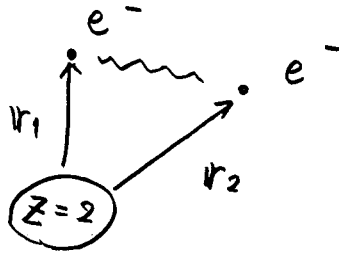


4. 原子, 分子の構造

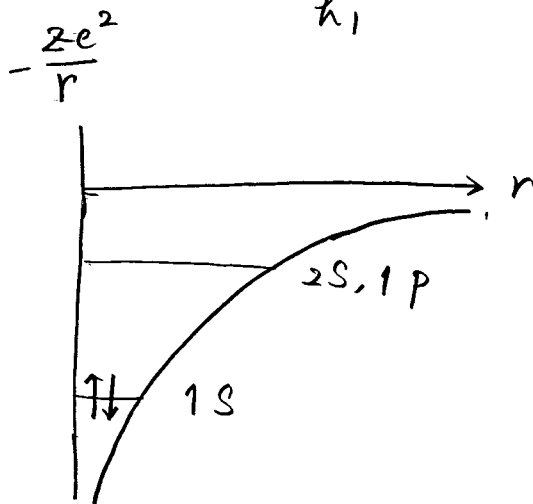
4.1. He 原子の構造



$$H = \frac{P_1^2}{2m} + \frac{P_2^2}{2m} - \frac{ze^2}{r_1} - \frac{ze^2}{r_2} + \frac{e^2}{|r_1 - r_2|}$$

。電子間相互作用を無視する近似

$$H_0 = \underbrace{\frac{P_1^2}{2m} - \frac{ze^2}{r_1}}_{h_1} + \underbrace{\frac{P_2^2}{2m} - \frac{ze^2}{r_2}}_{h_2}$$



$$\psi_{1s}(r) = 2 \left(\frac{Z}{a_0} \right)^{3/2} e^{-Zr/a_0} \times \frac{1}{\sqrt{4\pi}}$$

$$a_0 = \hbar / m c \alpha$$

$$\Psi_{gs}^{(0)}(r_1, r_2) = \psi_{1s}(r_1) \psi_{1s}(r_2) \cdot \underbrace{\frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)}_{|S=0\rangle}$$

↑

反対称化、スピン波動関数

$$E_{gs}^{(0)} = 2 E_{1s} = -m c^2 (2\alpha)^2 = -108.8 \text{ eV}$$

(note) ハミルトニアンを正確に解くと (数値計算)

$$E_{gs} = -78.975 \text{ eV}$$

・電子間相互作用の効果

摂動論で見積ると,

$$\begin{aligned}\Delta E &= \langle \Psi_{gs}^{(0)} | V_{ee} | \Psi_{gs}^{(0)} \rangle \\ &= \int d\mathbf{r}_1 d\mathbf{r}_2 \Psi_{gs}^{(0)*}(\mathbf{r}_1, \mathbf{r}_2) V_{ee}(\mathbf{r}_1, \mathbf{r}_2) \Psi_{gs}^{(0)}(\mathbf{r}_1, \mathbf{r}_2) \\ &= \int d\mathbf{r}_1 d\mathbf{r}_2 |\psi_{1s}(\mathbf{r}_1)|^2 |\psi_{1s}(\mathbf{r}_2)|^2 \cdot \frac{e^2}{|\mathbf{r}_1 - \mathbf{r}_2|} \\ &= \dots = \frac{5}{8} \frac{ze^2}{a_0}\end{aligned}$$

↓

$$\begin{aligned}E_{gs}^{(1)} &= E_{gs}^{(0)} + \Delta E = -108.8 + 34 \text{ eV} \\ &= -74.8 \text{ eV}\end{aligned}$$

○ よりよい近似

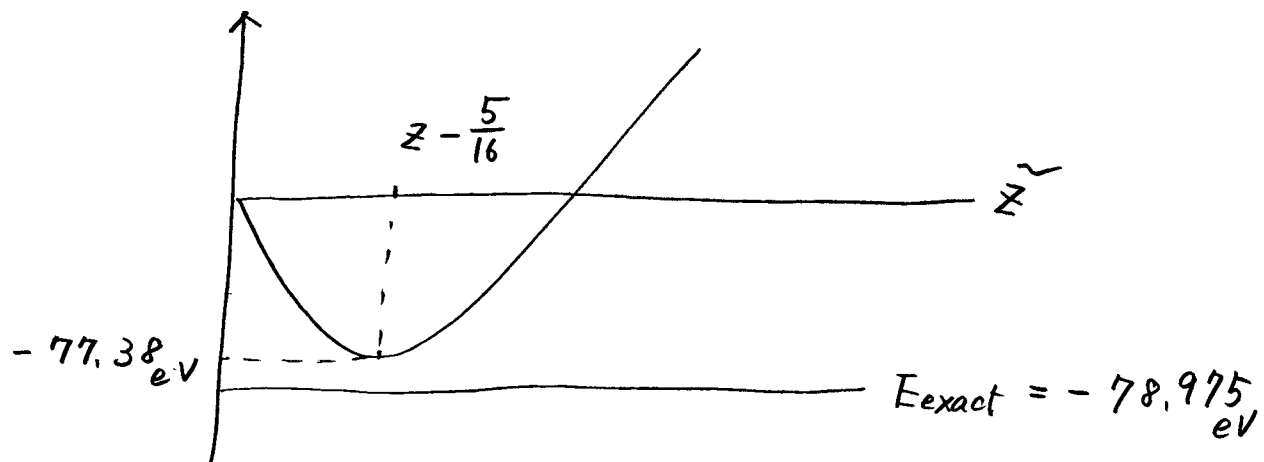
$$\Psi_{Egs}(r_1, r_2) \sim \tilde{\Psi}_{1s}(r_1) \tilde{\Psi}_{1s}(r_2) \quad |S=0\rangle$$

$$\tilde{\Psi}_{1s}(r) = 2 \left(\frac{\tilde{z}}{a_0} \right)^{3/2} e^{-\tilde{z}r/a_0} \cdot \frac{1}{\sqrt{4\pi}}$$

別の電子がいることにより原子核の電荷が
「無効」されるかもしれない
↓ $z \rightarrow \tilde{z} \quad \left[-\frac{ze^2}{r} \rightarrow -\frac{\tilde{z}e^2}{r} \right]$

\tilde{z} は変分法で決める

$$E(\tilde{z}) = \langle \Psi_{Egs} | H | \Psi_{Egs} \rangle$$



○ さらによい近似

r ごとに \tilde{z} の値を変える (ハートリー-フォック近似)

$$-\frac{ze^2}{r} \rightarrow -\frac{\tilde{z}(r)e^2}{r}$$

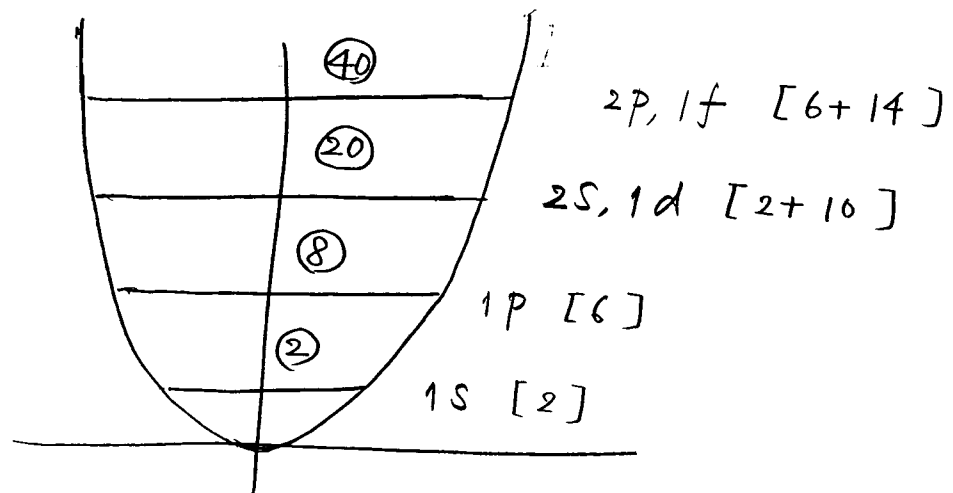
$$\downarrow E_{Egs}^{(HF)} = -77.87 \text{ eV}$$

(余談) 原子核の構造

魔法の数: 2, 8, 20, 28, 50, 82, 126

調和振動子ポテンシャルに核子をよめる

$$V(r) = \frac{1}{2} m \omega^2 r^2$$



$$E_{nl} = \left(2(n-1) + l + \frac{3}{2} \right) \hbar \omega$$

← $\hbar \omega$

縮退度: $2 \times (2l+1)$

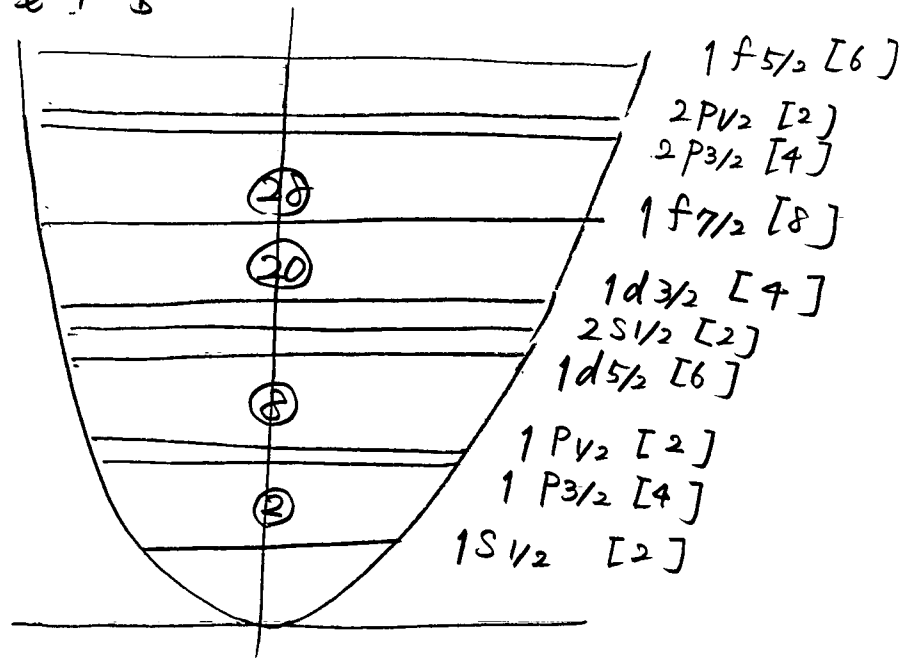
* $\hbar \omega$ - 軌道力が"入るな役割"

$$V(r) = \frac{1}{2} m \omega^2 r^2 + \alpha \vec{l} \cdot \vec{s}$$

(note) $\vec{j} = \vec{l} + \vec{s} \quad \rightarrow \quad \vec{l} \cdot \vec{s} = \frac{1}{2} (j^2 - l^2 - s^2)$

$$j = l \pm \frac{1}{2}$$

$$J = L + S$$

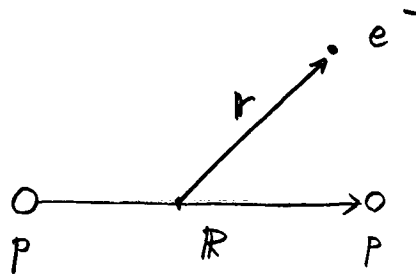


$$1f \rightarrow 1f_{7/2} \text{ \& } 1f_{5/2} \text{ 14 = 8 + 6) 分離}$$

4.2. 分子の構造

。断熱近似 (ボーン・オッペンハイマー近似)

最も単純な分子: H_2^+



$$H = -\frac{\hbar^2}{2\mu} \nabla_R^2 - \frac{\hbar^2}{2m} \nabla_r^2 + \frac{e^2}{R} - \frac{e^2}{|r + \frac{R}{2}|} - \frac{e^2}{|r - \frac{R}{2}|}$$

$$\mu = \frac{m_P}{2} \sim \frac{940}{2} \text{ (MeV/c}^2\text{)} = 470 \text{ (MeV/c}^2\text{)}$$

$$m \sim m_e = 0.5 \text{ (MeV/c}^2\text{)}$$

$\mu \gg m_e \rightarrow$ 陽子は相対距離 R で固定 (止まっている)

$$\downarrow$$

$$\left(-\frac{\hbar^2}{2m} \nabla_r^2 - \frac{e^2}{|r + \frac{R}{2}|} - \frac{e^2}{|r - \frac{R}{2}|} \right) \Psi_R(r) = E(R) \Psi_R(r)$$

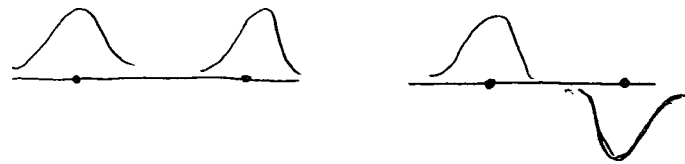
III.
he

○ 線形結合近似

$$|\varphi_R^{(\pm)}\rangle \sim \left| \begin{array}{c} \text{IS} \\ \odot \text{---} \bullet \end{array} \right\rangle \pm \left| \begin{array}{c} \text{IS} \\ \bullet \text{---} \odot \end{array} \right\rangle$$

$$= C_{\pm}(R) \left(\underbrace{\psi_{1s}(r - \frac{R}{2})}_{\psi_{1s}(L)} \pm \underbrace{\psi_{1s}(r + \frac{R}{2})}_{\psi_{1s}(R)} \right)$$

規格化:



$$1 = \langle \varphi_R^{(\pm)} | \varphi_R^{(\pm)} \rangle$$

$$= |C_{\pm}(R)|^2 \langle \psi_{1s}(L) \pm \psi_{1s}(R) | \psi_{1s}(L) \pm \psi_{1s}(R) \rangle$$

$$= |C_{\pm}(R)|^2 \left\{ \langle \psi_{1s}(L) | \psi_{1s}(L) \rangle + \langle \psi_{1s}(R) | \psi_{1s}(R) \rangle \right. \\ \left. \pm \langle \psi_{1s}(L) | \psi_{1s}(R) \rangle + \langle \psi_{1s}(R) | \psi_{1s}(L) \rangle \right\}$$

$$= |C_{\pm}(R)|^2 \left\{ 2 \pm 2 \int dr \psi_{1s}(r - \frac{R}{2})^* \psi_{1s}(r + \frac{R}{2}) \right\}$$

$$\equiv S(R)$$

$$= \dots = \left(1 + \frac{R}{a_0} + \frac{R^2}{3a_0^2}\right) e^{-R/a_0}$$

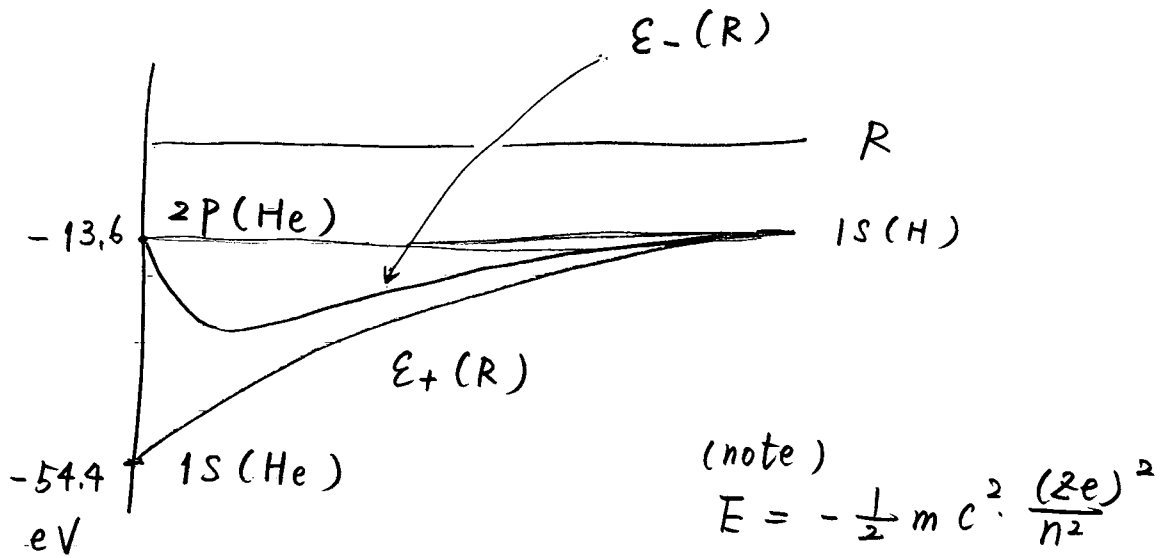
$$\downarrow$$

$$|C_{\pm}(R)|^2 = \frac{1}{2(1 \pm S(R))}$$

$$\hookrightarrow \begin{cases} \rightarrow 0 & (R \rightarrow 0) \\ \rightarrow 1 & (R \rightarrow \infty) \end{cases}$$

エネルギー

$$E_{\pm}(R) = |C_{\pm}(R)|^2 \langle \psi_{1s}(L) \pm \psi_{1s}(R) | \hat{H} | \psi_{1s}(L) \pm \psi_{1s}(R) \rangle$$



(note) $|\psi_R^{(+)}\rangle \rightarrow |\psi_{1s}(\text{He})\rangle$
 $|\psi_R^{(-)}\rangle \rightarrow |\psi_{2p}(\text{He})\rangle$
 $R \rightarrow 0$

2つの陽子の入れかえに対し反対称
(パリティ: マイナス)

• 基底状態

