac{\pi^2 \hbar^2}{2m} \left(\frac{3}{\pi} \rho \right)^{2/3}$$

Fermi gas model

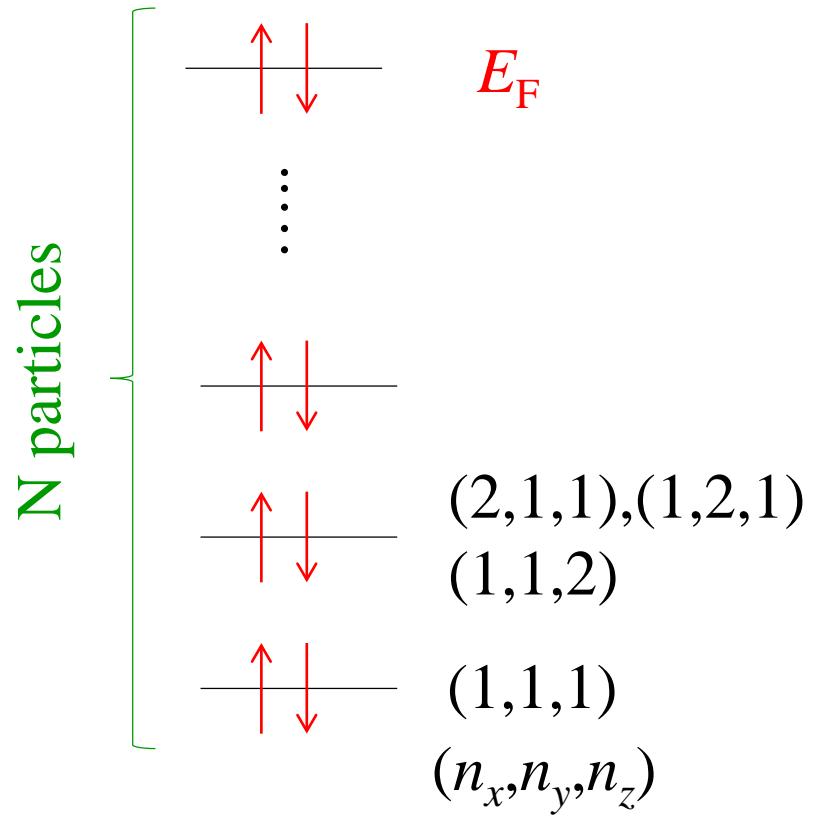
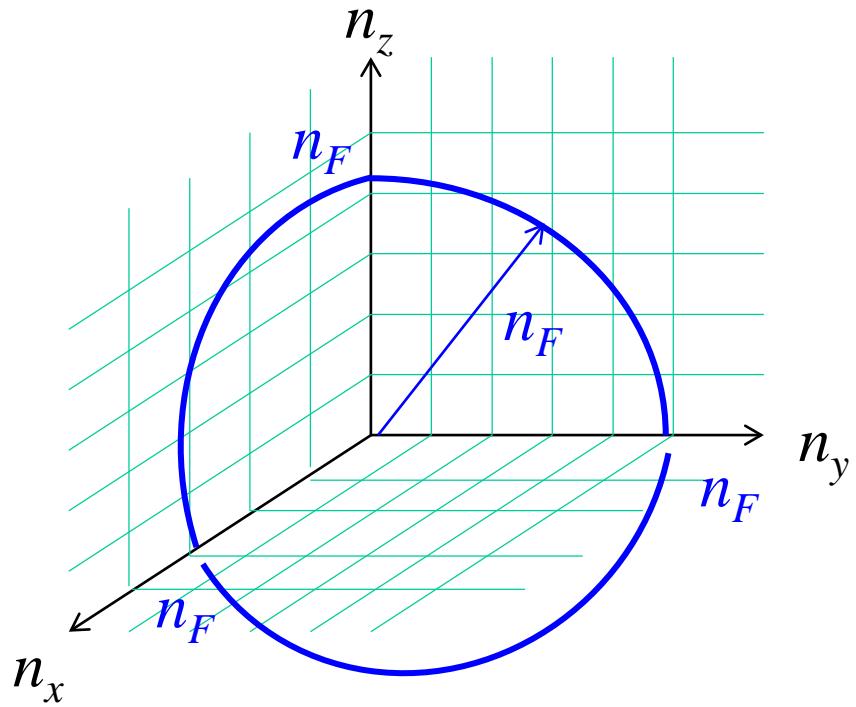
total energy

$$\begin{aligned} E_{\text{tot}} &= 2 \times \frac{1}{8} \int_{|\mathbf{n}| \leq n_F} E_{n_x n_y n_z} d^3 n \\ &= 2 \times \frac{1}{8} \int_{|\mathbf{n}| \leq n_F} \frac{\pi^2 \hbar^2}{2m L^2} (n_x^2 + n_y^2 + n_z^2) d^3 n \\ &= \dots = \frac{\hbar^2 \pi^3}{10m L^2} n_F^5 \end{aligned}$$

$$N = 2 \times \frac{1}{8} \times \frac{4}{3} \pi n_F^3$$

$$\longrightarrow \boxed{E_{\text{tot}} = \frac{3}{5} N \cdot \frac{\hbar^2}{2m} \left(3\pi^2 \frac{N}{V} \right)^{2/3} = \frac{3}{5} N E_F}$$

Fermi gas model



$$N = 2 \times \frac{1}{8} \times \frac{4}{3} \pi n_F^3$$

$$\begin{aligned} E_{\text{tot}} &= 2 \times \frac{1}{8} \int_{|\mathbf{n}| \leq n_F} E_{n_x n_y n_z} d^3 n \\ &= \dots = \frac{3}{5} N \cdot \frac{\hbar^2}{2m} \left(3\pi^2 \frac{N}{V} \right)^{2/3} \end{aligned}$$