

# Application of beyond-mean-field approach to hypernuclei and heavy-ion fusion reactions



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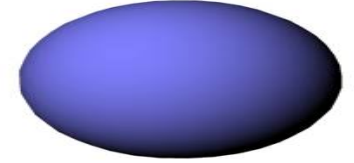
堯江明 (西南大学(重慶)・東北大学(仙台)

→ 北 Carolina U.)



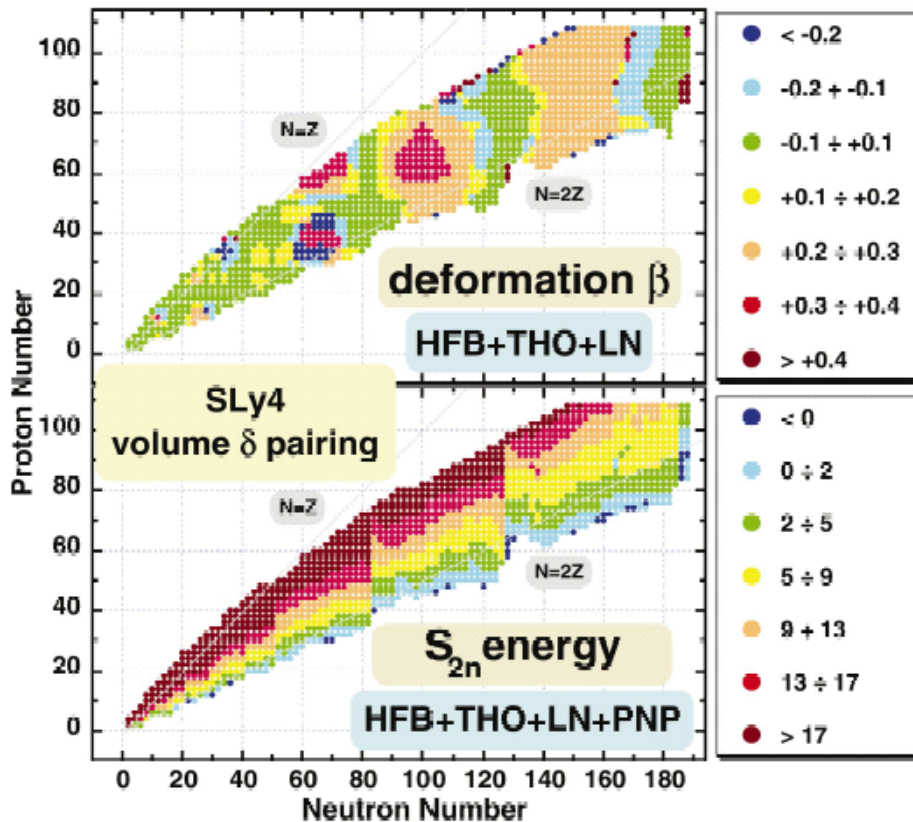
- 1. Introduction: mean-field approximation and beyond*
- 2. Application to heavy-ion fusion reactions*
- 3. Microscopic particle-rotor model for hypernuclei*
- 4. Summary*

# Mean-field approximation and beyond



## Self-consistent mean-field (Hartree-Fock) method:

- independent particles in a mean-field potential
- global theory for **the whole nuclear chart**
- body-fixed frame → intuitive picture for nuclear deformation
- optimized shape can be automatically determined



# Mean-field approximation and beyond

- ✓ body-fixed frame formalism → intuitive picture of nuclear def.

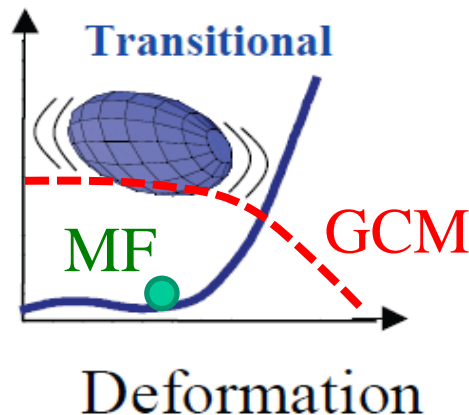
## Drawbacks of the mean-field approximation : nuclear spectrum

- ✓ spectrum: lab-frame ← transformation from intrinsic to lab. frames
- nuclear spectrum: requires to go beyond the mean-field approximation

$$|\Psi_{IM}(\beta)\rangle = \hat{P}_{MK}^I \hat{P}^N \hat{P}^Z |\Psi_{MF}(\beta)\rangle$$

angular momentum + particle number projections

- ✓ quantum fluctuation



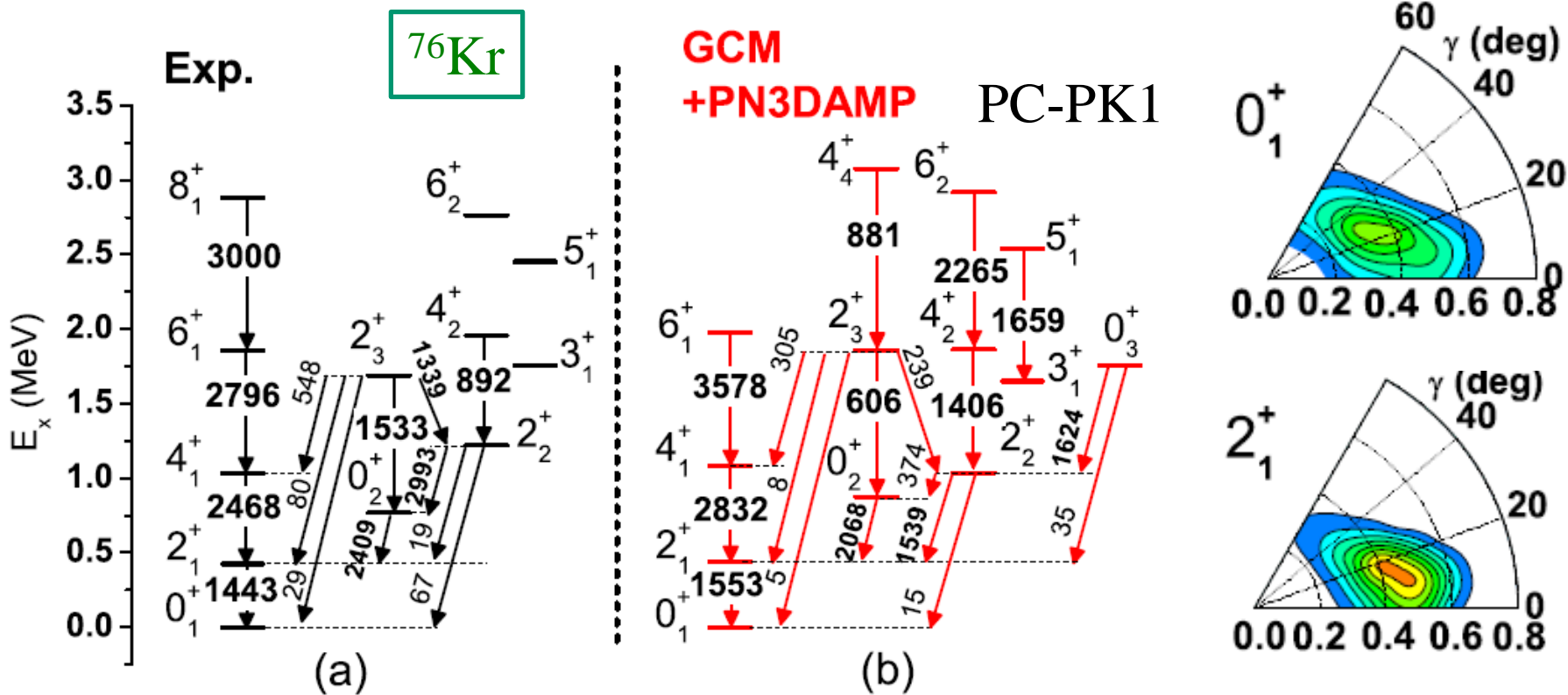
$$|\Phi_{IM}\rangle = \int d\beta f(\beta) |\Psi_{IM}(\beta)\rangle$$

generator coordinate method (GCM)

single → multi Slater determinants

# beyond mean-field approximation

- ✓ angular momentum + particle number projections
- ✓ quantum fluctuation (GCM)



J.M. Yao, K.H., Z.P. Li, J. Meng, and P. Ring, PRC89 ('14) 054306

cf. collective coordinates: Quadrupole moments (local operators)

## beyond mean-field approximation

- ✓ angular momentum + particle number projections
- ✓ quantum fluctuation (GCM)

In this talk:

applications to

1. Heavy-ion fusion reactions at sub-barrier energies
2. Low-lying spectrum of  $\Lambda$ -hypernuclei

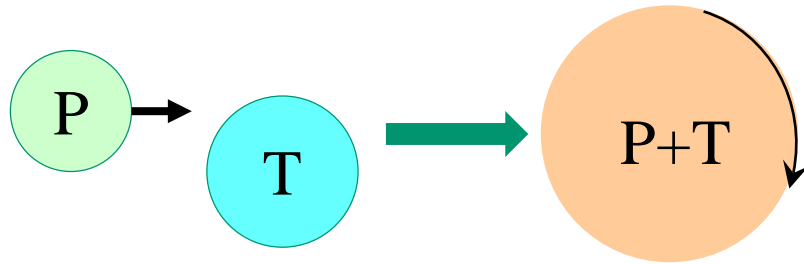
# Heavy-ion fusion reactions at sub-barrier energies

K.H. and J.M. Yao, PRC91 ('15) 064606

J.M. Yao and K.H., PRC94 ('16) 011303(R)

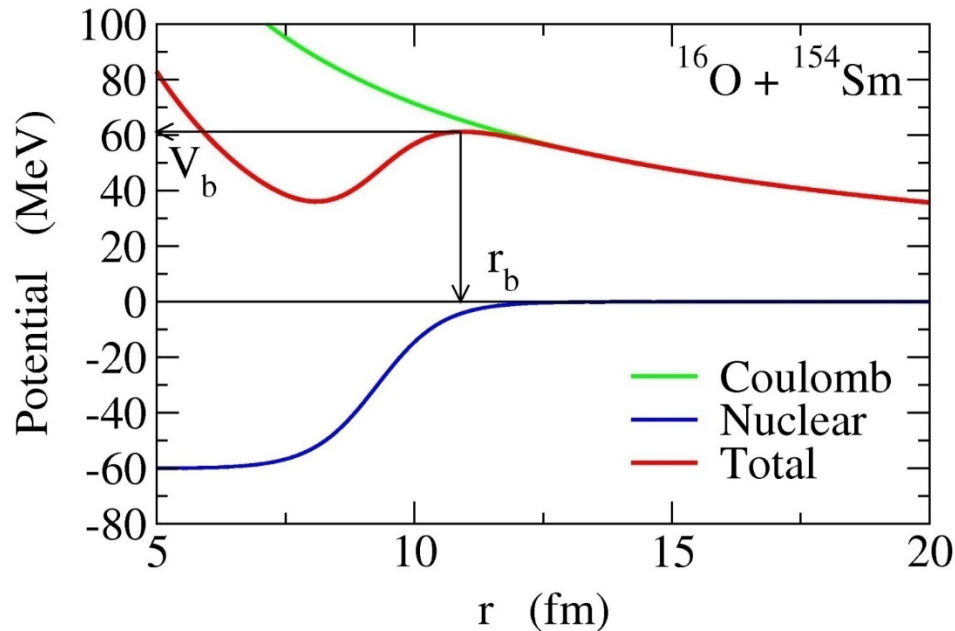
# Heavy-ion fusion reactions

Fusion: compound nucleus formation



$$A_{\text{CN}} = A_{\text{P}} + A_{\text{T}}$$

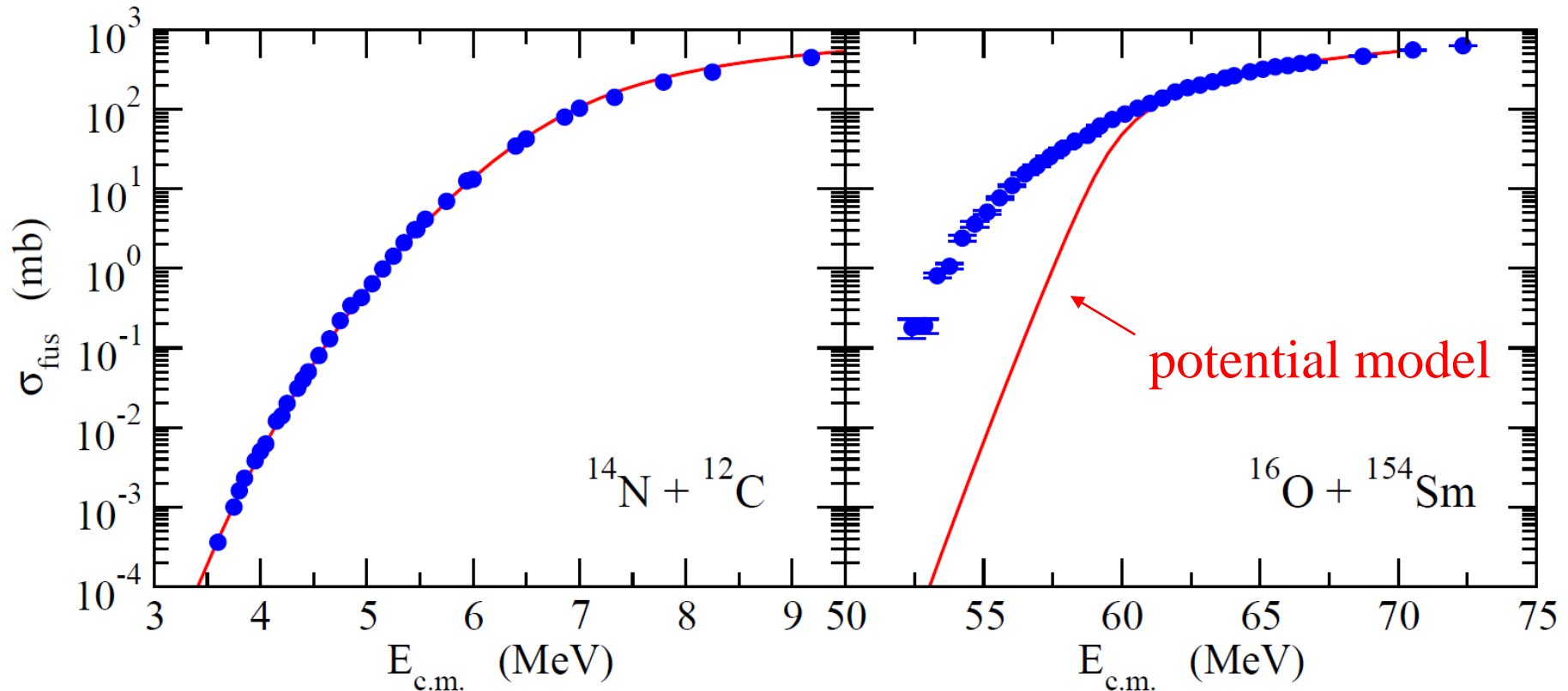
potential model:  $V(r) + \text{absorption}$



$$\sigma_{\text{fus}}(E) = \frac{\pi}{k^2} \sum_l (2l + 1) (1 - |S_l(E)|^2)$$

# Discovery of large sub-barrier enhancement of $\sigma_{\text{fus}}$

potential model:  $V(r) + \text{absorption}$



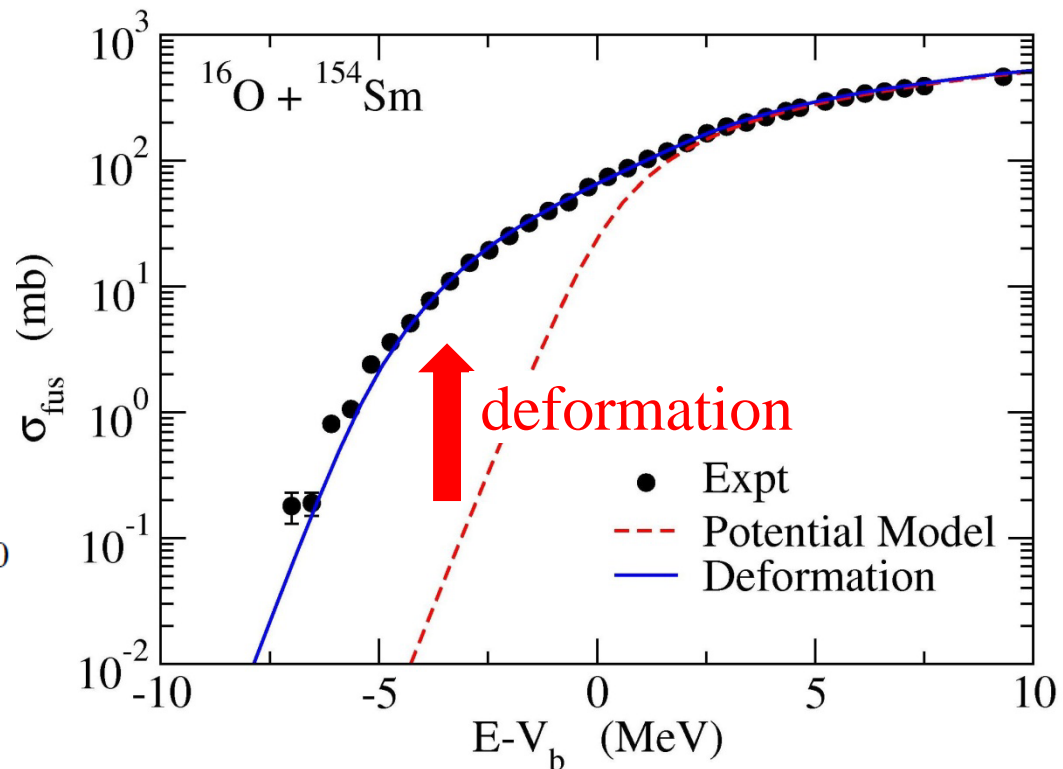
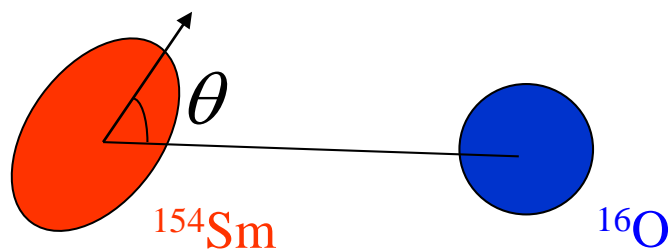
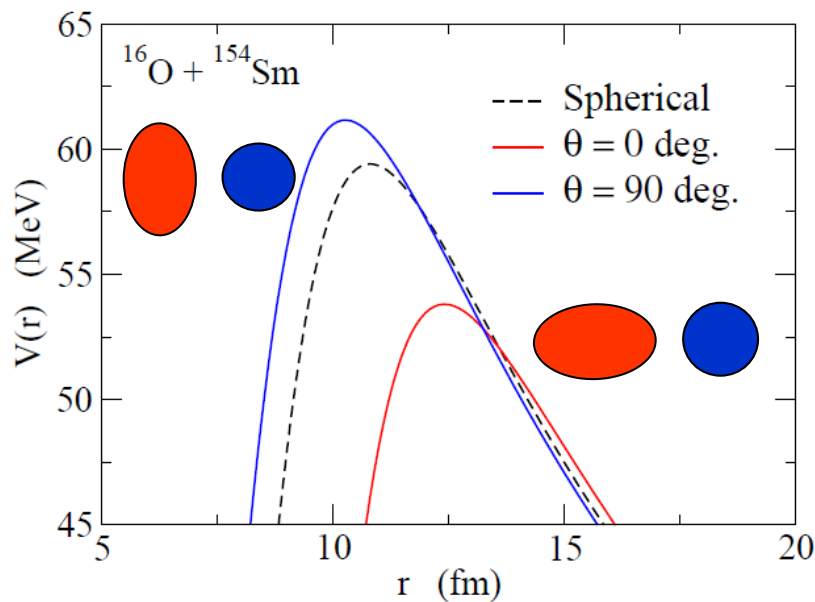
cf. seminal work:

R.G. Stokstad et al., PRL41('78) 465



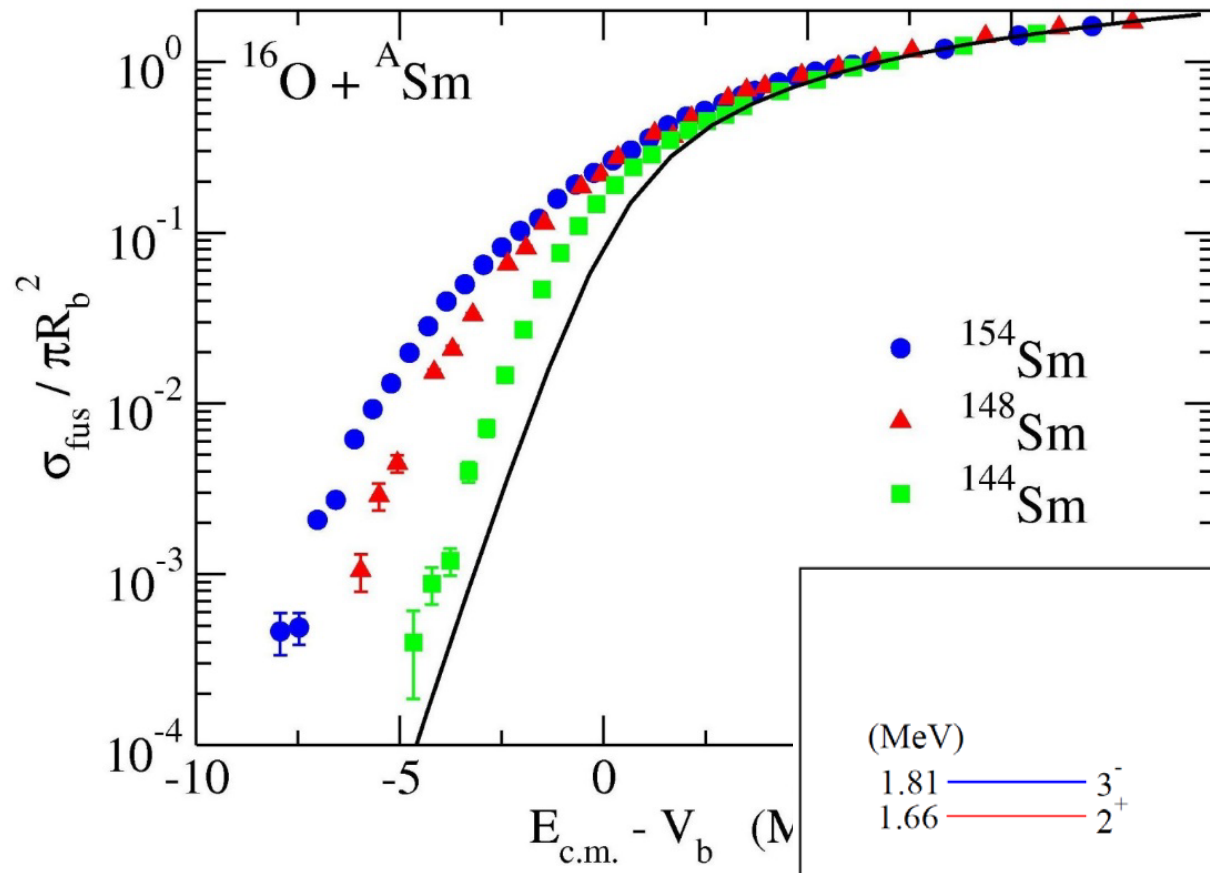
# Effect of nuclear deformation

$^{154}\text{Sm}$  : a deformed nucleus with  $\beta_2 \sim 0.3$

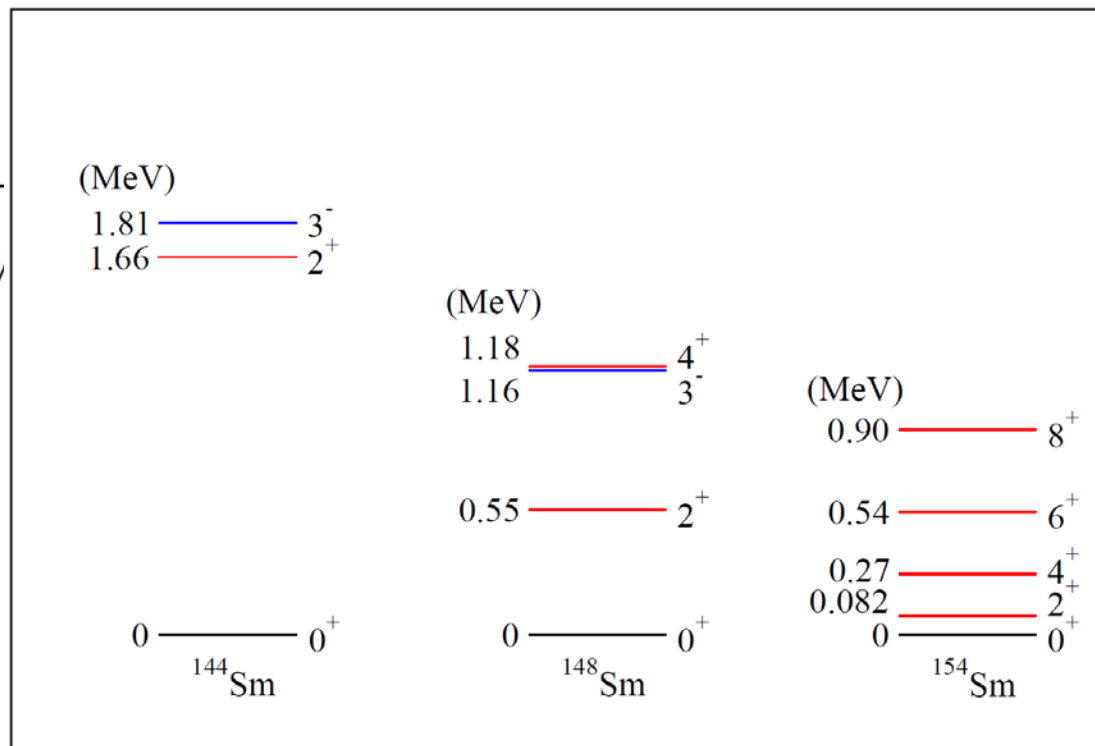


$$\sigma_{\text{fus}}(E) = \int_0^1 d(\cos \theta) \sigma_{\text{fus}}(E; \theta)$$

**Fusion: strong interplay between nuclear structure and nuclear reaction**

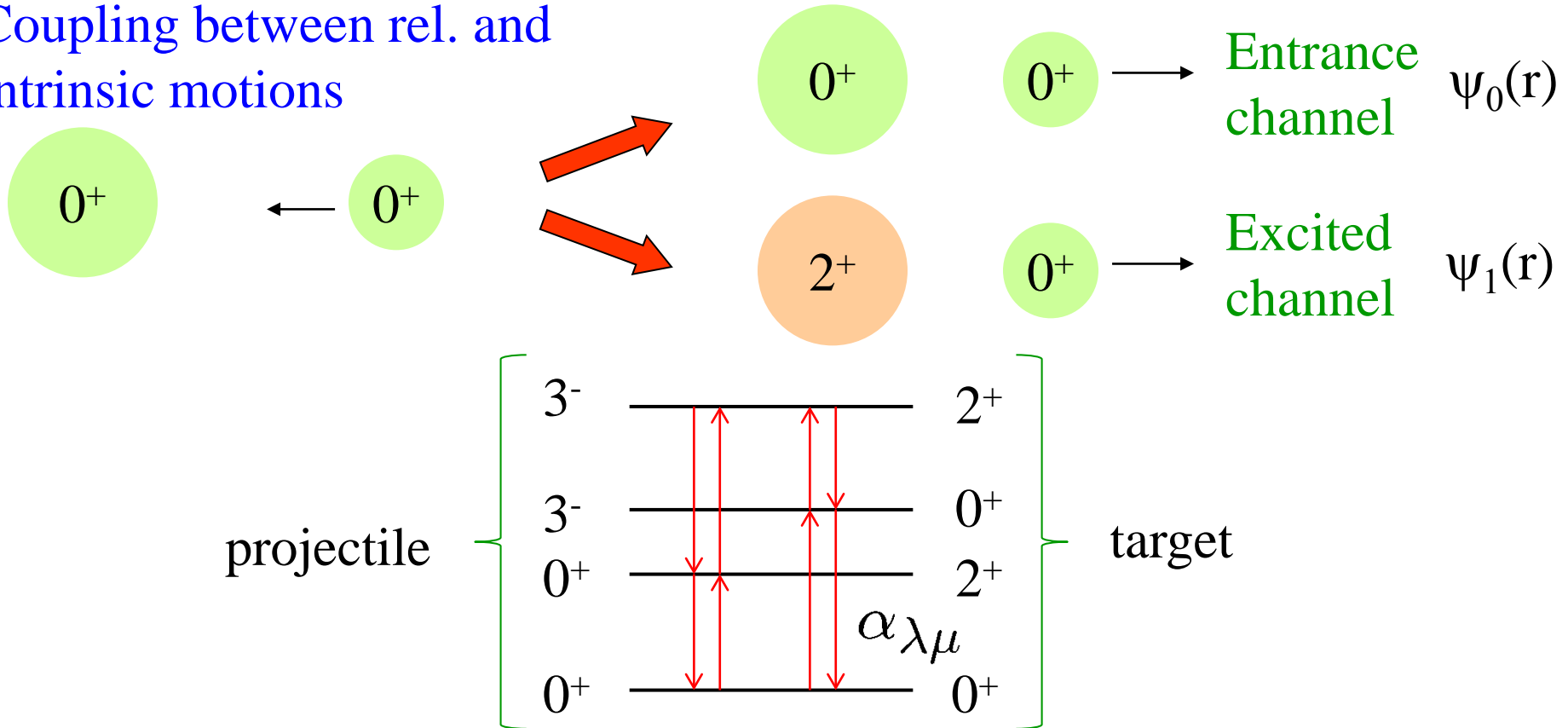


Strong target dependence  
at  $E < V_b$



# Coupled-Channels method

Coupling between rel. and intrinsic motions



$$\Psi(\mathbf{r}, \xi) = \sum_k \psi_k(\mathbf{r}) \phi_k(\xi)$$



coupled Schroedinger equations for  $\psi_k(\mathbf{r})$

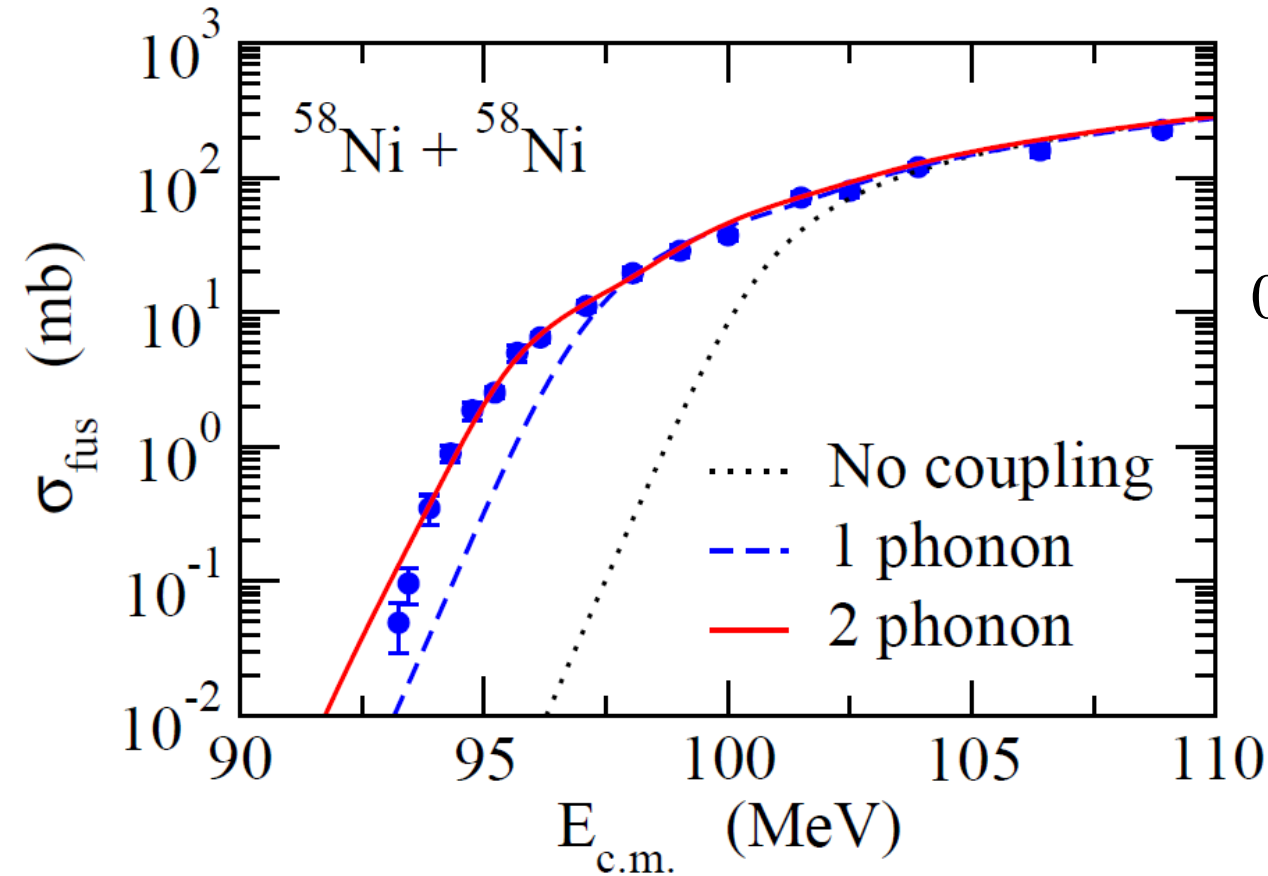
$$V_{\text{coup}}(r, \xi) \sim -R_T \frac{dV_N(r)}{dr} \alpha_\lambda Y_\lambda(\hat{\mathbf{r}})$$

excitation operator

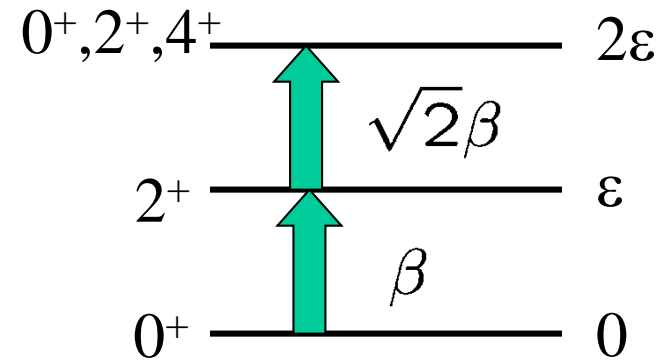
# Semi-microscopic modeling of sub-barrier fusion

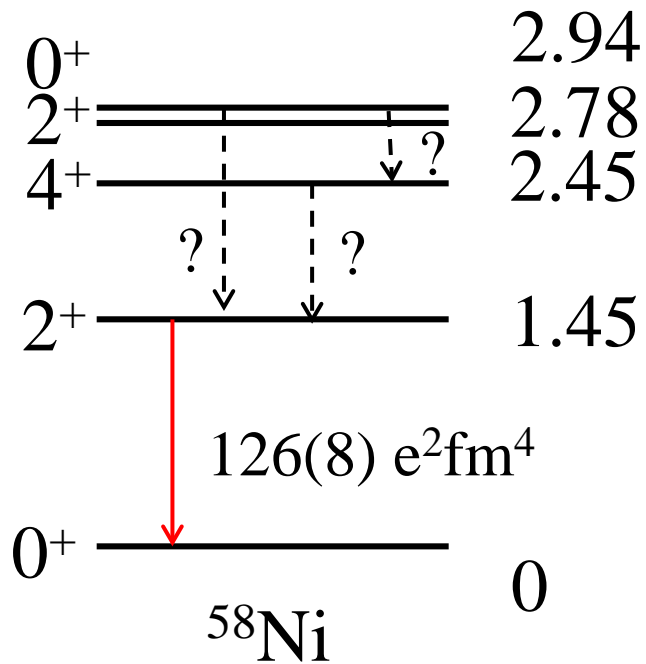
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multi-phonon excitations



simple harmonic oscillator





$$Q(2_1^+) = -10 \pm 6 e\text{fm}^2$$

Simple harmonic oscillator  
: justifiable?

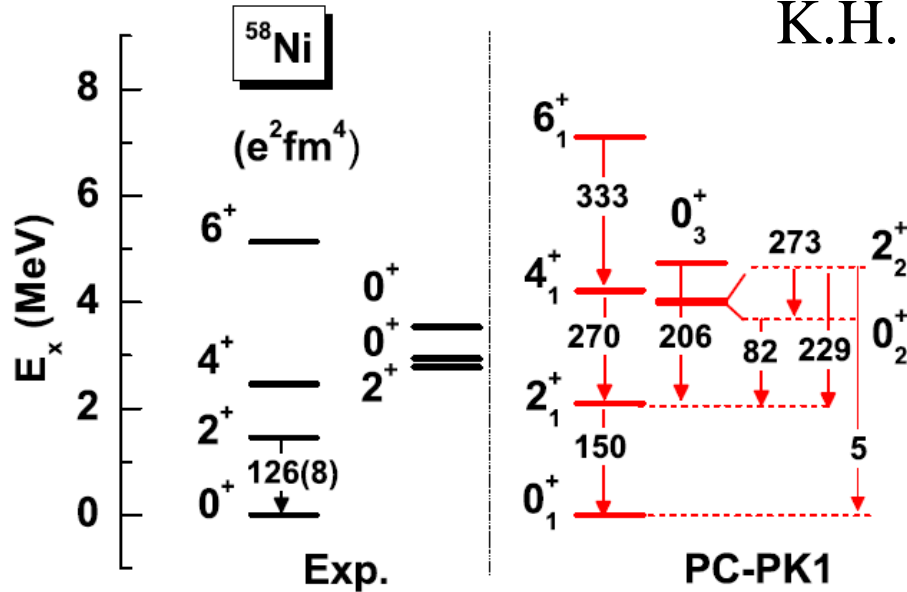


beyond mean-field approach  
to account for anharmonicity

cf. GCM: large amplitude  
collective motions

# Semi-microscopic coupled-channels model for sub-barrier fusion

K.H. and J.M. Yao, PRC91 ('15) 064606



microscopic  
multi-pole operator

$$\checkmark V_{\text{coup}} \sim -R_T \frac{dV_N}{dr} \alpha_\lambda \cdot Y_\lambda(\hat{r}) \rightarrow -R_T \frac{dV_N}{dr} Q_\lambda \cdot Y_\lambda(\hat{r})$$

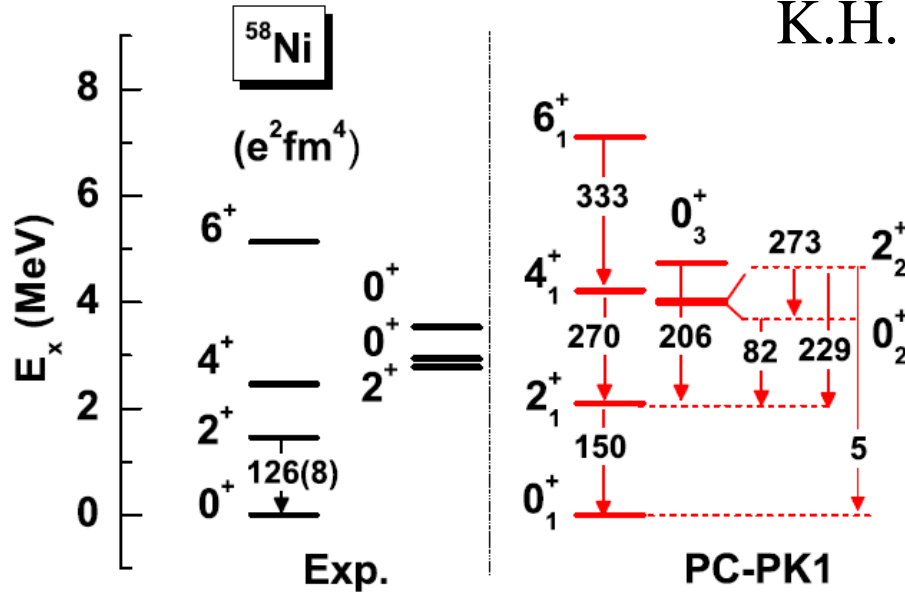
- ✓  $M(E2)$  from MR-DFT calculation ← among higher members of phonon states
- ✓ scale to the empirical  $B(E2; 2_1^+ \rightarrow 0_1^+)$  of phonon states

$$B(E2; I_i \rightarrow I_f) \rightarrow f \times B(E2; I_i \rightarrow I_f)$$

$$f = \frac{B(E2; 2^+ \rightarrow 0^+)_{\text{exp}}}{B(E2; 2^+ \rightarrow 0^+)_{\text{calc.}}}$$

# Semi-microscopic coupled-channels model for sub-barrier fusion

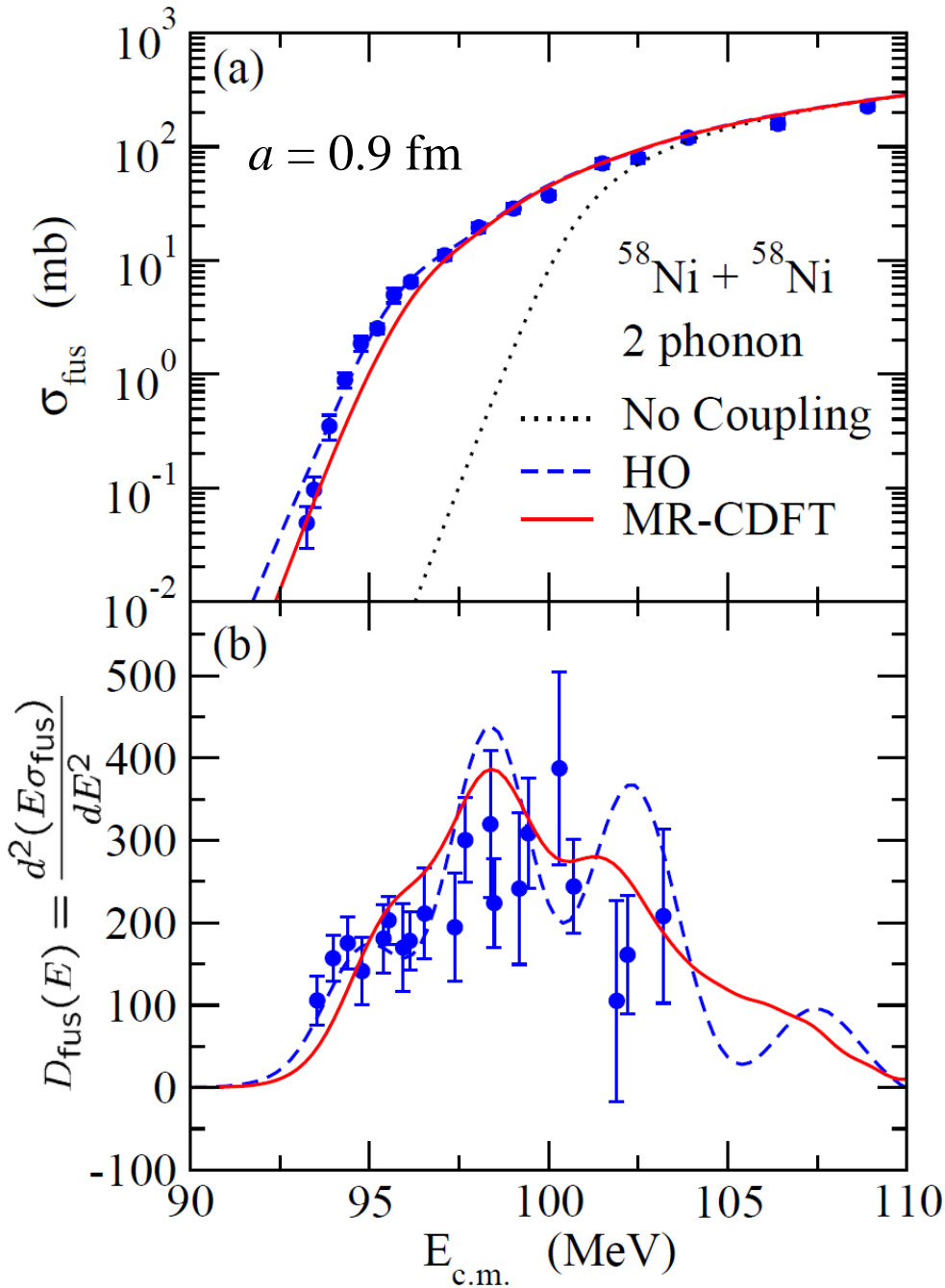
K.H. and J.M. Yao, PRC91 ('15) 064606



microscopic  
multi-pole operator

$$\checkmark \quad V_{\text{coup}} \sim -R_T \frac{dV_N}{dr} \alpha_\lambda \cdot Y_\lambda(\hat{r}) \rightarrow -R_T \frac{dV_N}{dr} Q_\lambda \cdot Y_\lambda(\hat{r})$$

- ✓  $M(E2)$  from MR-DFT calculation ← among higher members of phonon states
- ✓ scale to the empirical  $B(E2; 2_1^+ \rightarrow 0_1^+)$
- ✓ still use a phenomenological potential
- ✓ use the experimental values for  $E_x$
- ✓  $\beta_N$  and  $\beta_C$  from  $M_n/M_p$  for each transition
- ✓ axial symmetry (no  $3^+$  state)



$^{58}\text{Ni} + ^{58}\text{Ni}$

anharmonicity of  $2^+$  phonon  
 $\rightarrow$  only a minor improvement

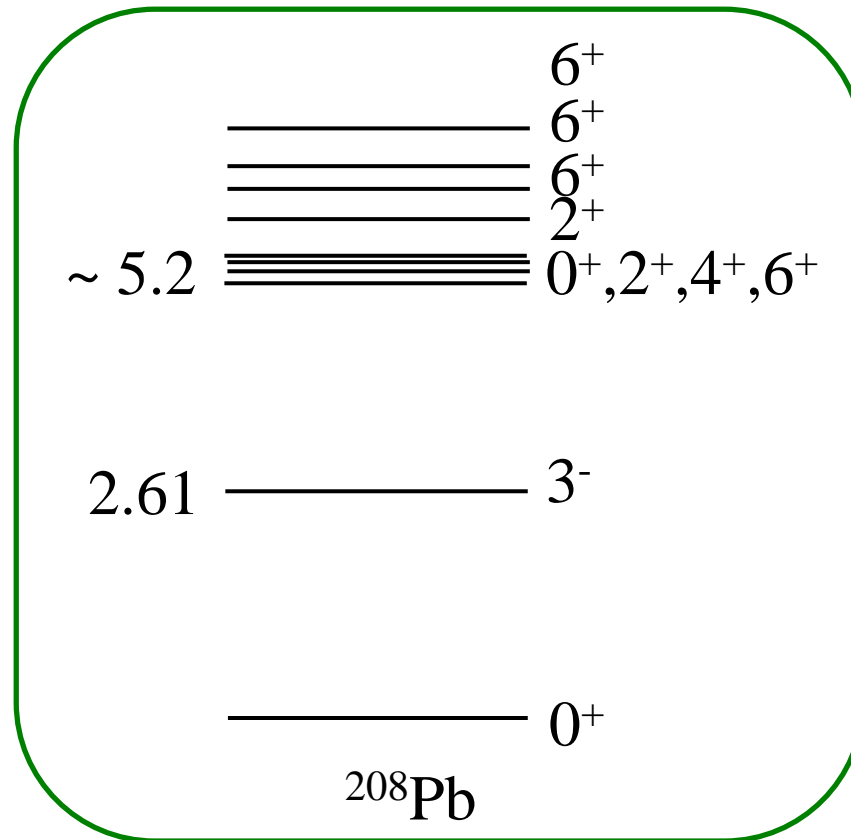


Next, more non-trivial case  
 with  $2^+ - 3^-$  coupling:  
 anharmonicity of oct. vib.  
 in  $^{208}\text{Pb}$



# Application to $^{16}\text{O} + ^{208}\text{Pb}$ fusion reaction

double-octupole phonon states in  $^{208}\text{Pb}$



M. Yeh, M. Kadi, P.E. Garrett et al., PRC57 ('98) R2085

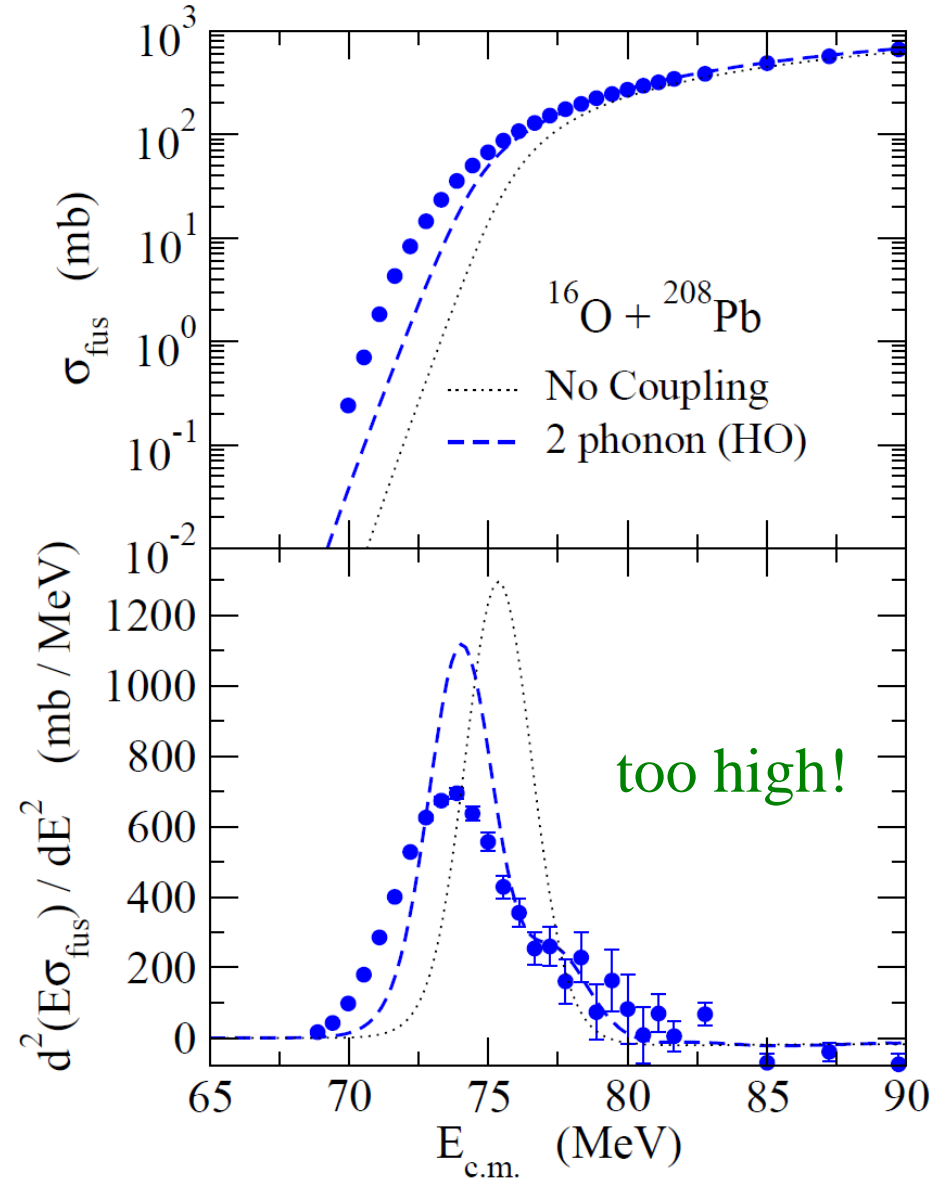
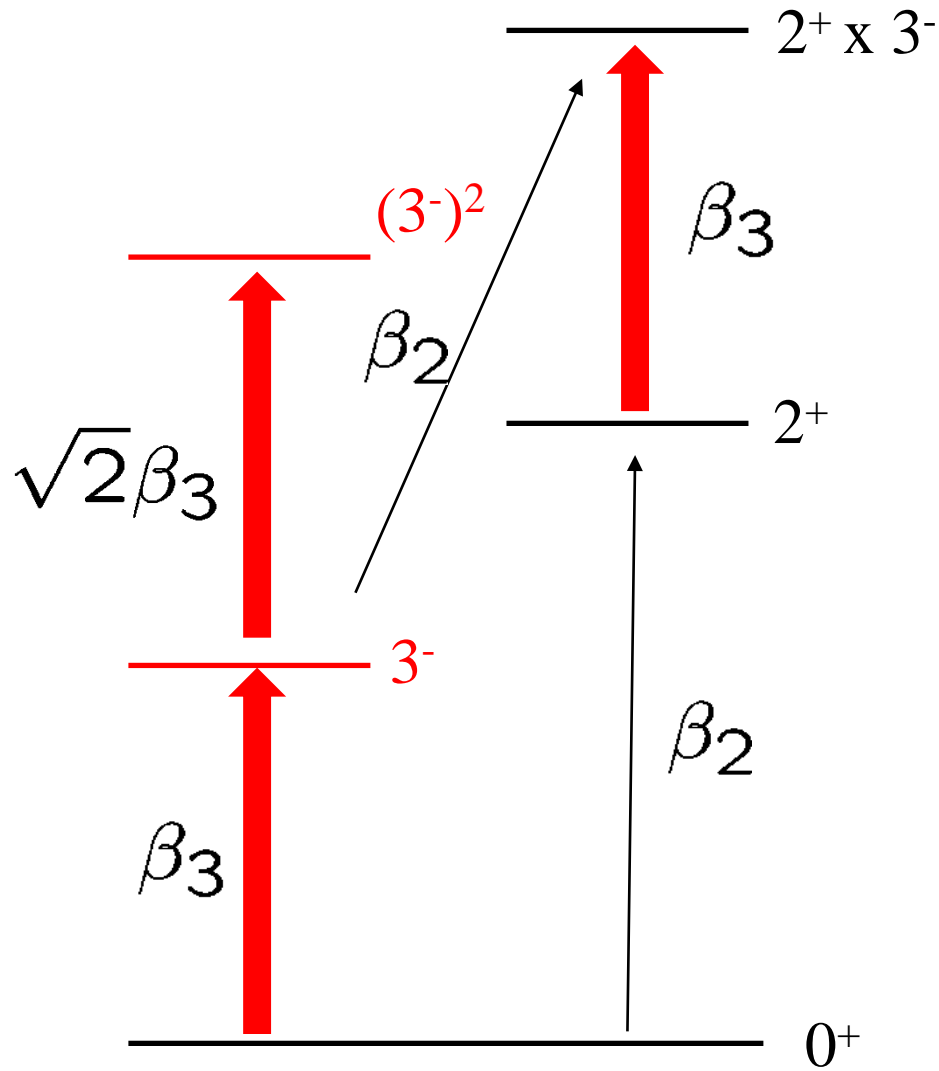
K. Vetter, A.O. Macchiavelli et al., PRC58 ('98) R2631

V. Yu. Pnomarev and P. von Neumann-Cosel, PRL82 ('99) 501

B.A. Brown, PRL85 ('00) 5300

large fragmentations, especially  $6^+$  state

# Application to $^{16}\text{O} + ^{208}\text{Pb}$ fusion reaction



cf. C.R. Morton et al., PRC60('99) 044608

expt. data

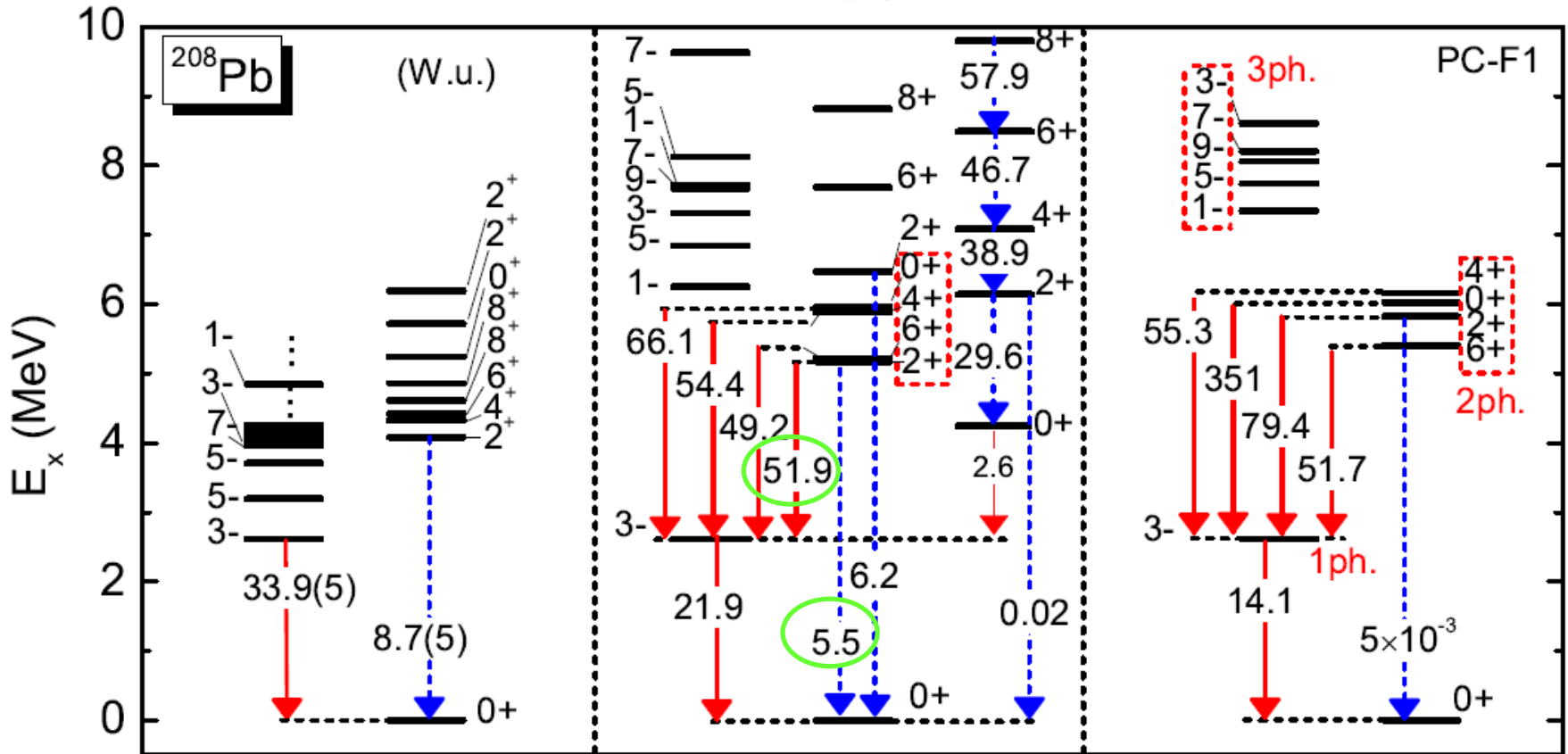
fluctuation both  
in  $\beta_3$  and  $\beta_2$

fluctuation in  $\beta_3$   
frozen at  $\beta_2=0$

(a) Exp.

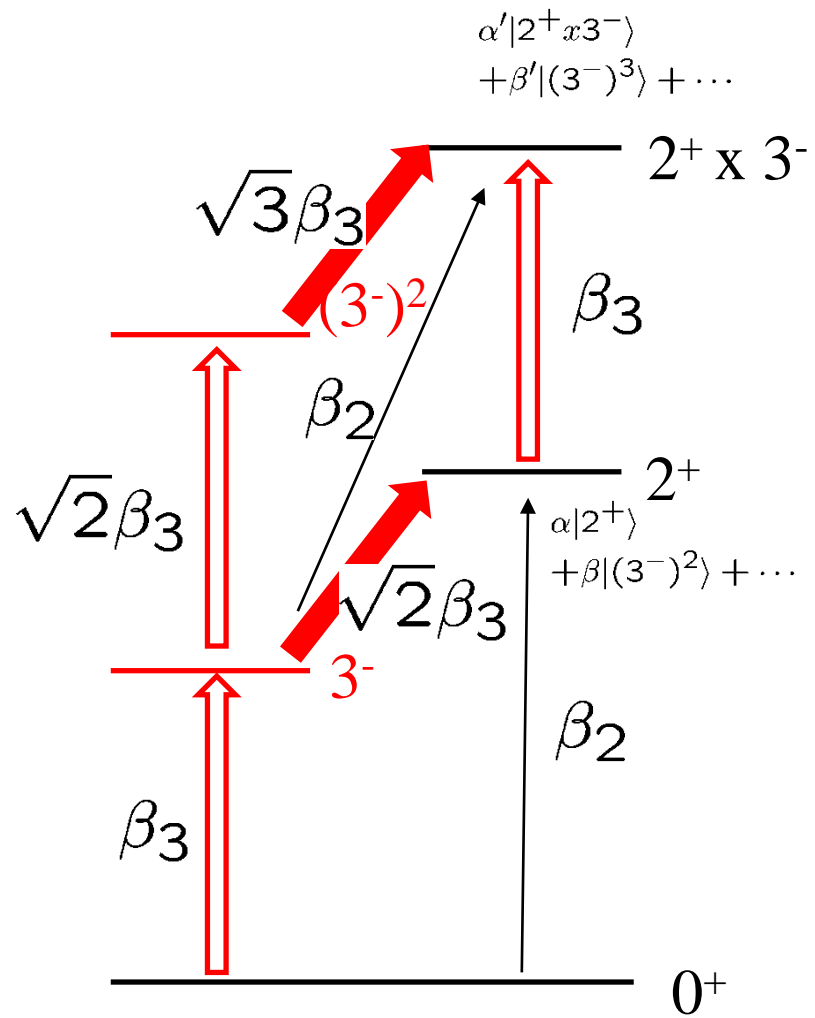
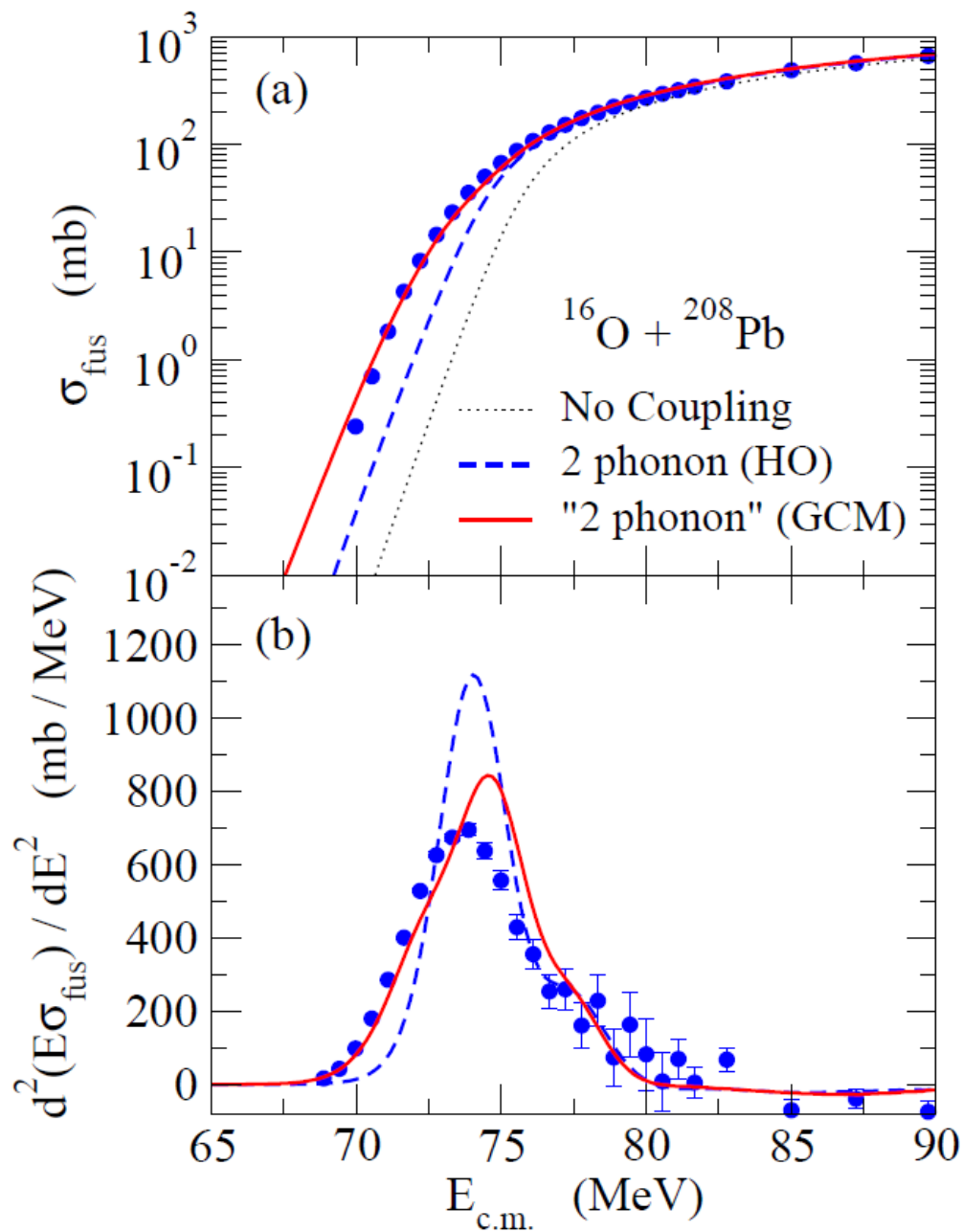
(b) GCM ( $\beta_2$ - $\beta_3$ )

(c) GCM ( $\beta_3$ )



$2_1^+$  state: strong coupling both to g.s. and  $3_1^-$

$$\longrightarrow |2_1^+\rangle = \alpha|2^+\rangle_{\text{HO}} + \beta|[3^- \otimes 3^-]^{(I=2)}\rangle_{\text{HO}} + \dots$$



J.M. Yao and K.H.,  
PRC94 ('16) 11303(R)

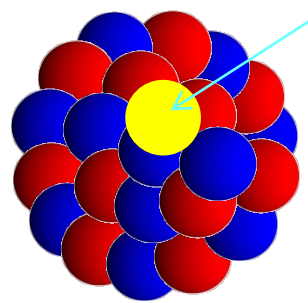
# Low-lying spectra of $\Lambda$ -hypernuclei

H. Mei, K.H., J.M. Yao, T. Motoba, PRC90 ('14) 064302  
PRC91 ('15) 064305  
PRC93 ('16) 044307  
two more papers  
in preparation  
H. Mei, K.H., J.M. Yao, PRC93 ('16) 011301(R)

# Introduction

## $\Lambda$ hypernucleus

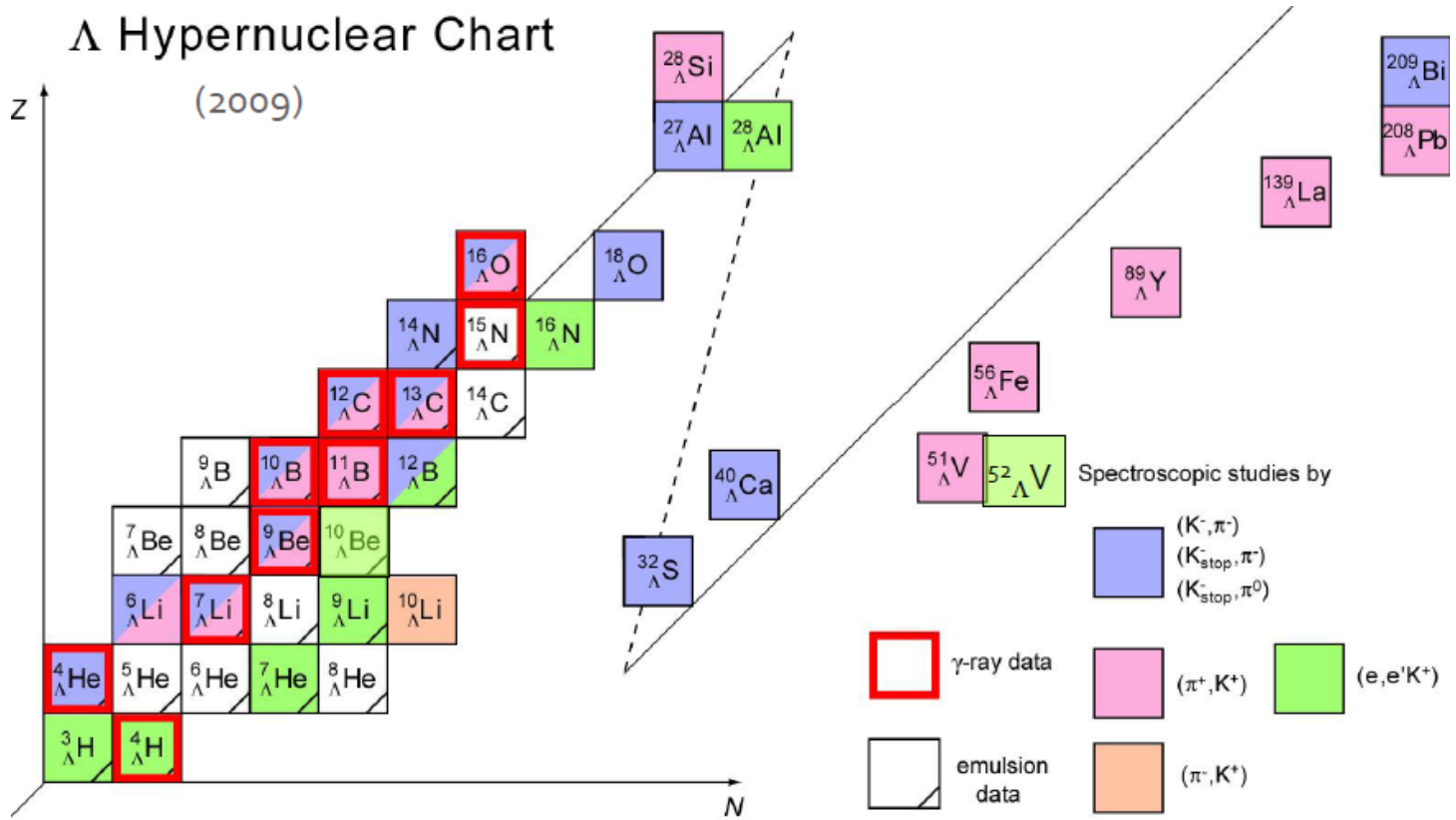
$\Lambda$  particle: the lightest hyperon  
(no charge, no isospin)



proton

neutron

\*no Pauli principle between nucleons and a  $\Lambda$  particle



# Introduction

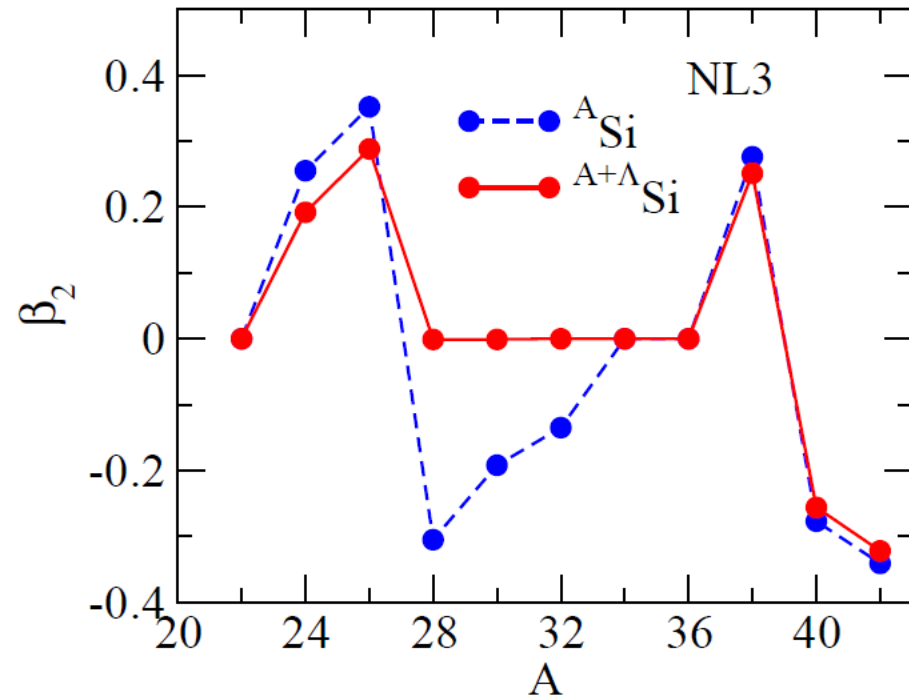
**Impurity effects:** one of the main interests of hypernuclear physics

**how does  $\Lambda$  affect several properties of atomic nuclei?**

- size, shape, density distribution, single-particle energy, shell structure, fission barrier.....

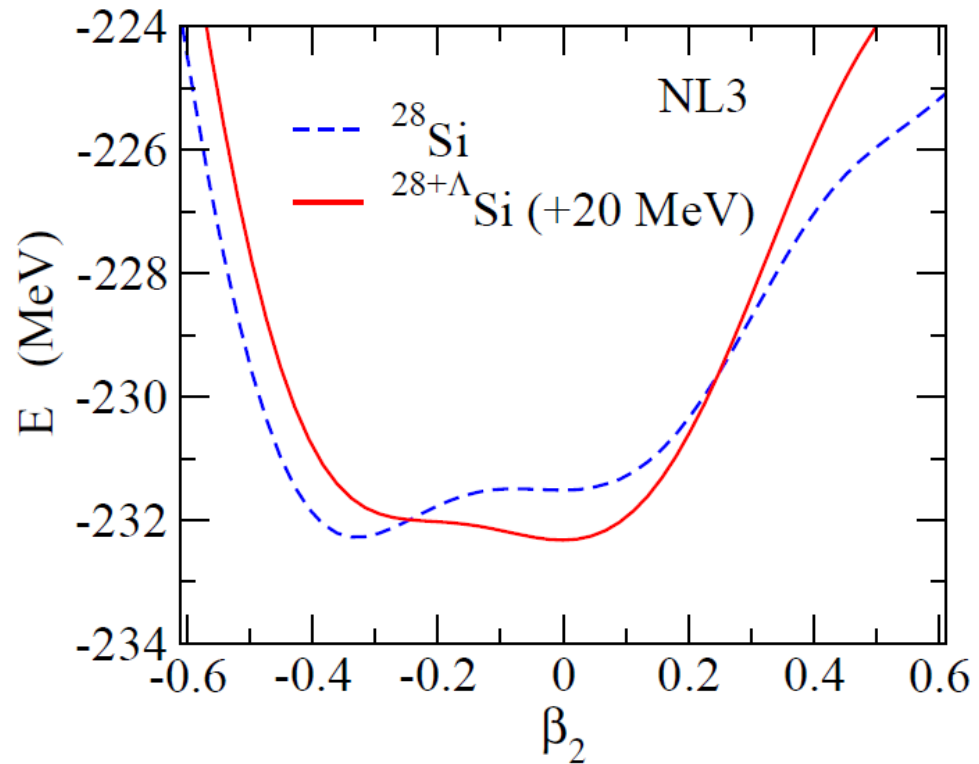
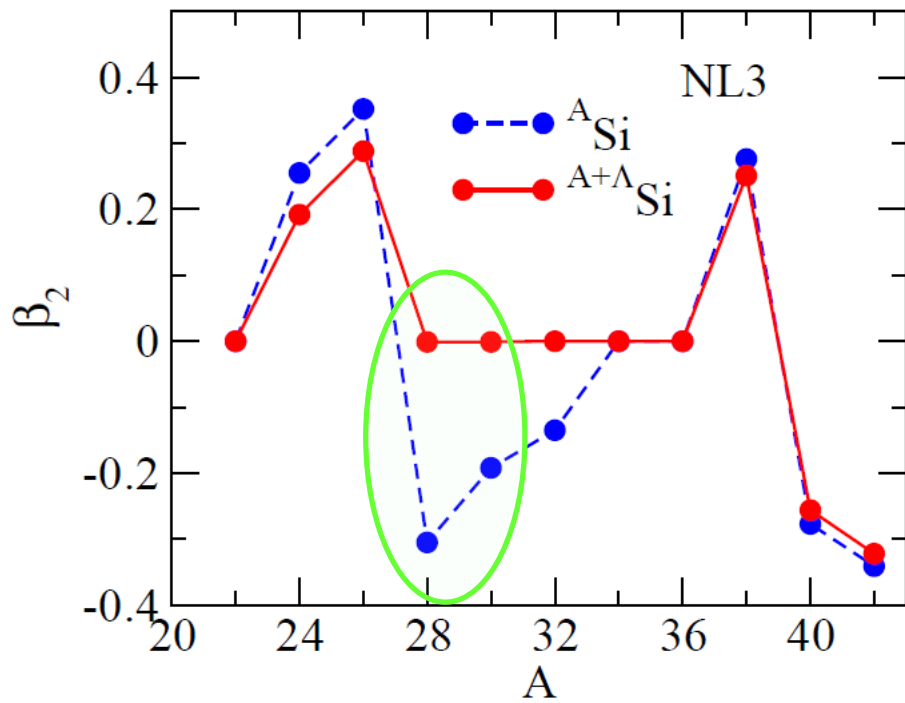
Theoretical approaches:

- ✓ cluster model
- ✓ shell model
- ✓ AMD
- ➔ ✓ self-consistent mean-field models



Myaing Thi Win and K.H.,  
PRC78('08)054311

# Si isotopes



Myaing Thi Win and K.Hagino, PRC78('08)054311

soft energy surface  $\rightarrow$  role of beyond-mean-field effect?





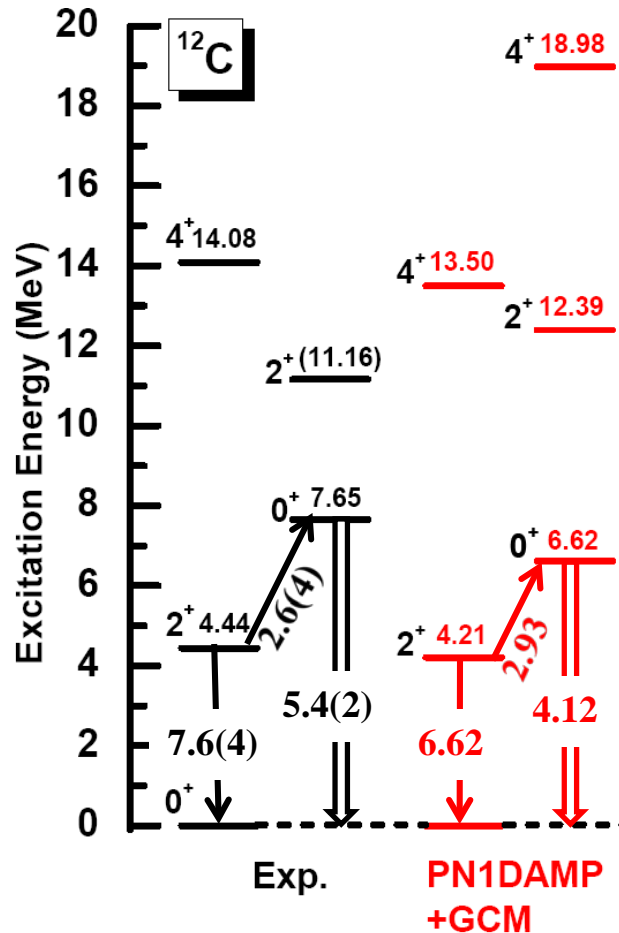


# Microscopic Particle-Rotor Model for $\Lambda$ hypernuclei

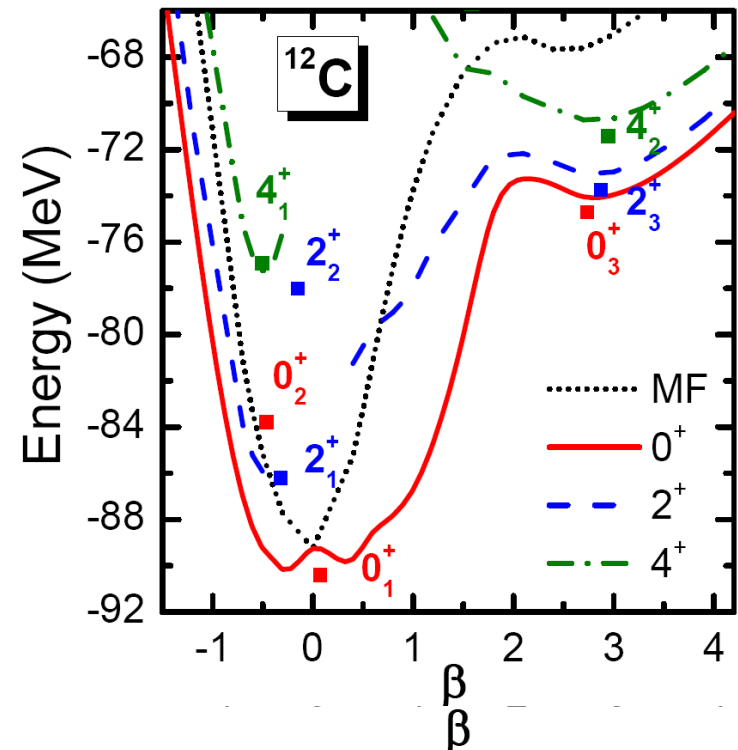
Example:  $^{13}_{\Lambda}\text{C}$

i) beyond mean-field calculations for e-e core ( $^{12}\text{C}$ ): GCM + projections

$$|\Phi_{I_c M_c}\rangle = \int d\beta f(\beta) |\Psi_{I_c M_c}(\beta)\rangle = \int d\beta f(\beta) \hat{P}_{M_c K_c}^{I_c} \hat{P}^N \hat{P}^Z |\Psi_{\text{MF}}(\beta)\rangle$$



- ✓ axial symmetry
- ✓ relativistic PC-F1



# Microscopic Particle-Rotor Model for $\Lambda$ hypernuclei

Example:  $^{13}_{\Lambda}\text{C}$

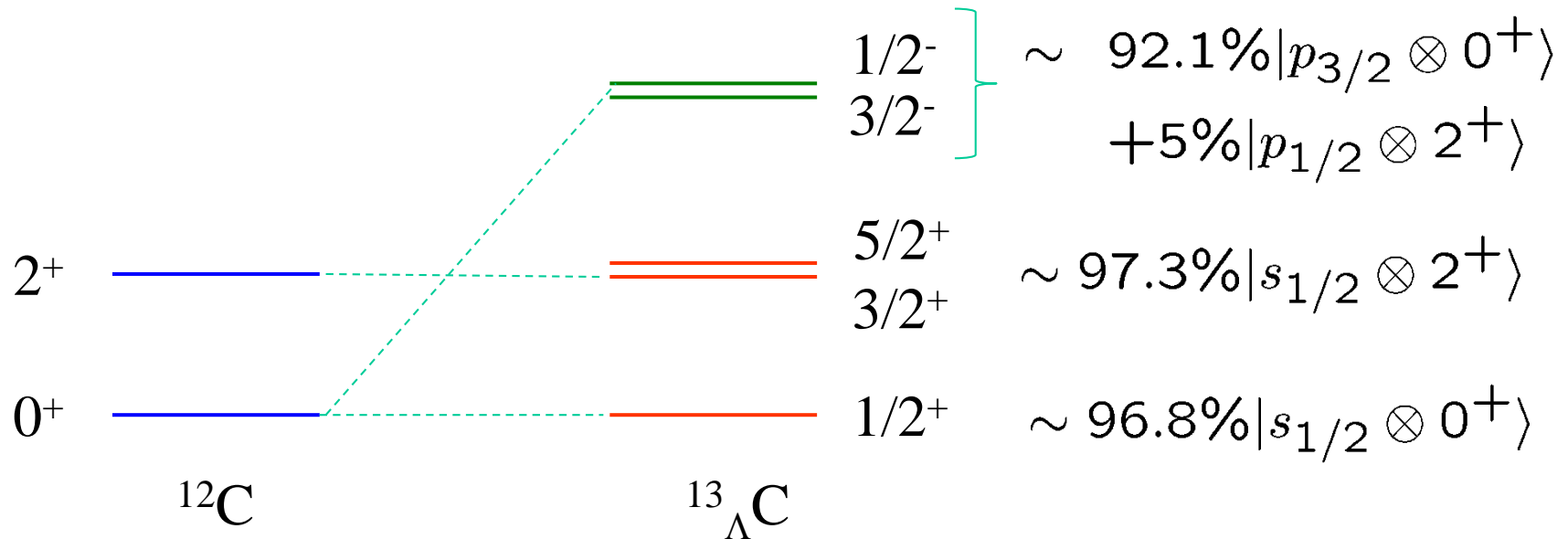
(i) beyond mean-field calculations for e-e core ( $^{12}\text{C}$ )

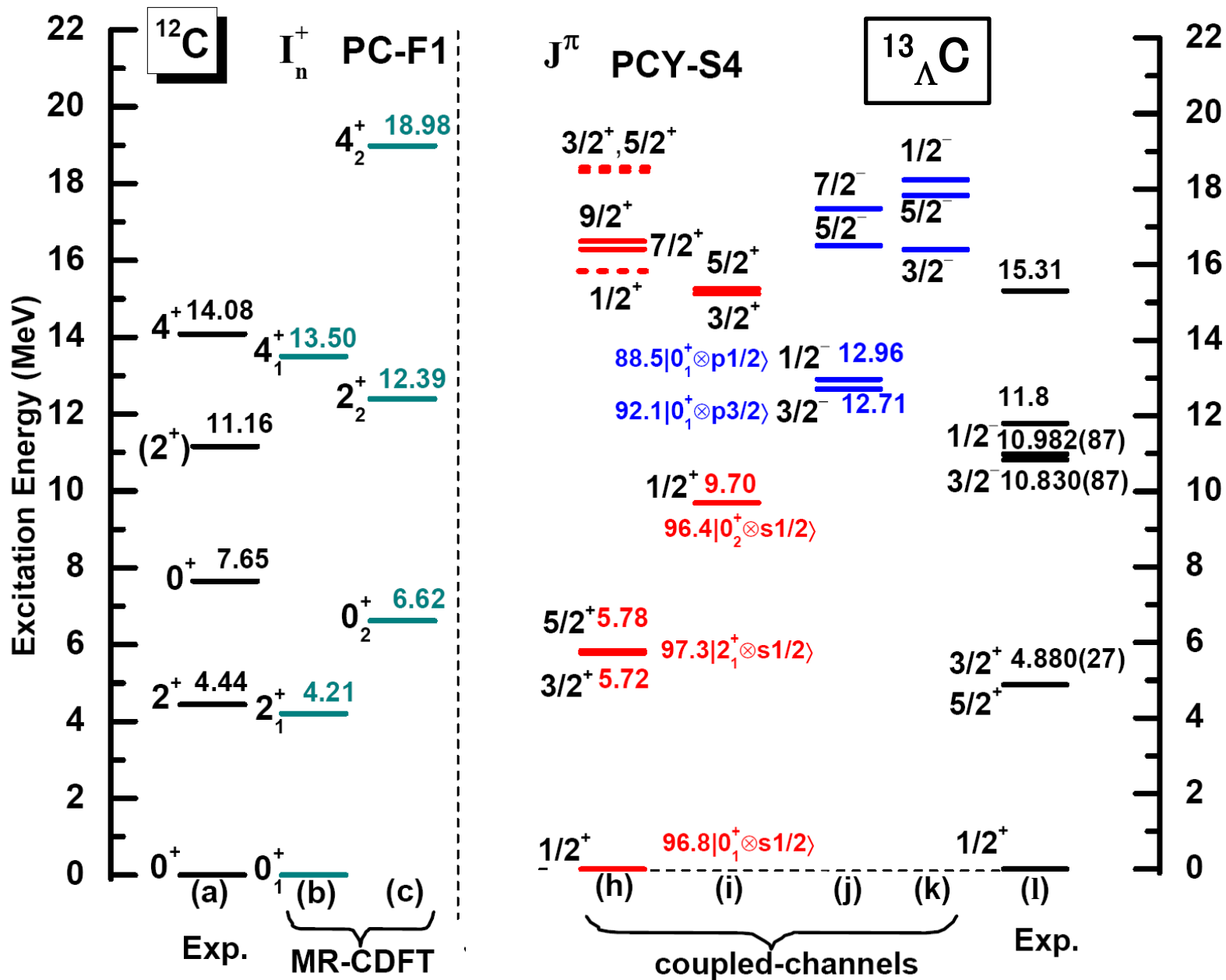
(ii) coupling of  $\Lambda$  to the core states

$$|\Phi_{IM}\rangle = \sum_{j,l,I_c} [\psi_{jl}(r_{\Lambda}) \otimes |\Phi_{I_c}\rangle]^{(IM)}$$

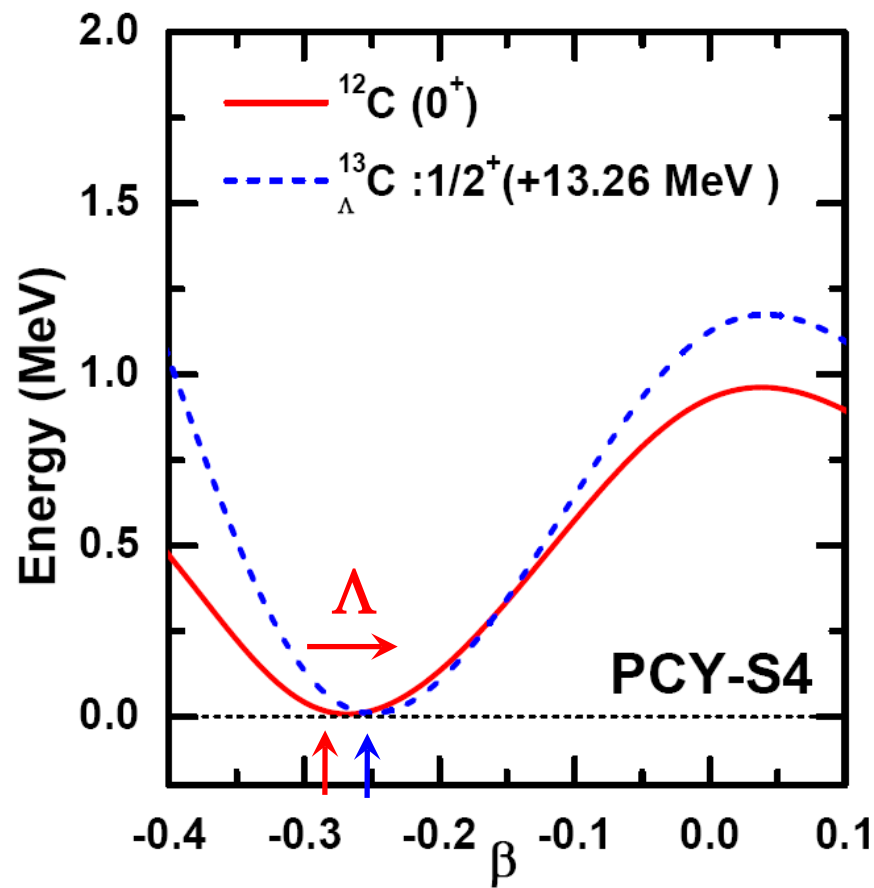
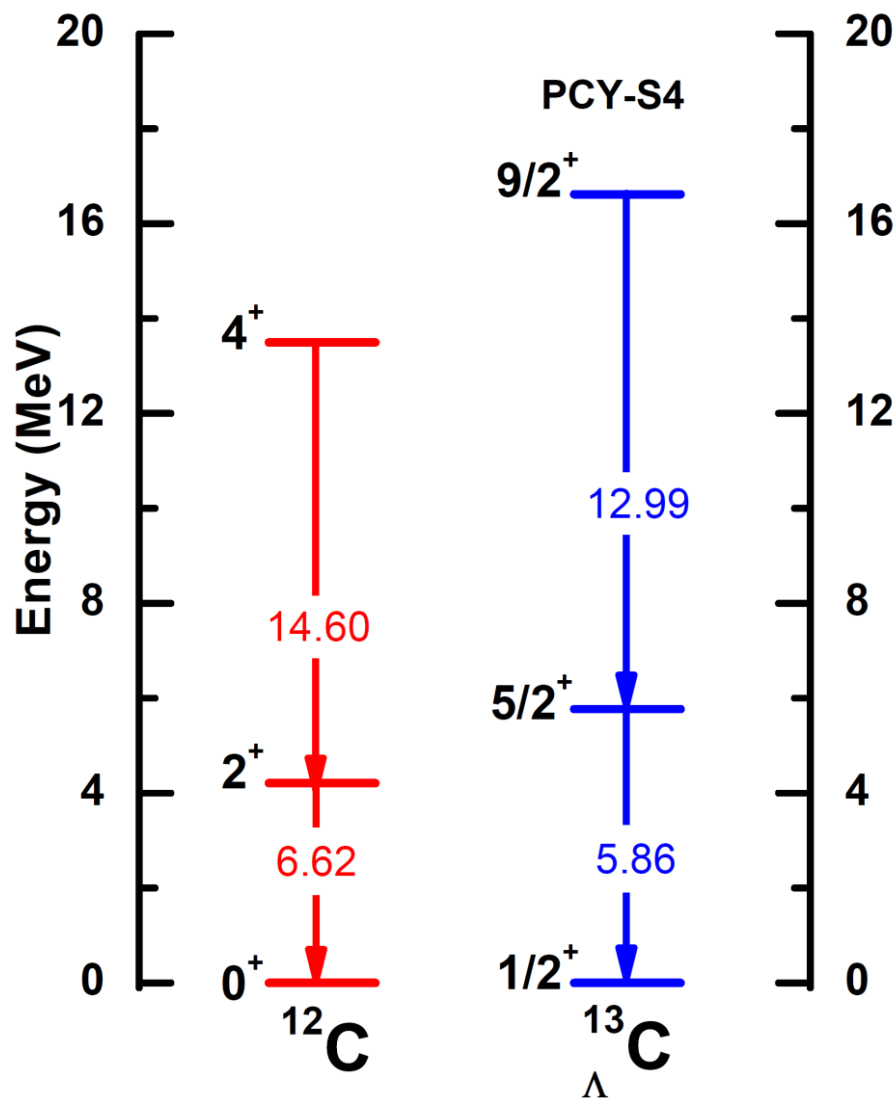
$\Lambda$

core





# B(E2) transition rates ( $e^2\text{fm}^4$ )

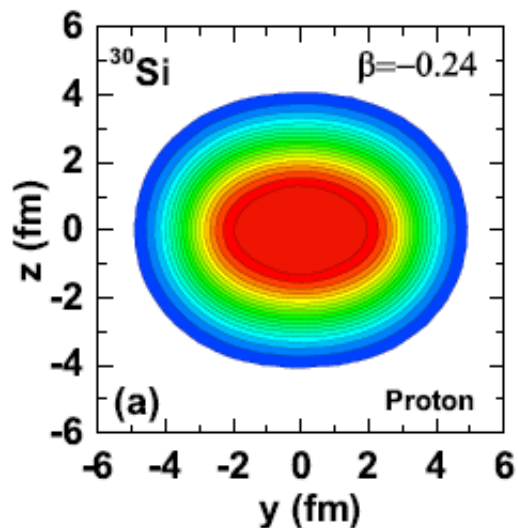


	$^{12}\text{C}$	$^{13}_{\Lambda}\text{C}$
$\beta$	-0.27	-0.25
$r_p$ (fm)	2.44	2.39

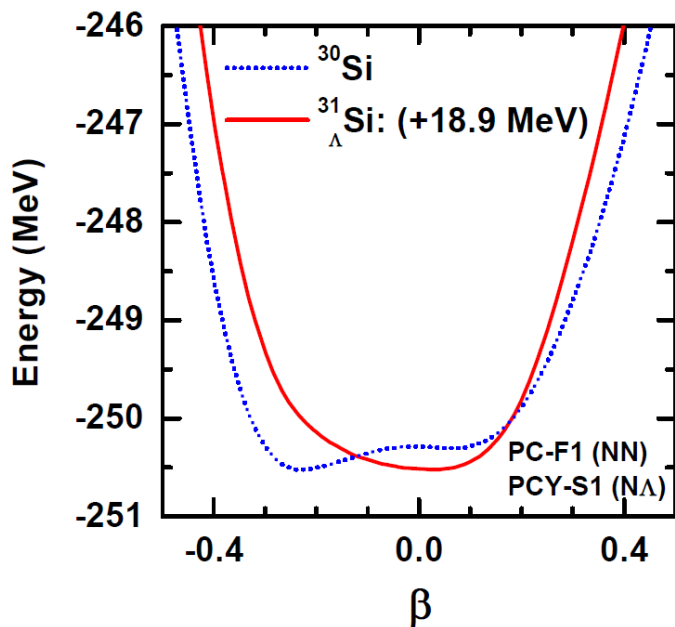
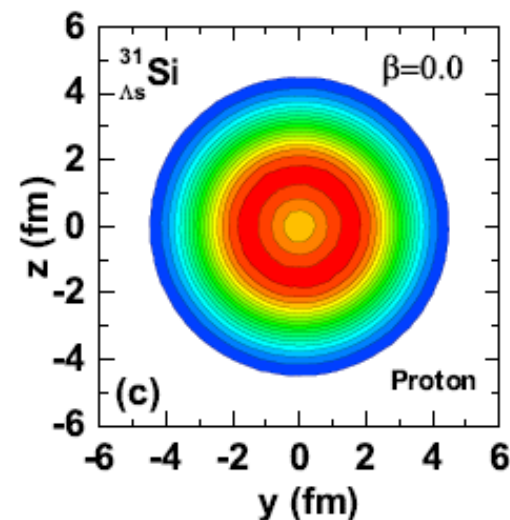
➤ B(E2) : ~ 11% reduction

# Beyond MF effect on deformation change in $^{31}_{\Lambda}\text{Si}$

Mean-field calculations:

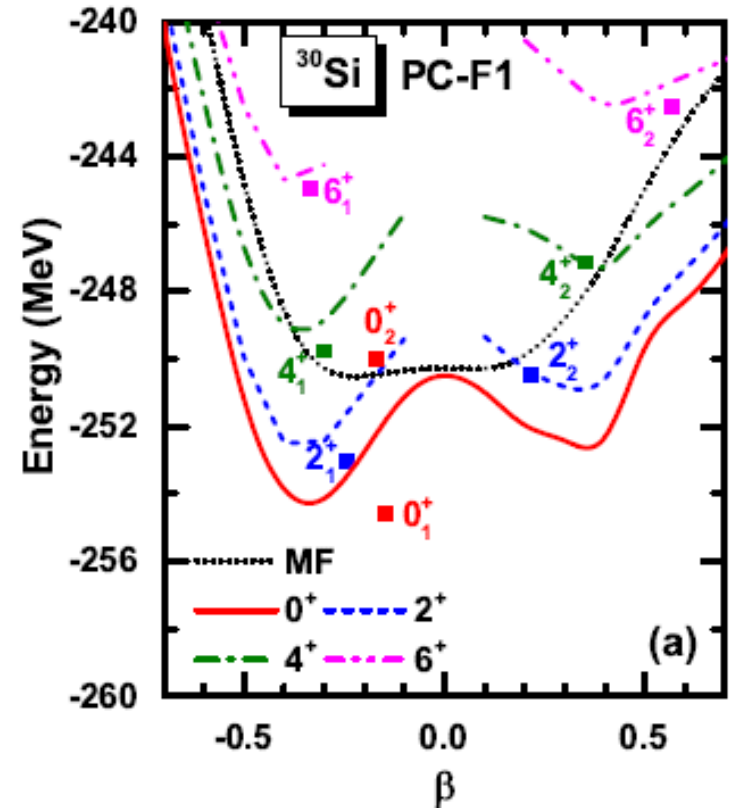
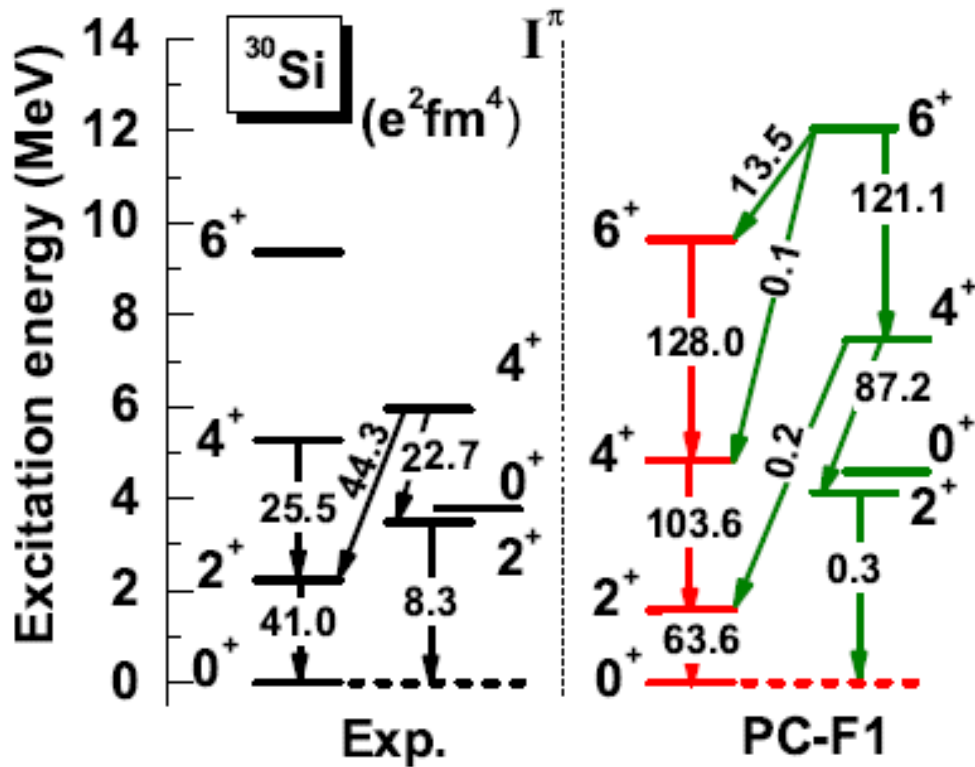


$\Lambda$   
oblate  $\rightarrow$  spherical



Soft potential energy surface  
 $\rightarrow$  beyond-mean-field effects?

# spectrum of $^{30}\text{Si}$



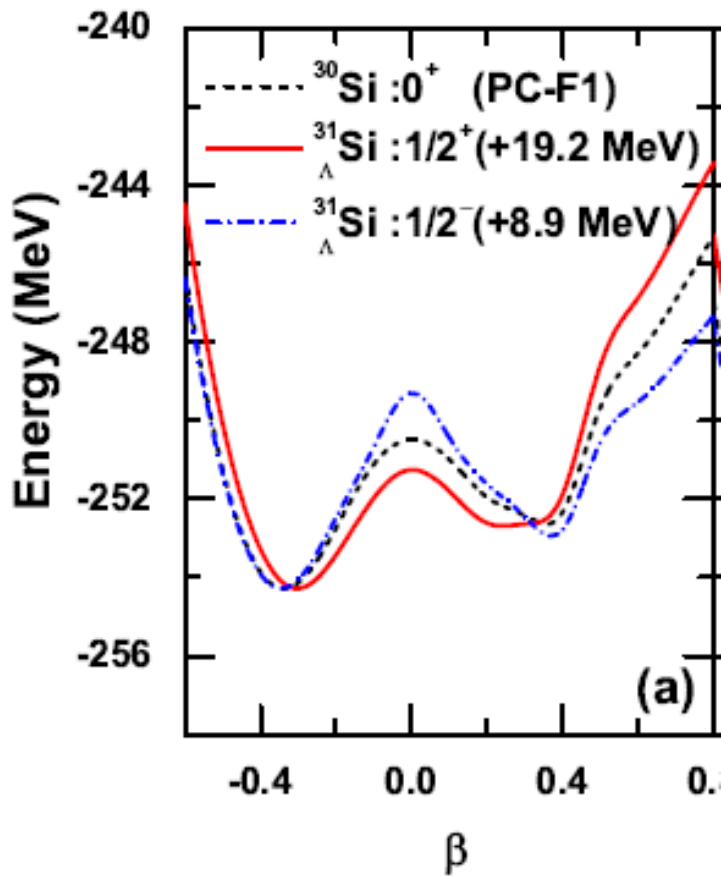
\* axial symmetry

$$R_{4/2} = E(4^+) / E(2^+) = 3.083$$

cf. Expt:  $R_{4/2} = 2.66$



# spectrum of $^{31}_{\Lambda}\text{Si}$

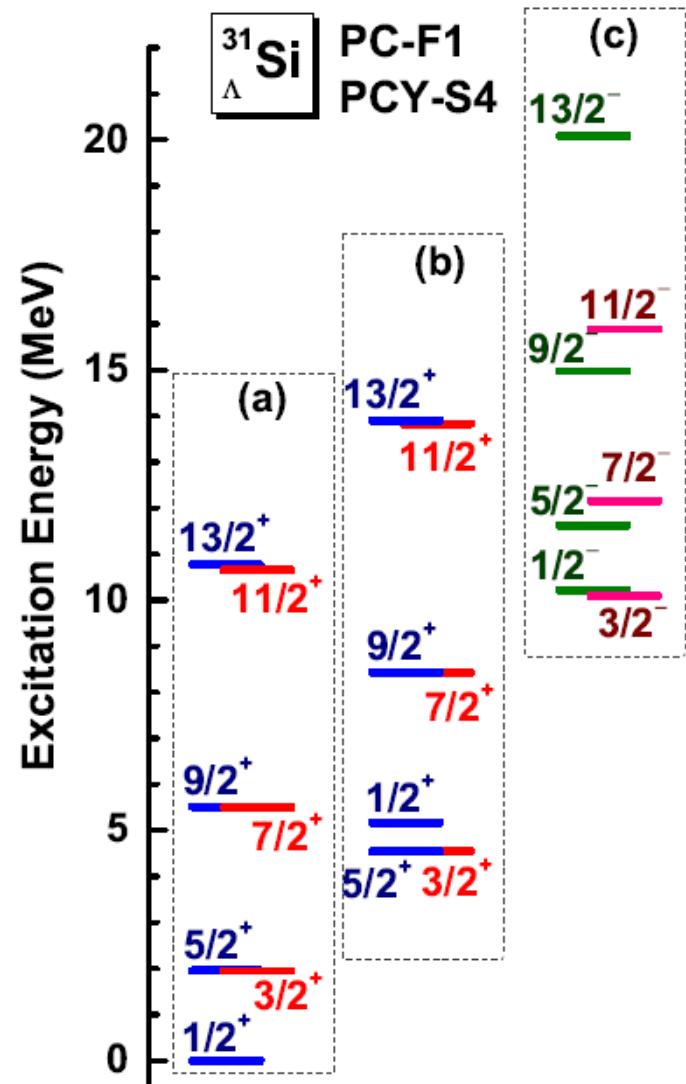


still rotational spectrum

[  $R_{4/2} = 3.08 \rightarrow 2.83$  ]

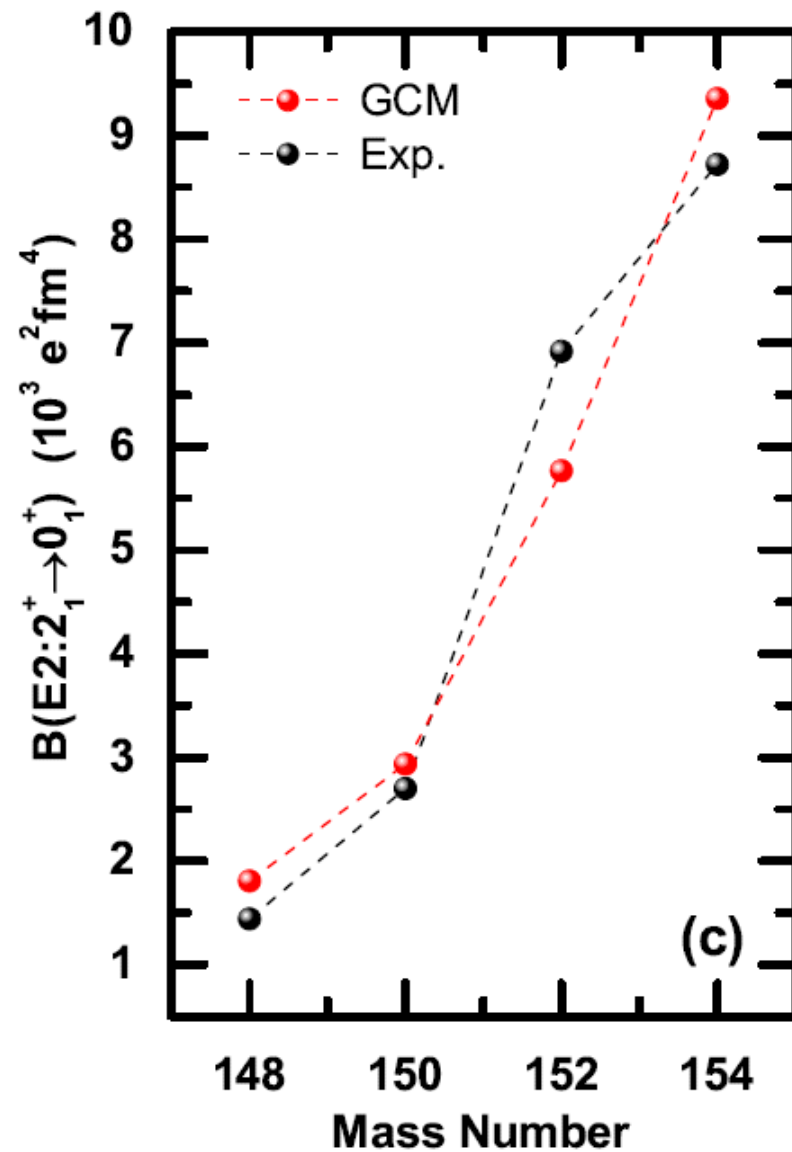
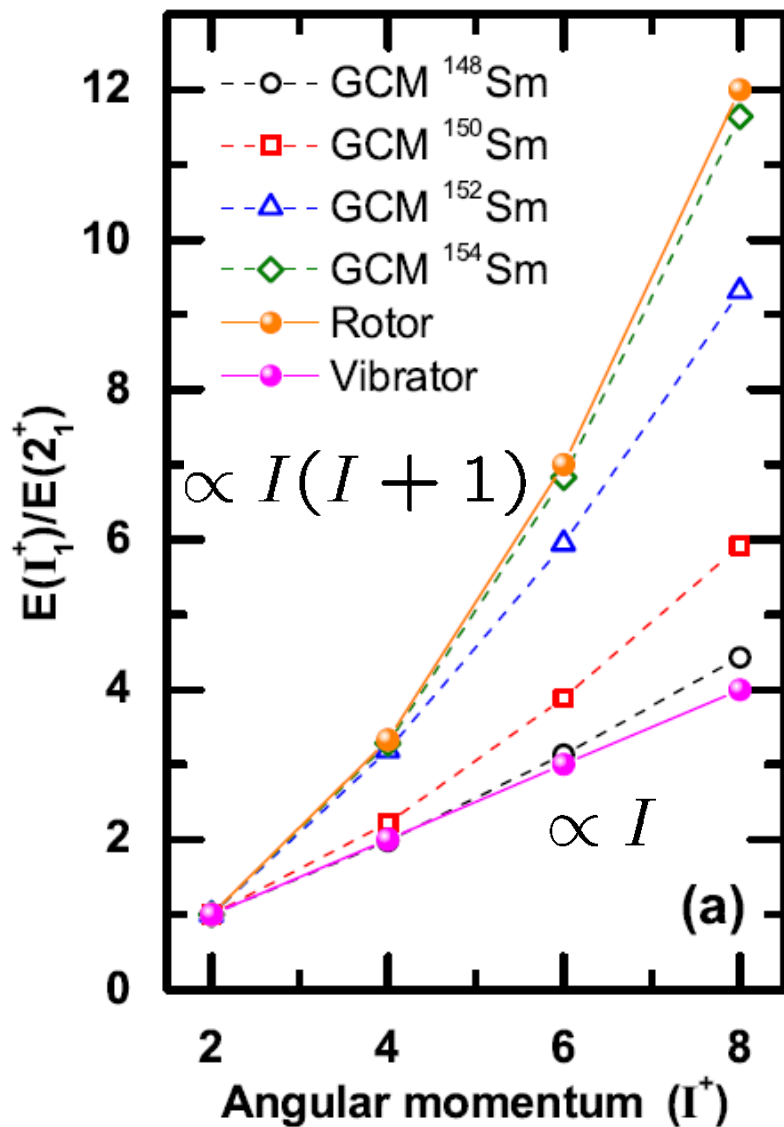
changes towards spherical, but not quite

→ large BMF effects



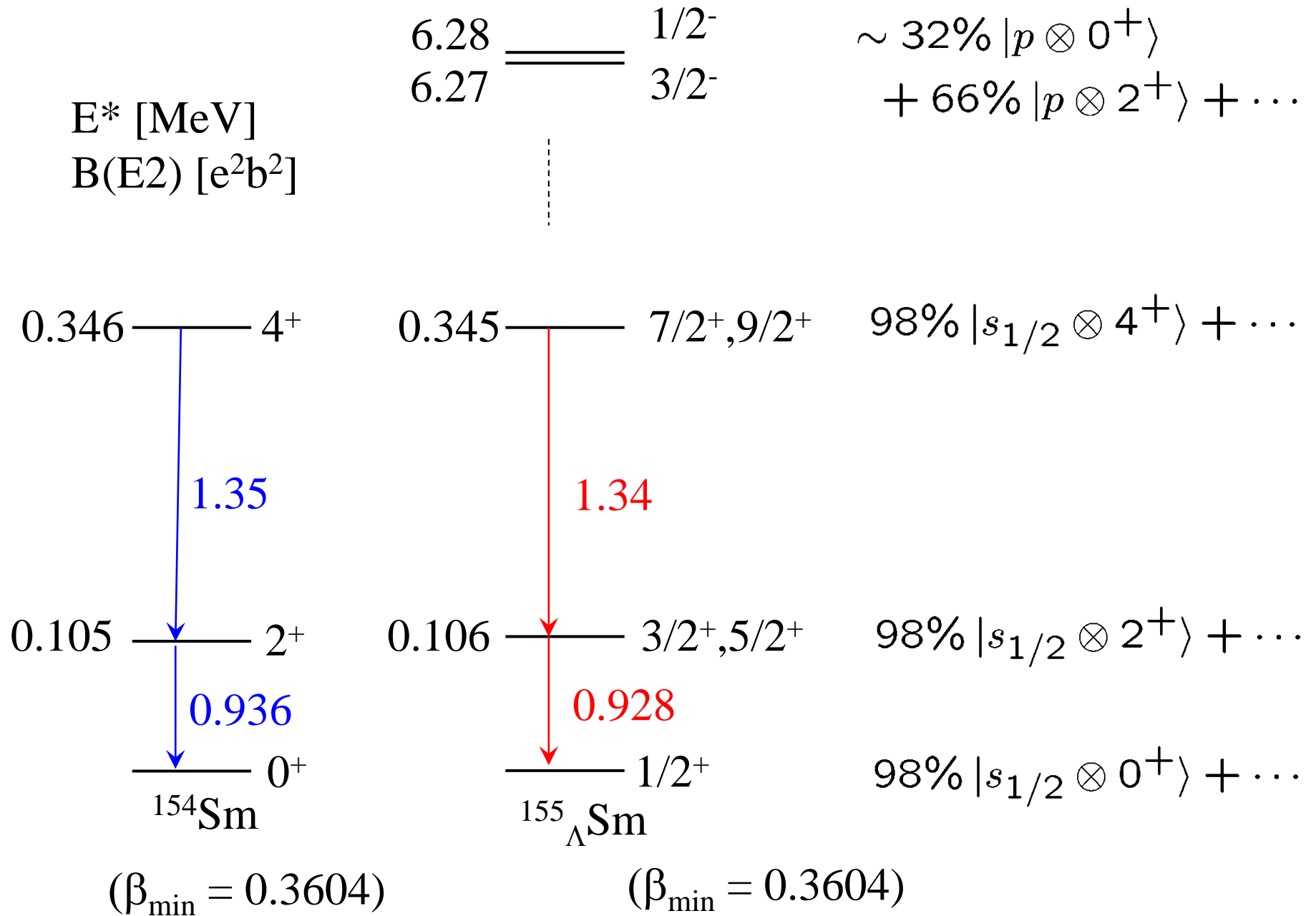
H. Mei, K.H., J.M. Yao, and T. Motoba,  
in preparation

# Application to Sm isotopes

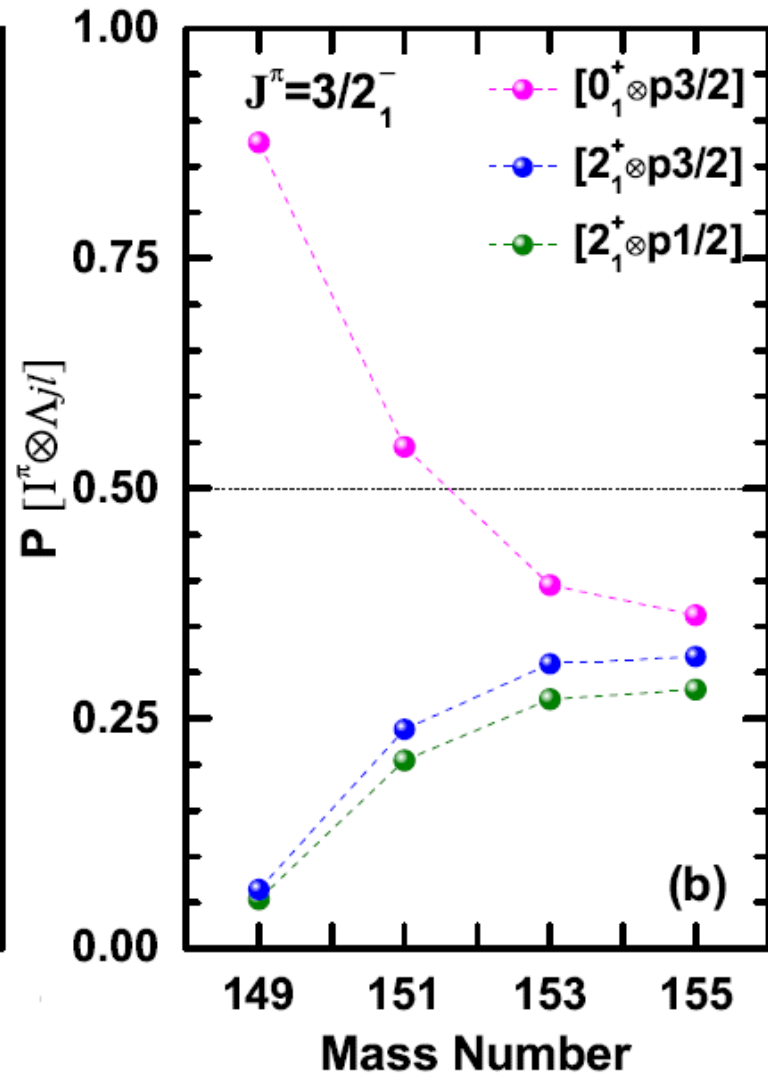
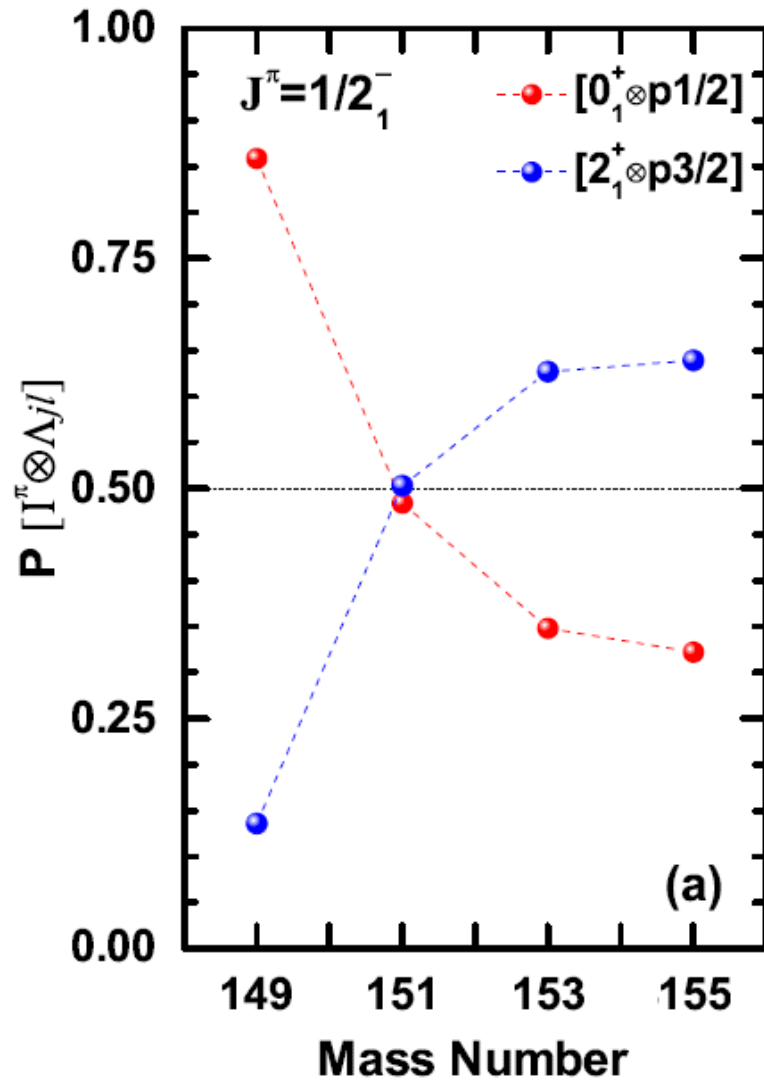


# Results for $^{155}_{\Lambda}\text{Sm}$

H. Mei, K.H., J.M. Yao, T. Motoba, PRC91('15) 064305



weak coupling  $\rightarrow$  strong coupling



H. Mei, K.H., J.M. Yao, and T. Motoba,  
in preparation

# Summary

## Applications of Beyond-Mean-Field method

### ➤ Heavy-ion subbarrier fusion reactions

- ✓ strong interplay between reaction and structure  
cf. fusion barrier distributions

### C.C. calculations with beyond-MF method

- ✓ anharmonicity
- ✓ octupole vibrations:  $^{16}\text{O} + ^{208}\text{Pb}$

### ➤ Low-lying spectra of $\Lambda$ hypernuclei

#### Microscopic particle-rotor model

- ✓  $\Lambda$  + GCM states for core
- ✓ first calculation for low-lying spectrum based on mean-field type calculations
- ✓ re-examination of “disappearance of deformation” in MF calc.

